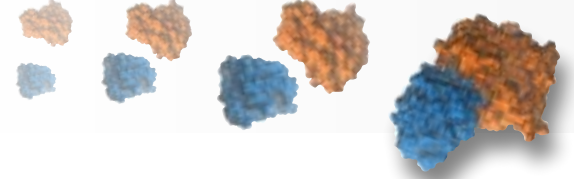


TRANSITION-PATH SAMPLING AND FREE-ENERGY CALCULATIONS

Chris Chipot

*Laboratoire International Associé CNRS UIUC,
Unité Mixte de Recherche n° 7565, Université de Lorraine*

*Beckman Institute for Advanced Science and Technology,
Department of Physics
University of Illinois at Urbana-Champaign*



INTRODUCTION

Why do we need reaction coordinates?

WHAT IS A GOOD REACTION-COORDINATE MODEL ?

- Reaction coordinate versus order parameter
- Committor distributions

THE STRING METHOD

- The basic string method
- The string method with swarms of trajectories

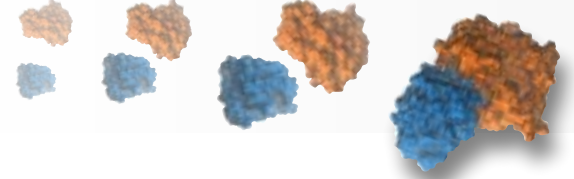
EXTRACTING THE FREE ENERGY FROM THE MINIMUM-ACTION PATH

- Path-collective variables
- Perturbative approach

APPLICATIONS

- Transition path of activation loop in c-Src kinase
- Chemomechanical coupling in V_1 -ATPase

RECONCILING THERMODYNAMICS AND KINETICS



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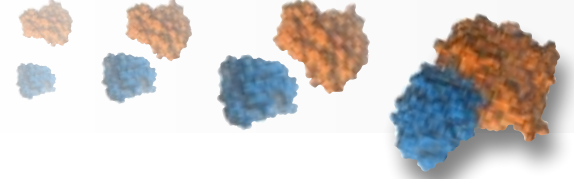
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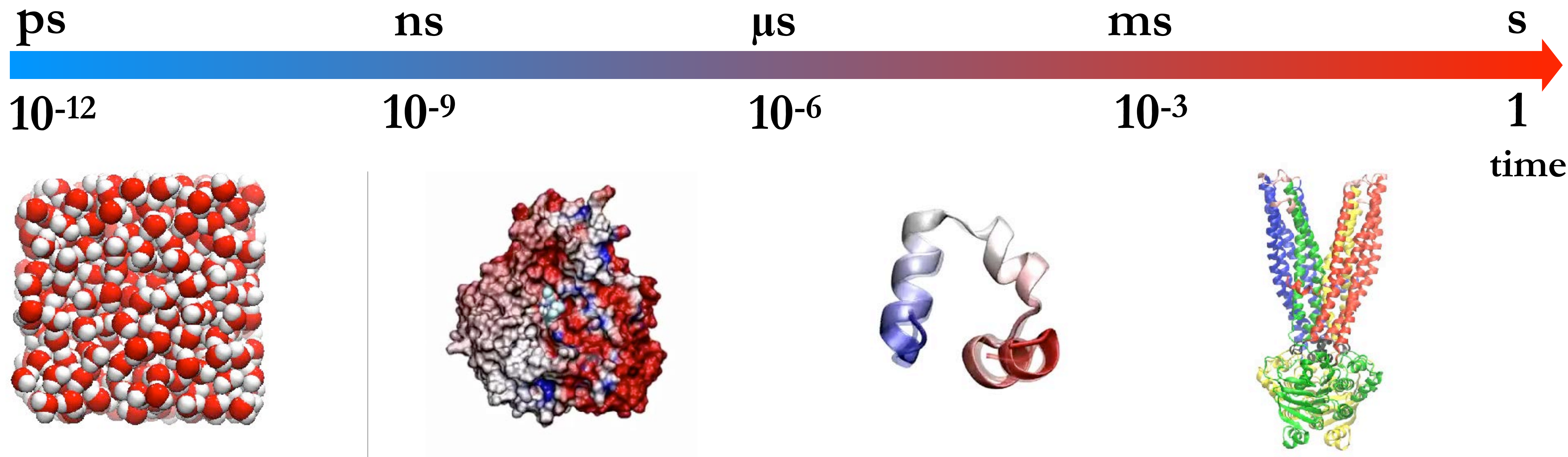
RECONCILING THERMODYNAMICS AND KINETICS



WHY DO WE NEED REACTION COORDINATES?

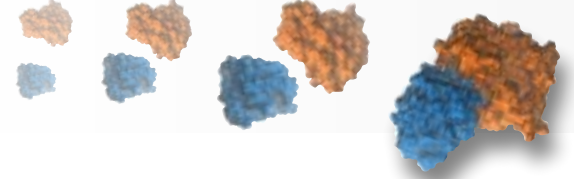
Capture both the kinetics and the thermodynamics of complex chemical and biological processes

Intricate transitions between metastable states often mirror substantial collectivity and rugged free-energy landscapes



De Donder, T. *L'affinité*. Gauthier-Villars: Paris, 1927

Kirkwood, J. G. *J. Chem. Phys.* **1935**, *3*, 300–313



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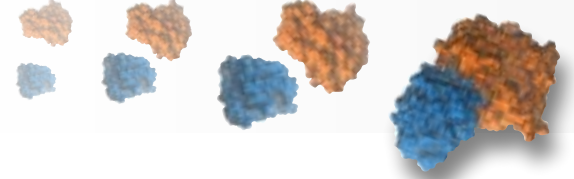
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WHAT IS A GOOD REACTION-COORDINATE MODEL ?



The *true* reaction coordinate generally refers to a unique mathematical object on \mathbb{R}^{3N} .

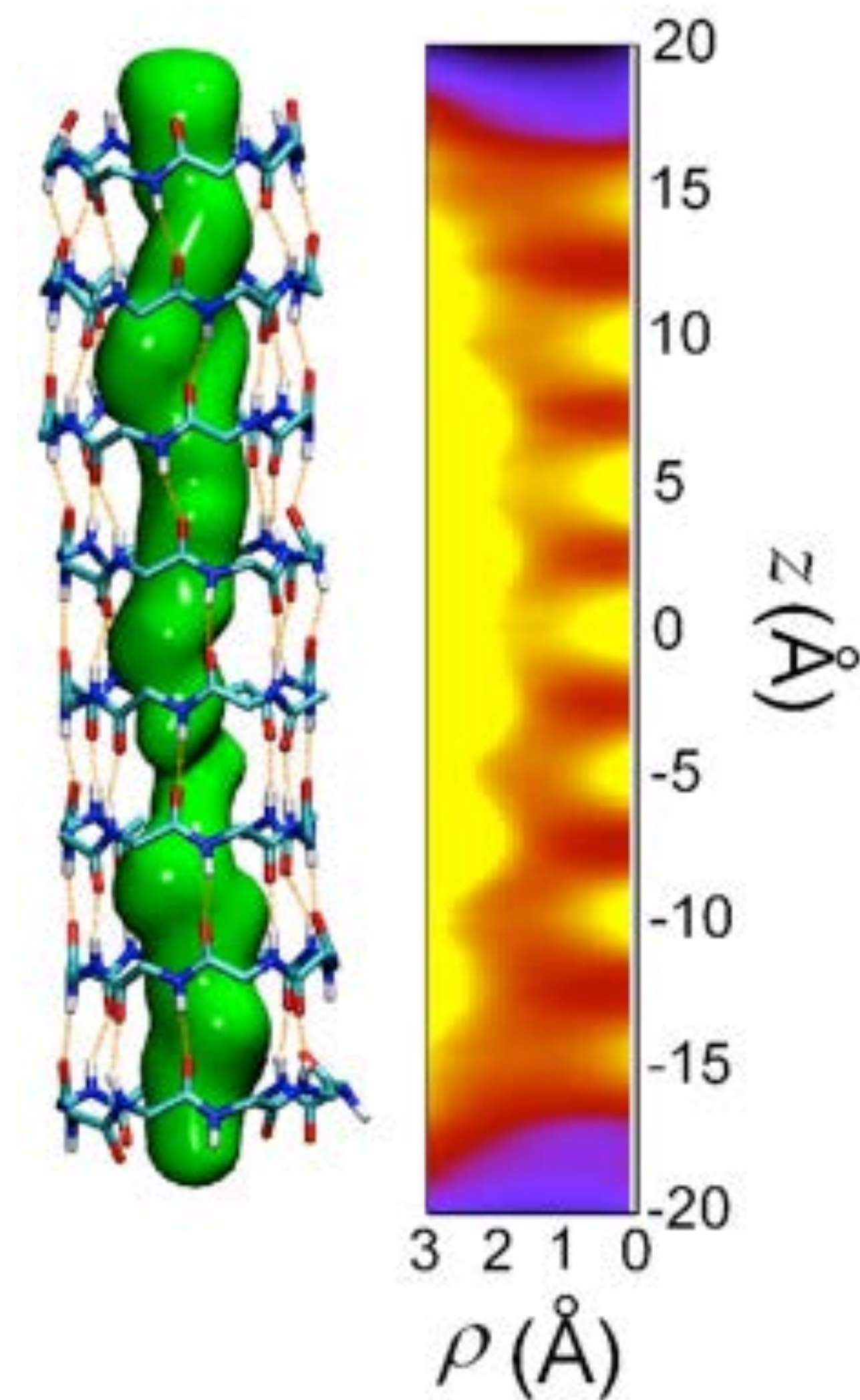
It defines the minimum free-energy pathway connecting the reference state to the target state of the transformation.

In practice, we coarse-grain the atomic detail:

$$\{x_1, x_2, \dots, x_N\} \longrightarrow \{z_1, z_2, \dots, z_n\}, \text{ where } n \ll N$$

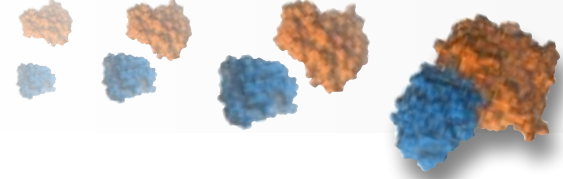
Cartesian
coordinates

collective
variables



Bolhuis, P. G.; Dellago, C.; Chandler, D. *Proc. Natl. Acad. Sci. U. S. A.* **2000**, *97*, 5877-5882

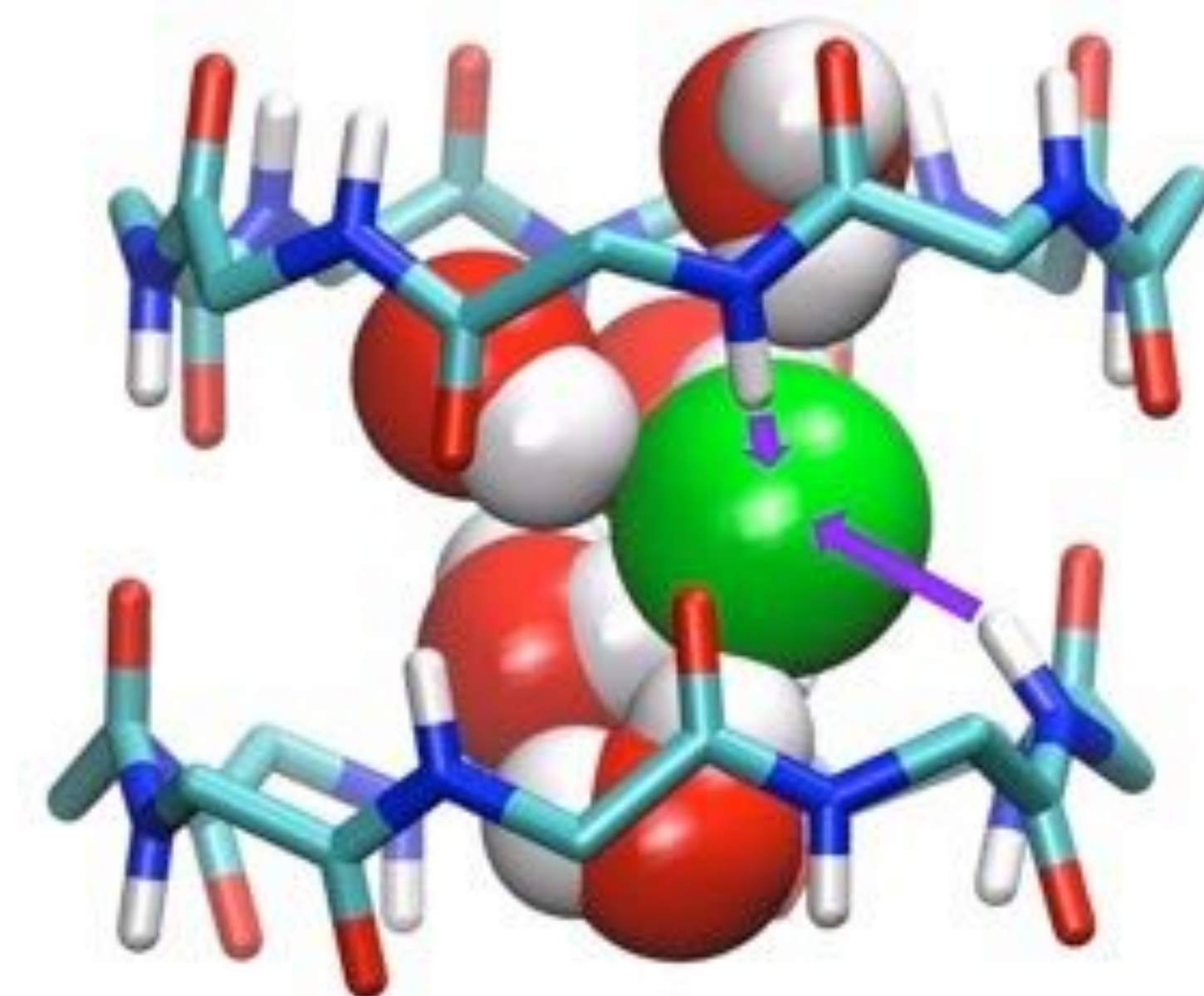
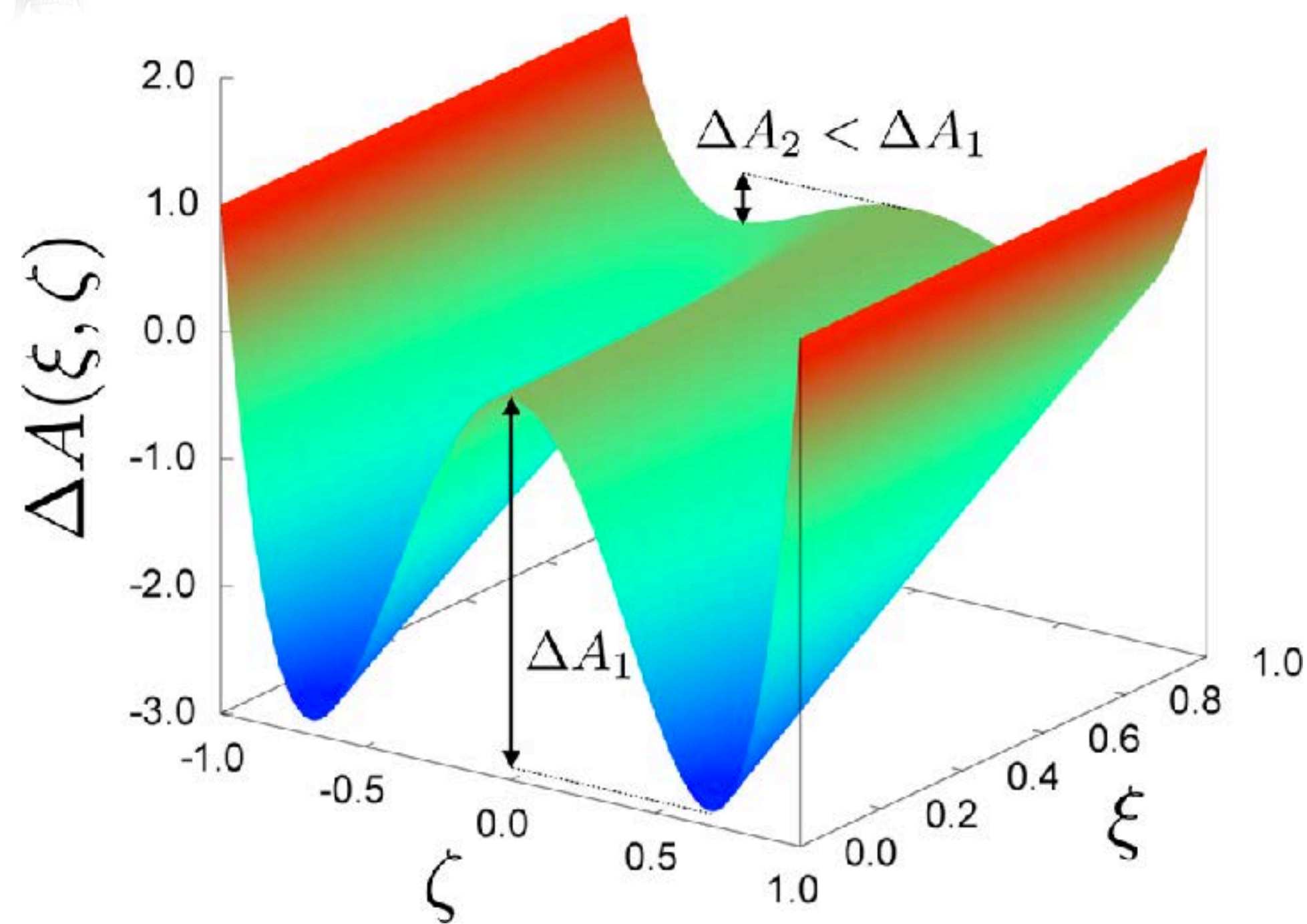
Bolhuis, P. G.; Chandler, D.; Dellago, C.; Geissler, P. *Ann. Rev. Phys. Chem.* **2002**, *59*, 291-318



WHAT IS A GOOD REACTION-COORDINATE MODEL ?



A one-dimensional order parameter, namely the long axis of the cavity, is not enough to describe ion conduction in a synthetic channel.

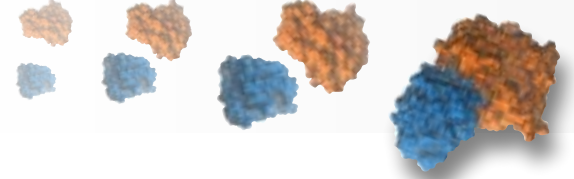


Stringent assumption of a one-dimensional geometric variable: Averaging of all other, fast and slow, degrees of freedom, which could not be further from the truth.

Bolhuis, P. G.; Dellago, C.; Chandler, D. *Proc. Natl. Acad. Sci. U. S. A.* **2000**, *97*, 5877-5882

Bolhuis, P. G.; Chandler, D.; Dellago, C.; Geissler, P. *Ann. Rev. Phys. Chem.* **2002**, *59*, 291-318

Chipot, C.; Lelièvre, T. *SIAM J. Appl. Math.* **2011**, *71*, 1673-1695



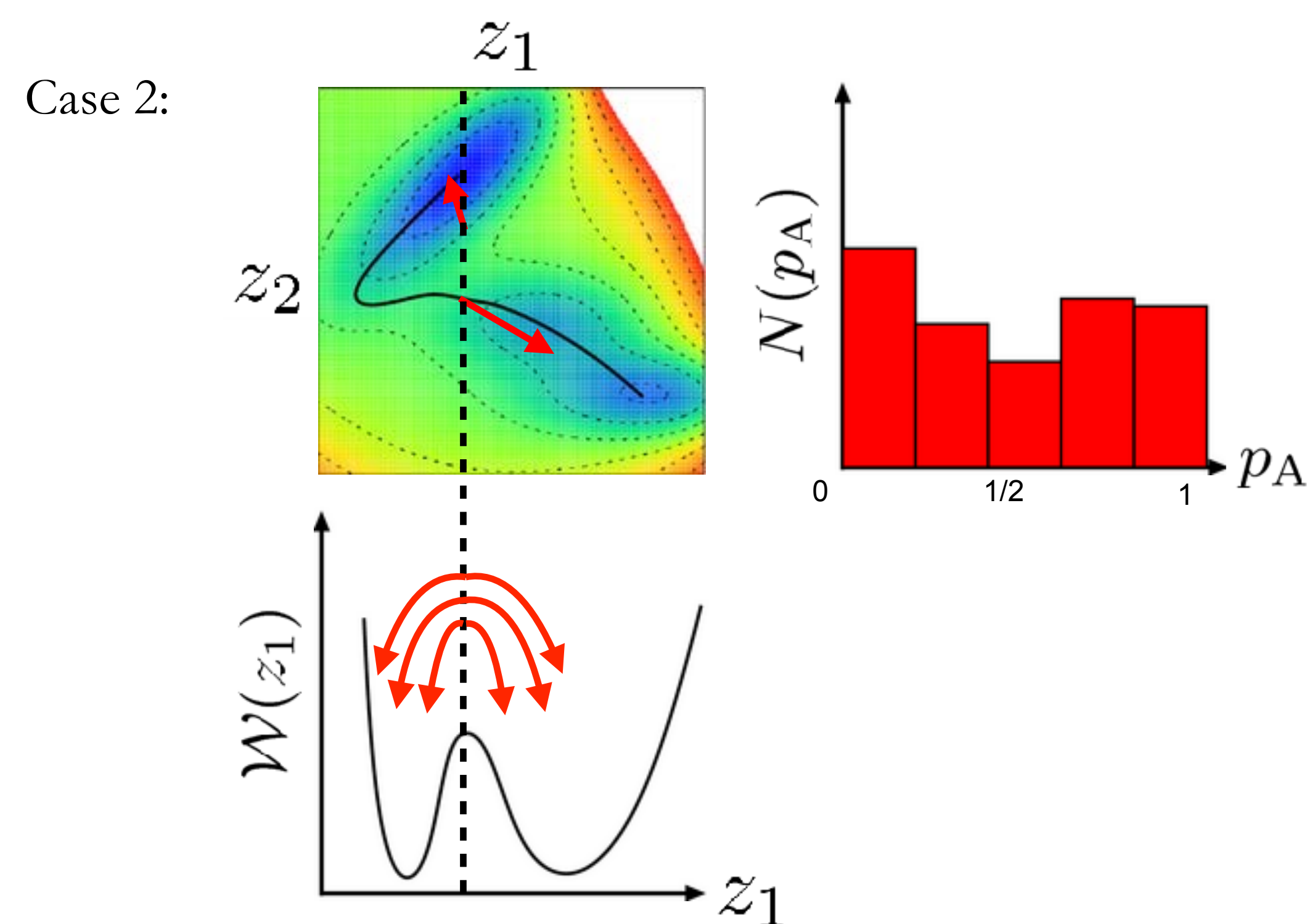
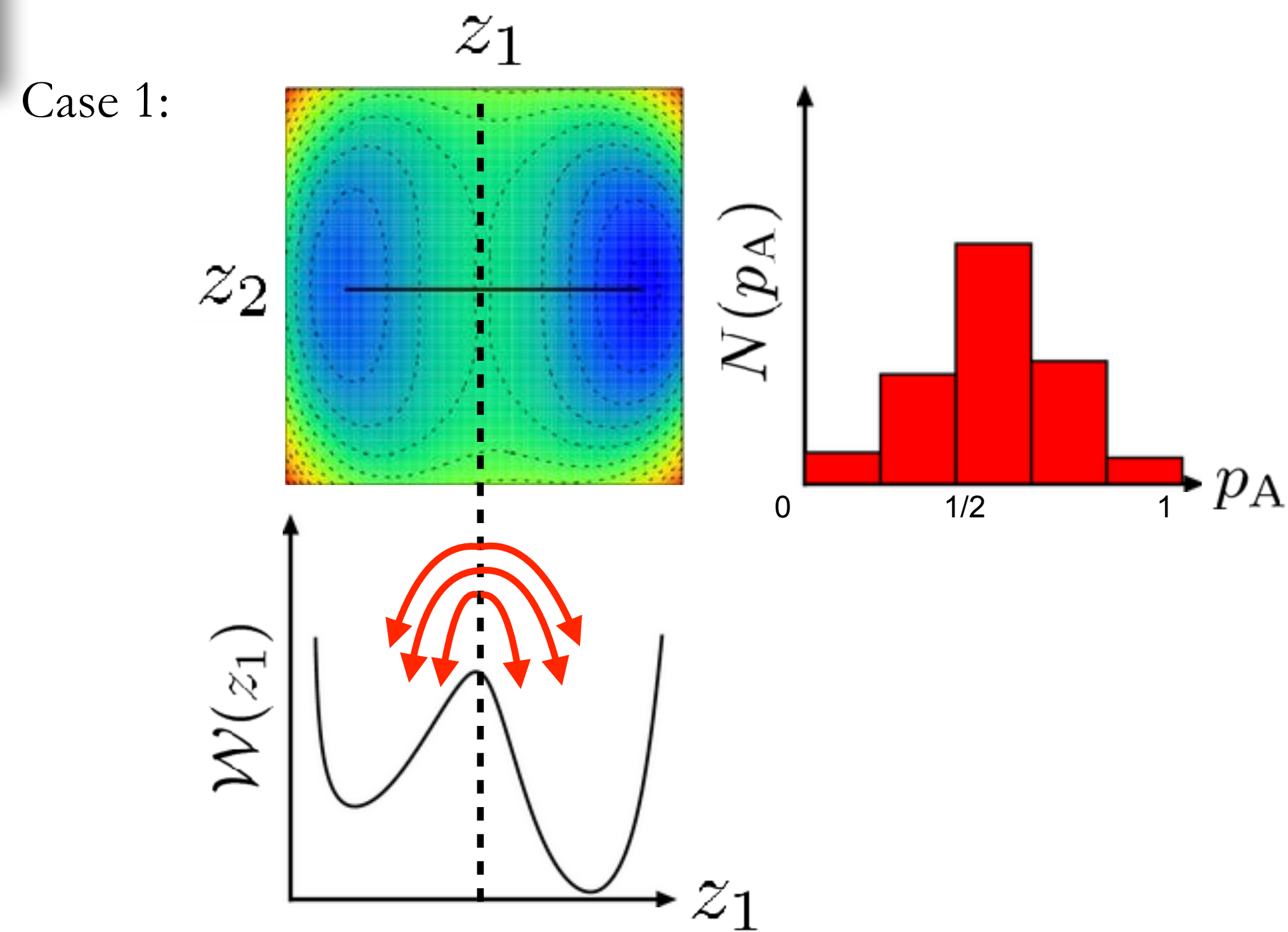
WHAT IS A GOOD REACTION-COORDINATE MODEL ?



Including relevant collective variables is absolutely crucial for finding true dynamical pathways.

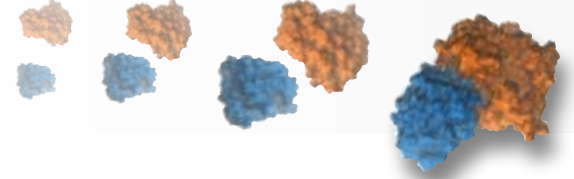
Committer — The probability to reach the target state before returning to the reference state.

The transition state surface for $A \rightarrow B$ transitions is formed of configurations where $p_B = p_A = 1/2$.



Bolhuis, P. G.; Dellago, C.; Chandler, D. *Proc. Natl. Acad. Sci. U. S. A.* **2000**, *97*, 5877-5882

Bolhuis, P. G.; Chandler, D.; Dellago, C.; Geissler, P. *Ann. Rev. Phys. Chem.* **2002**, *59*, 291-318



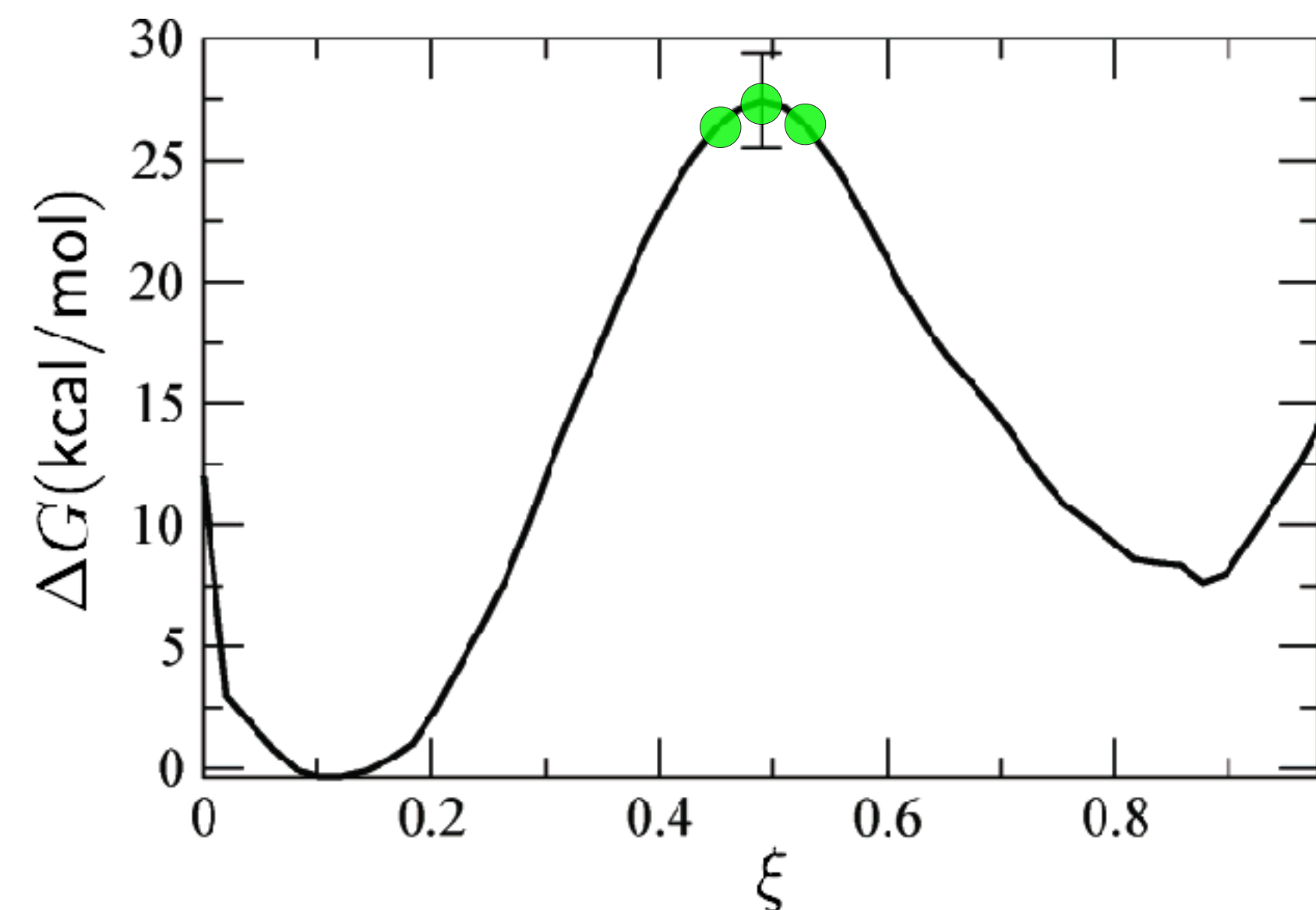
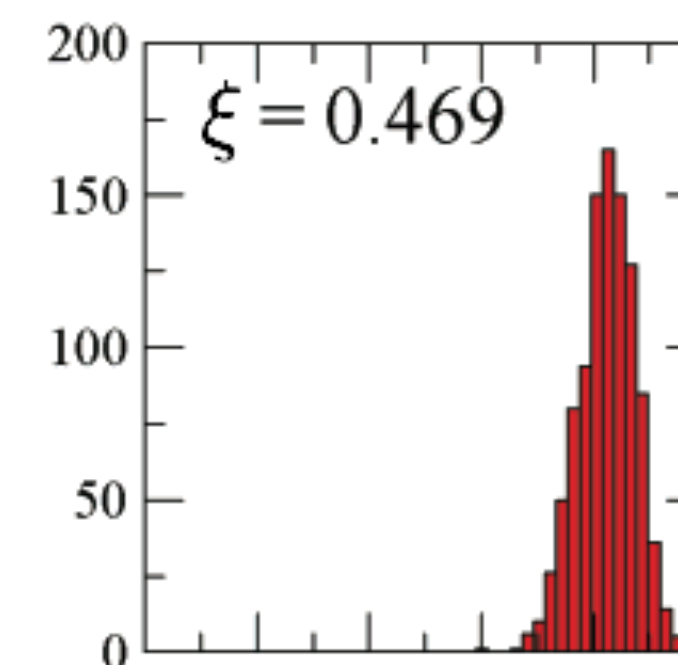
WHAT IS A GOOD REACTION-COORDINATE MODEL ?



While the choice of the reaction-coordinate model does not impact the thermodynamics of the process at hand, it modulates its kinetics.



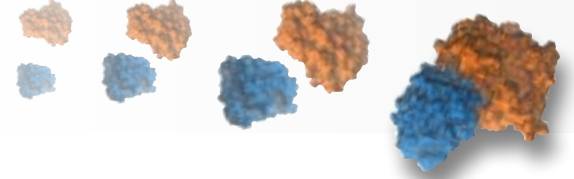
Option: Determine $N(p_A)$, the distribution of the committor probability, p_A , for the model of the reaction coordinate, ξ .



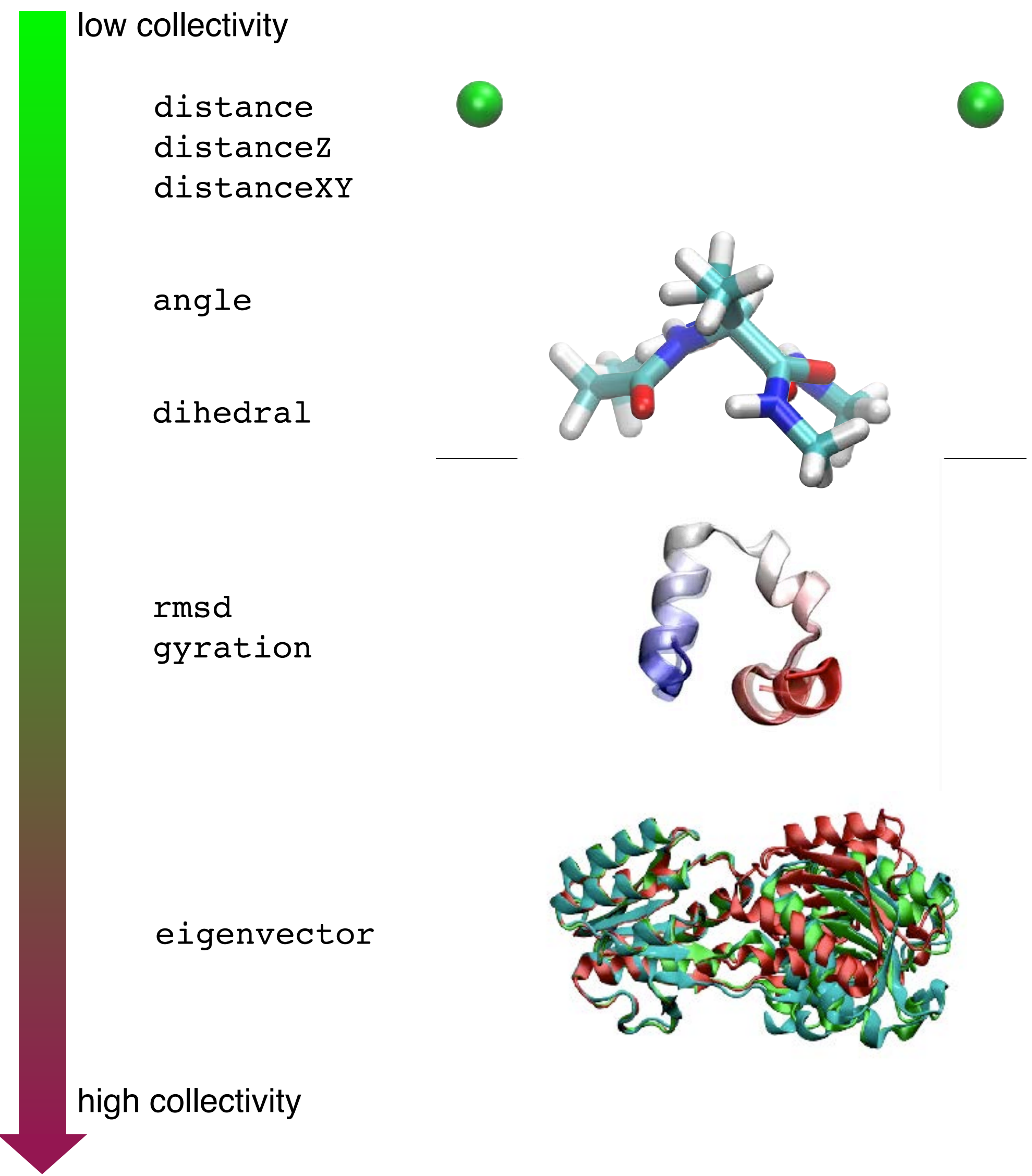
Run a series of molecular dynamics simulations from the putative maximum of the free-energy barrier and infer $N(p_A)$.

Bolhuis, P. G.; Dellago, C.; Chandler, D. *Proc. Natl. Acad. Sci. U. S. A.* **2000**, *97*, 5877-5882

Bolhuis, P. G.; Chandler, D.; Dellago, C.; Geissler, P. *Ann. Rev. Phys. Chem.* **2002**, *59*, 291-318



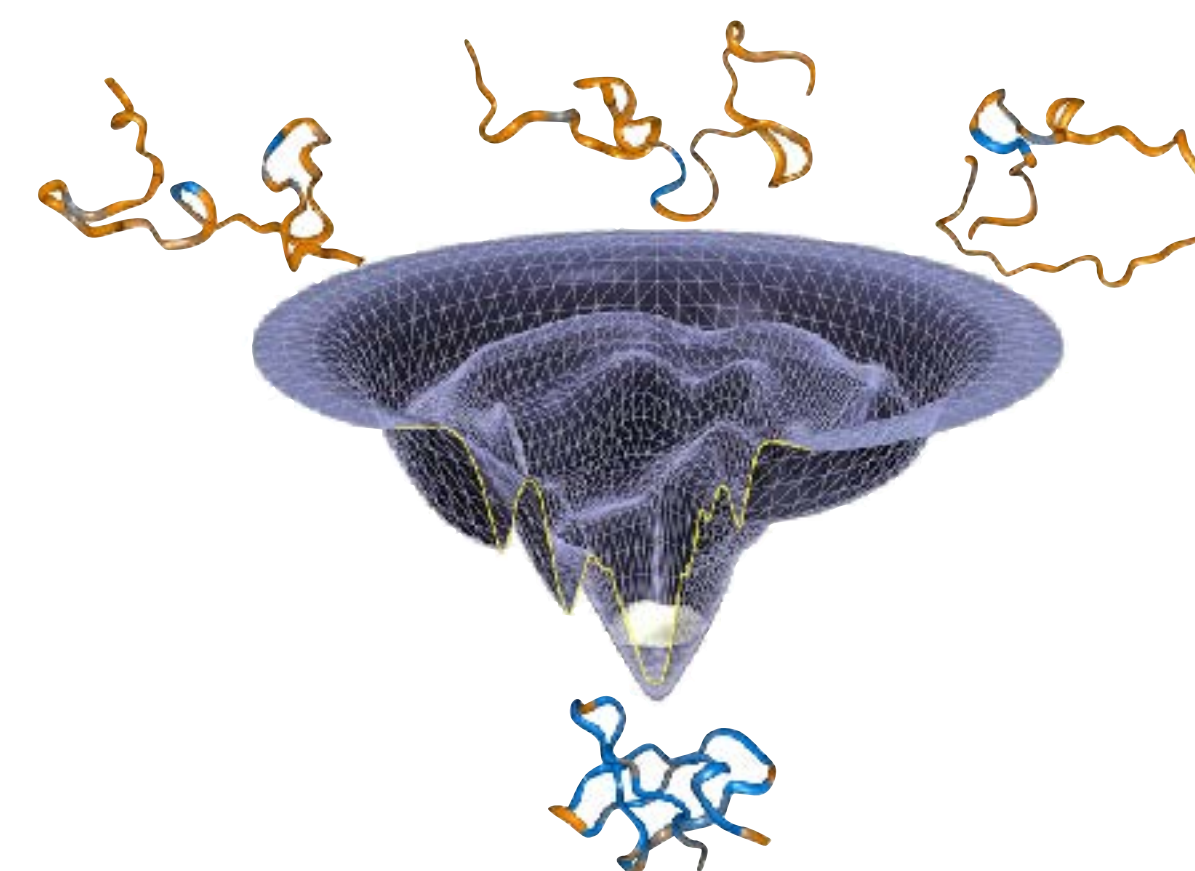
WHAT IS A GOOD REACTION-COORDINATE MODEL ?



Possible linear combination of variables

Degenerate variable

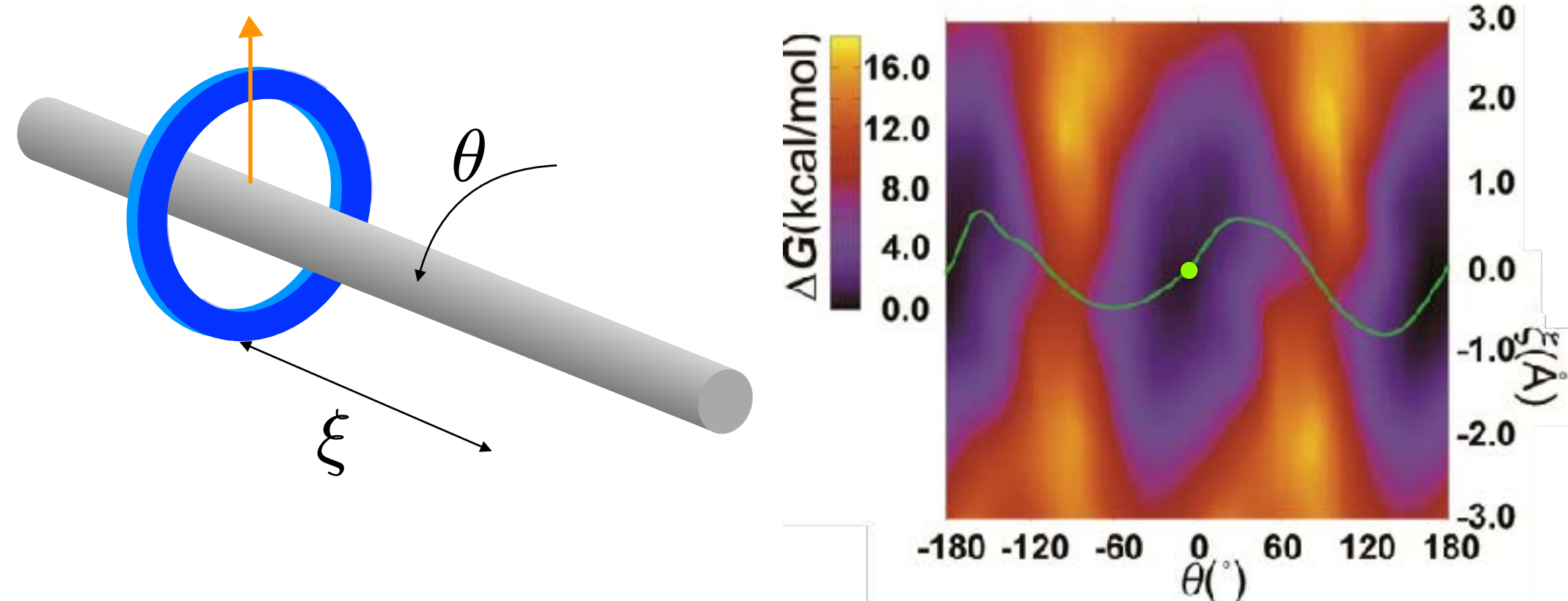
From normal mode or principal component analysis



Hénin, J.; Forin, G.; Chipot, C.; Klein, M. L. *J. Chem. Theor. Comput.* **2010**, *6*, 35-47

Fiorin, G.; Klein, M. L.; Hénin, J. *Mol. Phys.* **2013**, *111*, 3345-3362

WHAT IS A GOOD REACTION-COORDINATE MODEL ?



Movements in molecular objects can be more complex than suggested by chemical intuition.

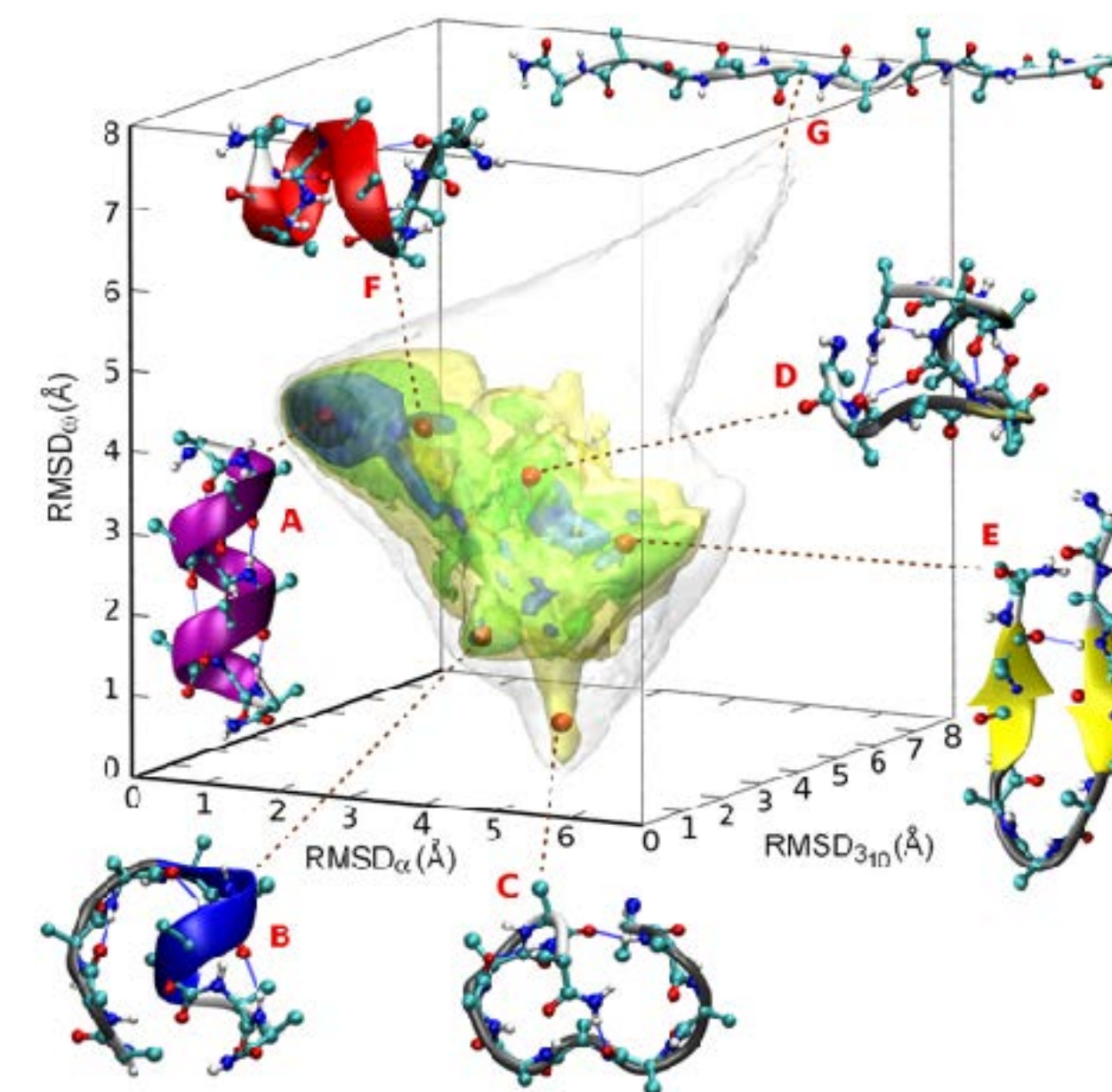
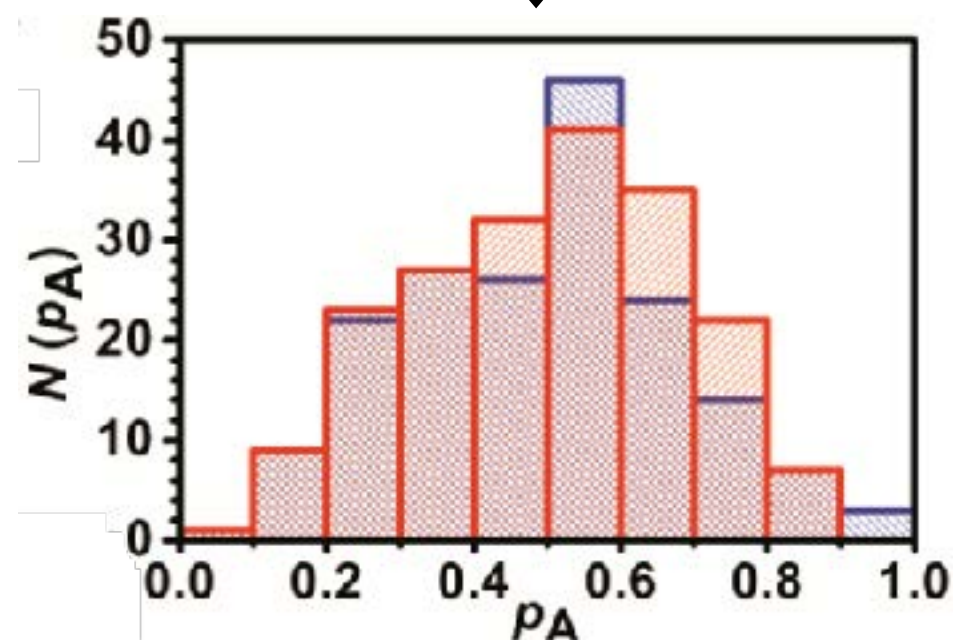
Define reaction coordinate model based on chemical intuition

Ascertain that the reaction coordinate model is a committor function

Increase dimensionality of the model

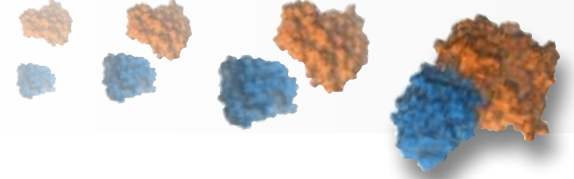
Turn to ergodic-sampling algorithms

Search for a minimum-action path



Liu, P.; Shao, X.; Chipot, C.; Cai, W. *Chem. Sci.* **2015**

Bolhuis, P. G.; Dellago, C.; Chandler, D. *Proc. Natl. Acad. Sci. U. S. A.* **2000**, *97*, 5877-5882



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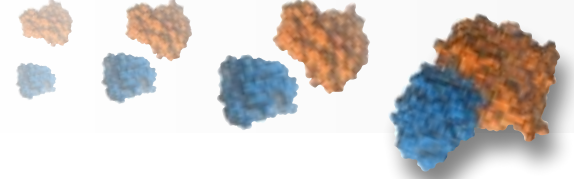
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RECONCILING THERMODYNAMICS AND KINETICS



THE BASIC STRING METHOD



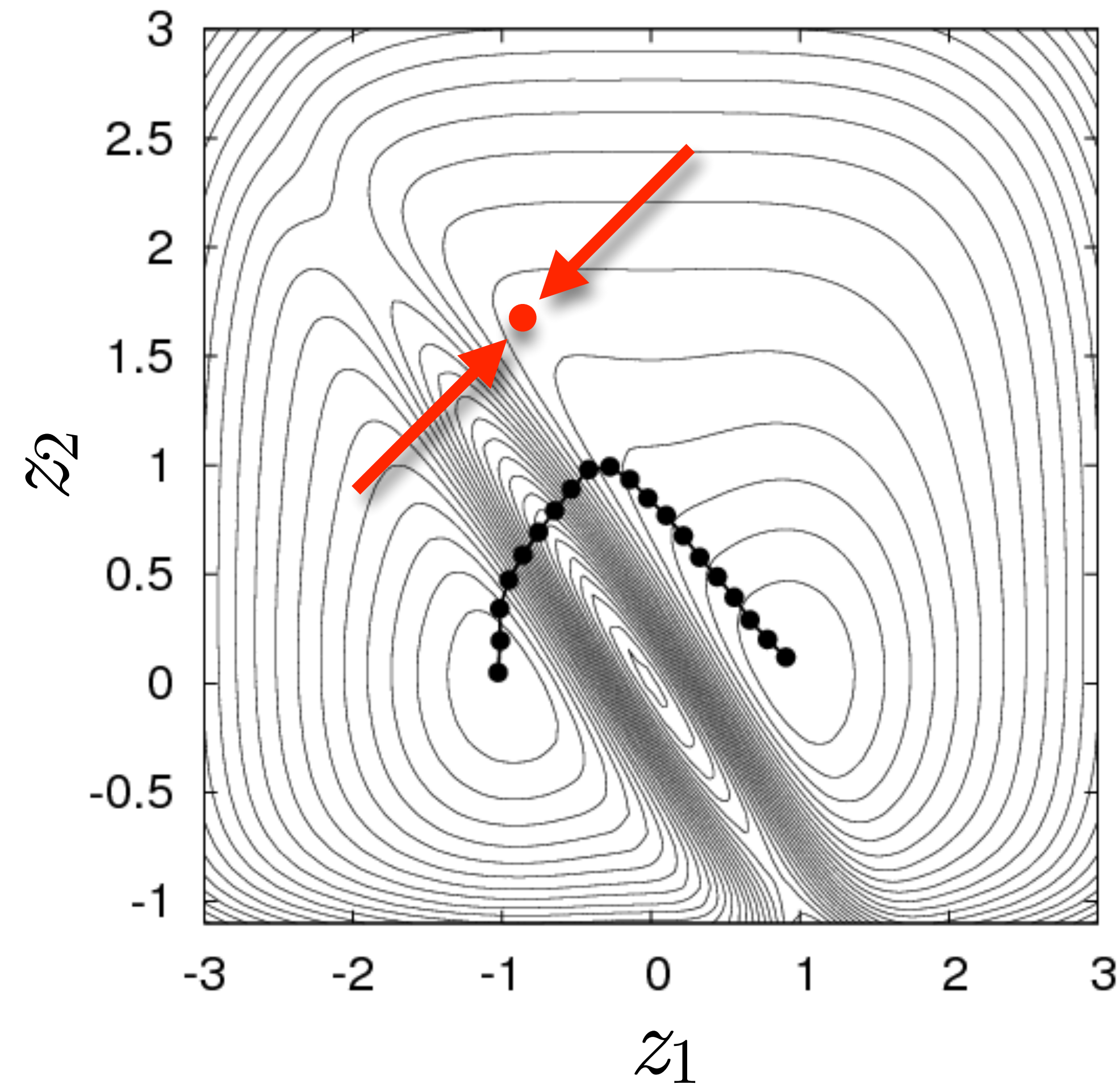
Let us assume some minimum-action (most probable) transition path connecting two basins of a free-energy landscape defined by a set of collective variables \mathbf{z} . Let us also consider the potential of mean force along this path,

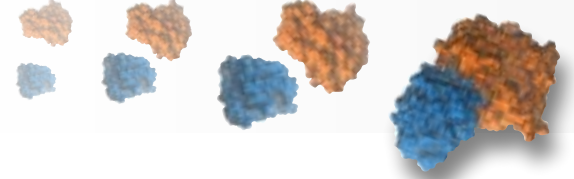
$$\exp[-\beta w(z)] = \frac{\int dx \delta[\mathbf{z} - \mathbf{z}'(\mathbf{x})] \exp[-\beta U(\mathbf{x})]}{\int dx \exp[-\beta U(\mathbf{x})]}$$

$$\mathbf{F} = -\nabla_z w(z)$$

Evolve the string of images until:

$$(\mathbf{DF})^\perp = 0$$





THE STRING METHOD WITH SWARMS OF TRAJECTORIES



Let us assume that the collective variables evolve on the free-energy landscape according to non-inertial Brownian dynamics,

$$z_i(\delta t) = z_i(0) + \sum_j \{ \beta D_{ij}[\mathbf{z}(0)] F_j[\mathbf{z}(0)] + \nabla_{z_j} D_{ij}[\mathbf{z}(0)] \} \delta t + R_i(0) \quad \text{where} \quad F_i = -\nabla_i w(\mathbf{z})$$

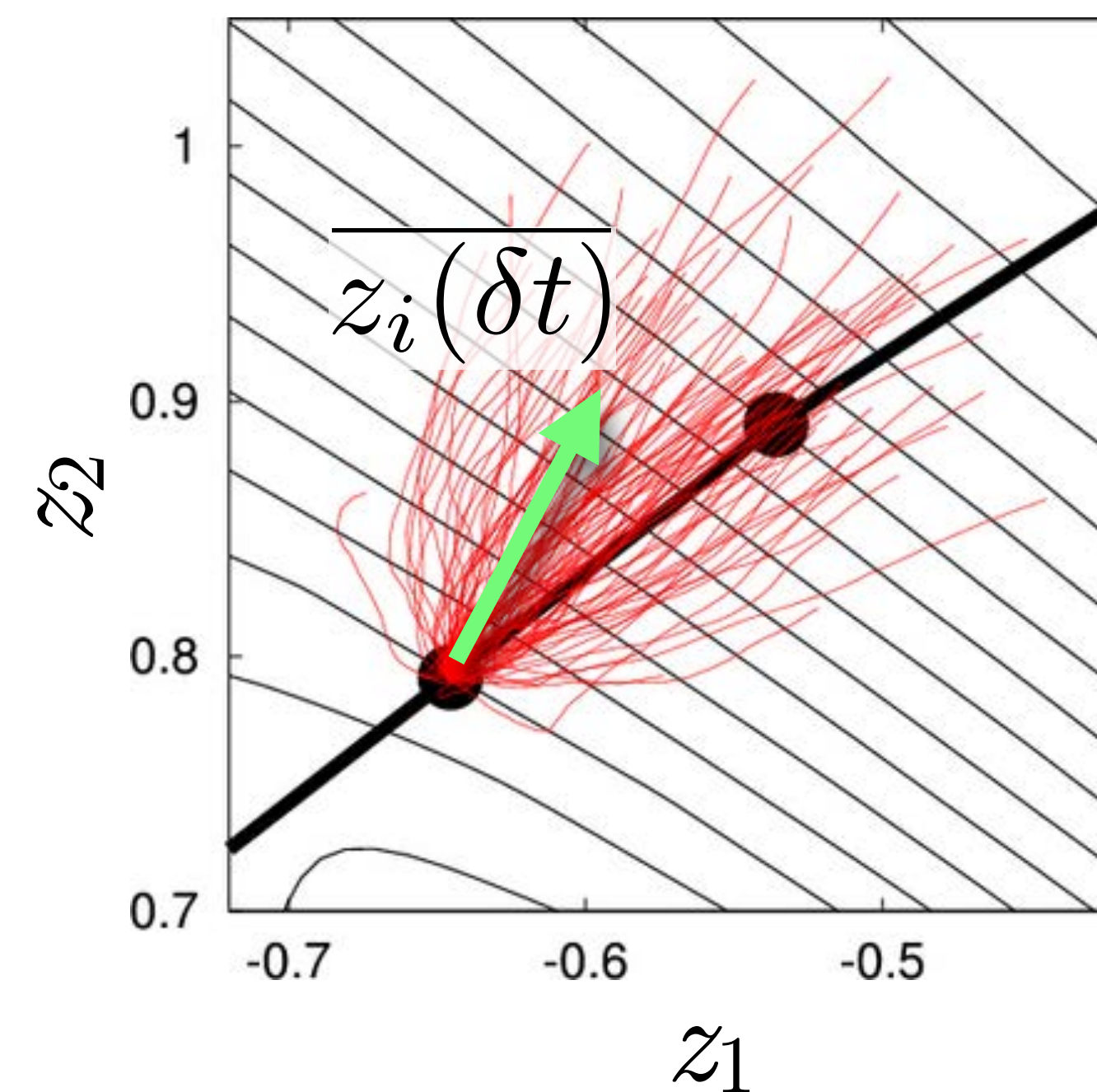
$$\langle R_i(0) R_i(\delta t) \rangle = 2D_{ij} \delta t$$

Let us consider a path $\mathbf{z}(\alpha)$ connecting the two basins, such that α varies between 0 and 1,

$$z_i(\alpha) = z_i(\alpha') + \sum_j \{ \beta D_{ij}[\mathbf{z}(0)] F_j[\mathbf{z}(0)] + \nabla_{z_j} D_{ij}[\mathbf{z}(0)] \} \delta t$$

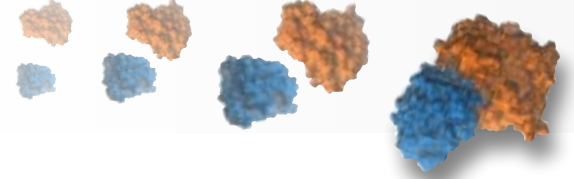
Average drift from an ensemble of unbiased trajectories of length δt ,

$$\overline{z_i(\delta t)} = \overline{z_i(\delta t) - z_i(0)} = \sum_j \{ \beta D_{ij}[\mathbf{z}(0)] F_j[\mathbf{z}(0)] + \nabla_{z_j} D_{ij}[\mathbf{z}(0)] \} \delta t$$

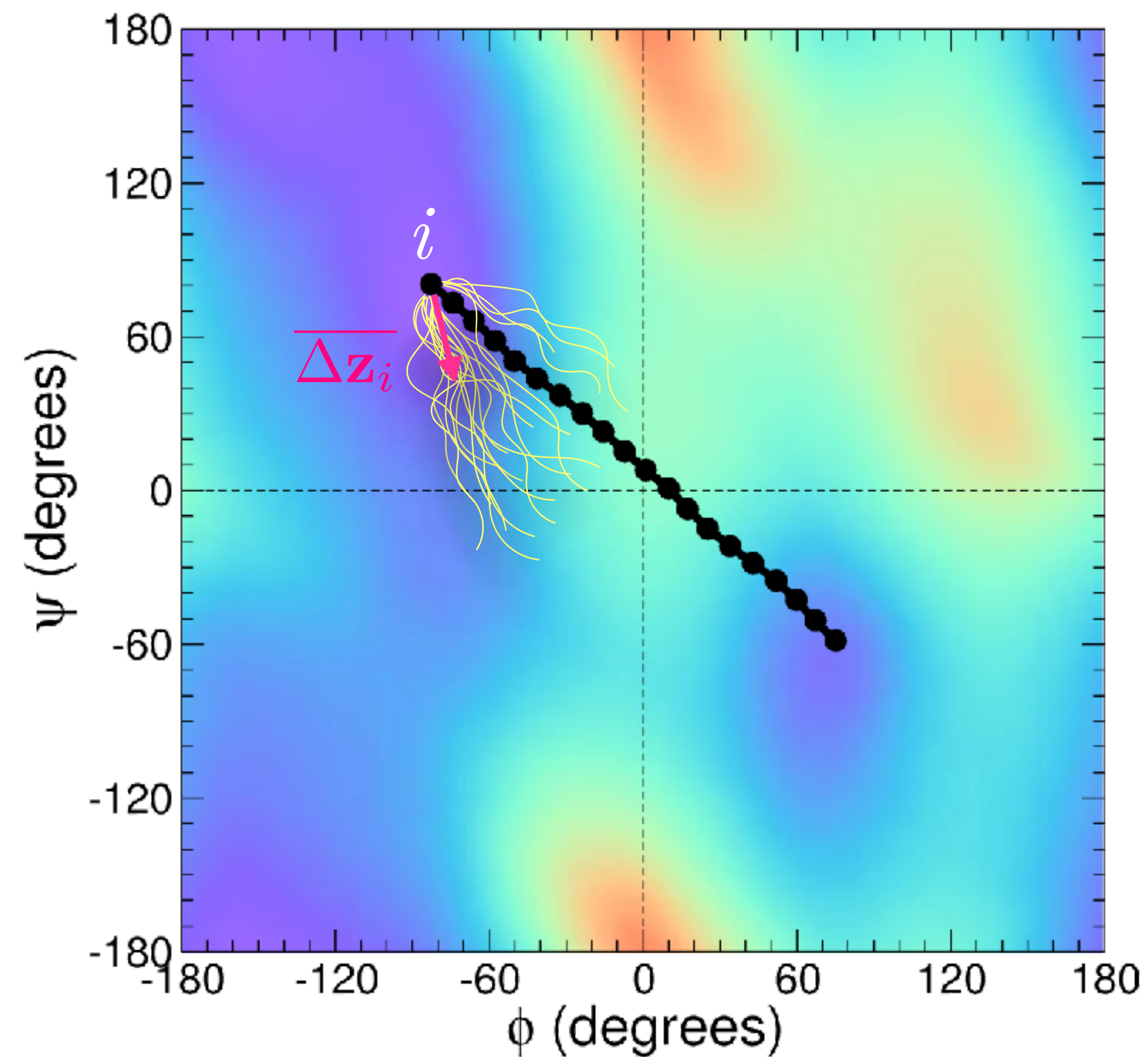


Pan, A. C.; Sezer, D.; Roux, B. *J. Phys. Chem. B* **2008**, *112*, 3432-3440

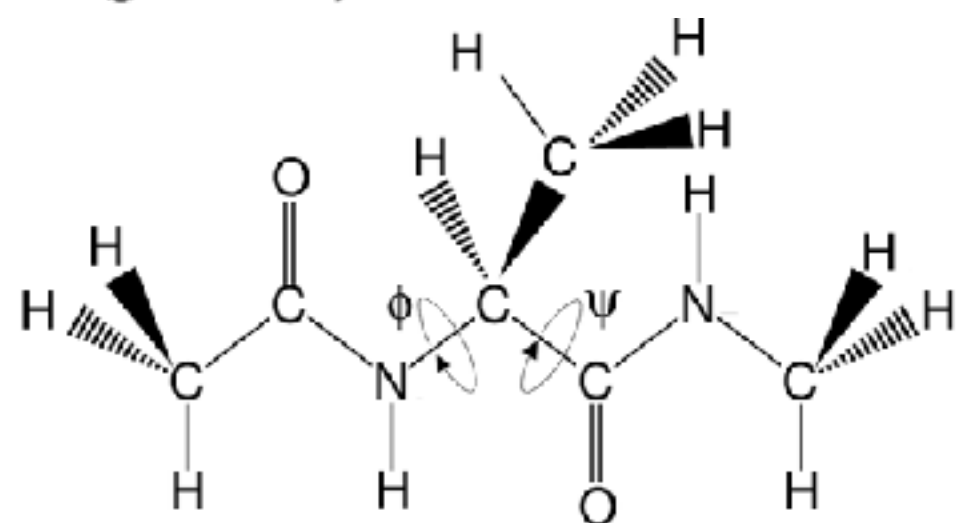
Ren, W.; Vanden-Eijnden, E.; Maragakis, P.; E, W. *J. Chem. Phys.* **2005**, *123*, 134109



THE STRING METHOD WITH SWARMS OF TRAJECTORIES



$$\begin{cases} C_{7eq} : -81^\circ, +81^\circ \\ C_{7ax} : +63^\circ, -81^\circ \end{cases}$$



Prepare a configuration for each one of the P images of the string, the corresponding collective variables of which are close to \mathbf{z}_i , for $i = 1, \dots, P$.

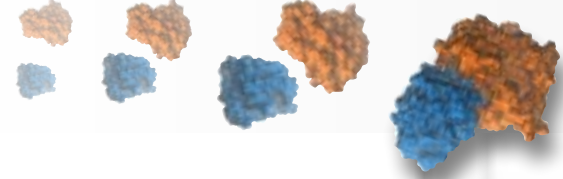
Generate an equilibrium trajectory for each image with \mathbf{z} restrained around \mathbf{z}_i .

From the equilibrium trajectory, generate a large number of short, unbiased trajectories for each image.

Use the resulting average displacement, $\overline{\Delta \mathbf{z}_i}$, to determine the position of the P images.

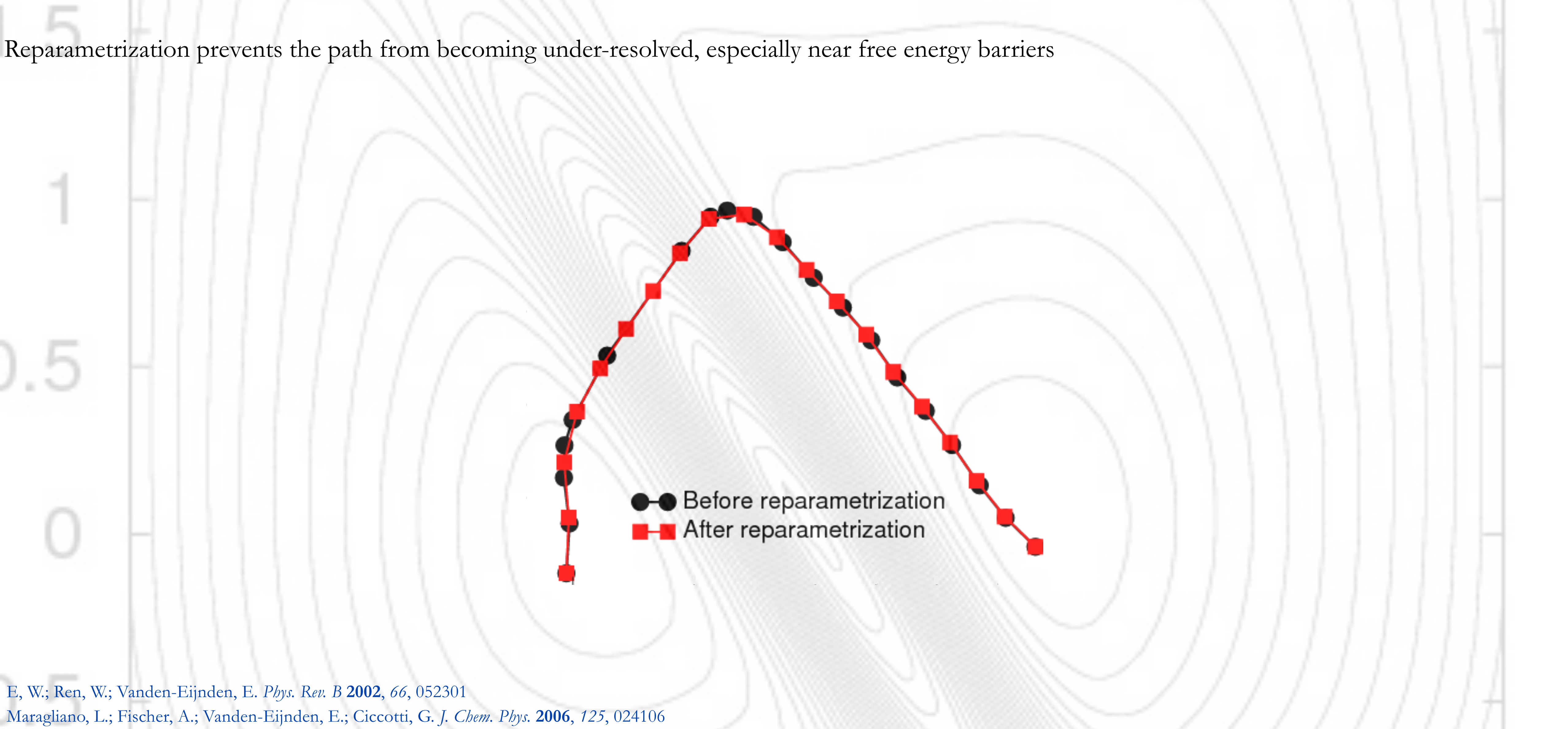
Parameterize the string to ensure that the images are equidistant in collective-variable space.

Pan, A. C.; Sezer, D.; Roux, B. *J. Phys. Chem. B* **2008**, *112*, 3432-3440



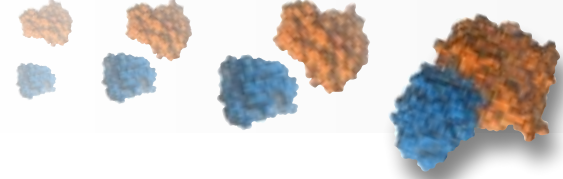
THE STRING METHOD WITH SWARMS OF TRAJECTORIES

Reparametrization prevents the path from becoming under-resolved, especially near free energy barriers

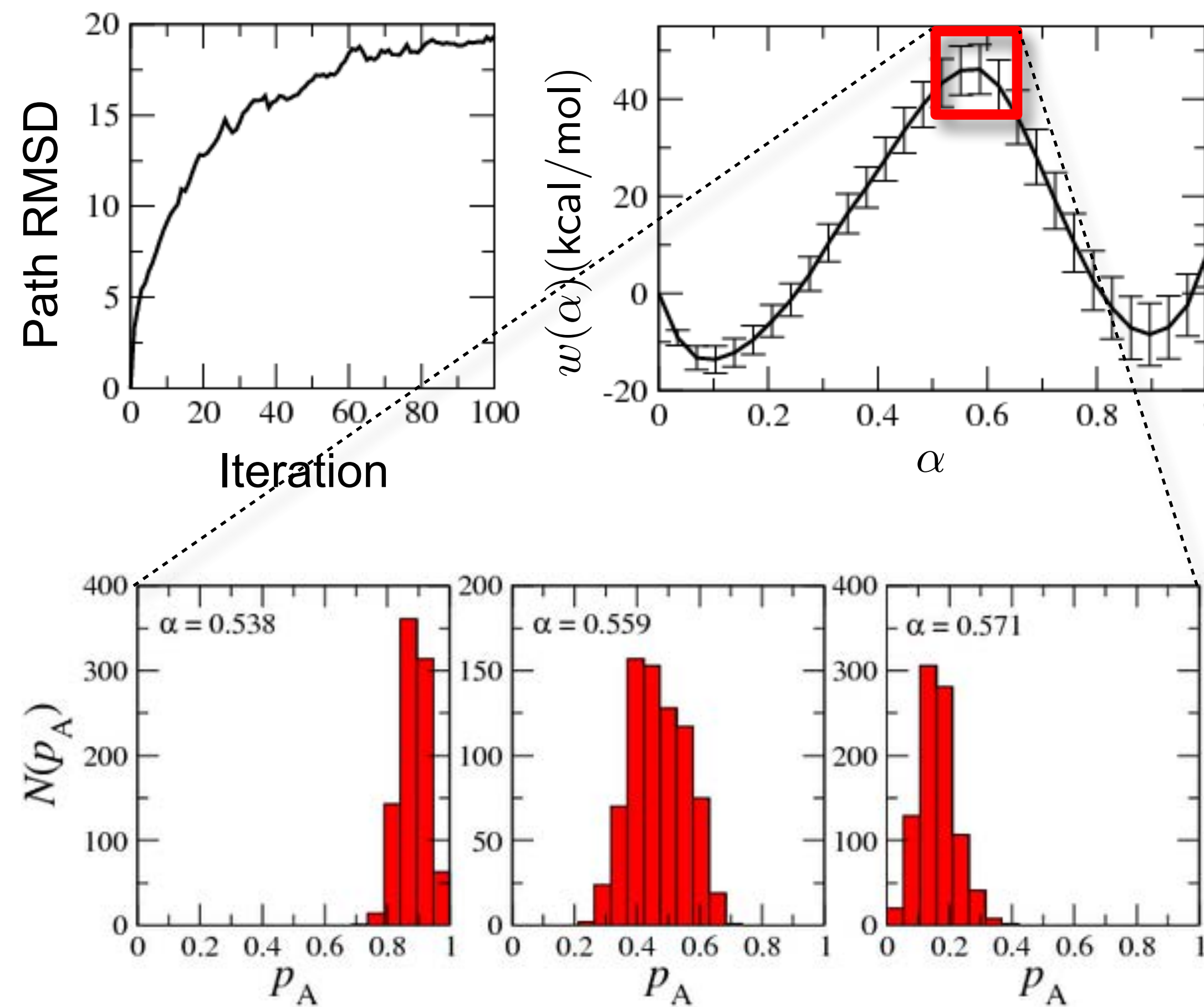
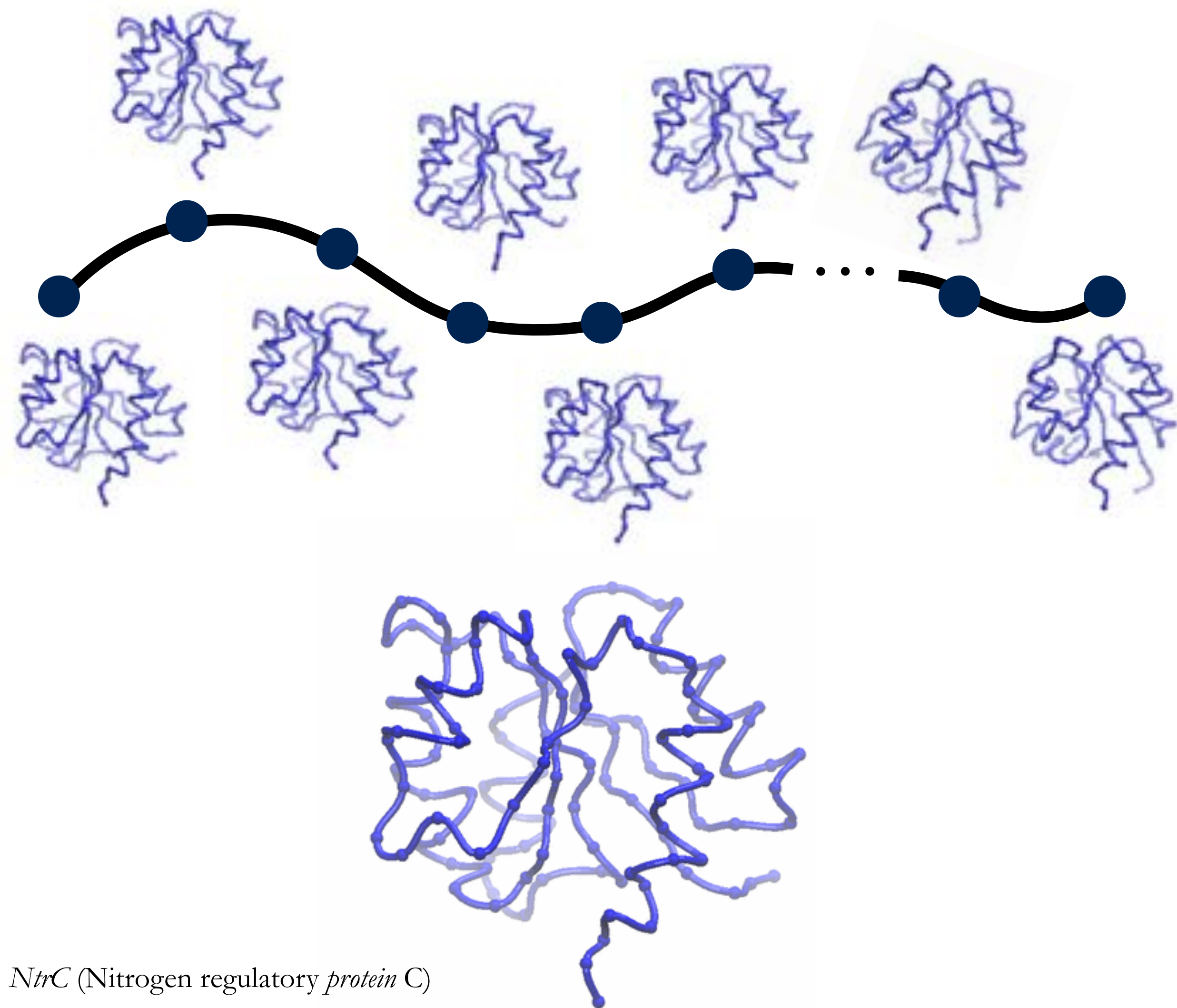


E, W.; Ren, W.; Vanden-Eijnden, E. *Phys. Rev. B* **2002**, *66*, 052301

Maragliano, L.; Fischer, A.; Vanden-Eijnden, E.; Ciccotti, G. *J. Chem. Phys.* **2006**, *125*, 024106

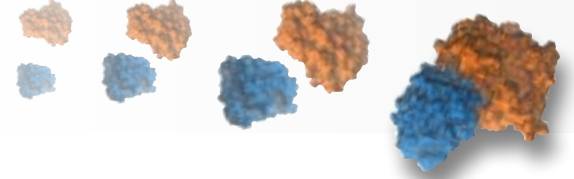


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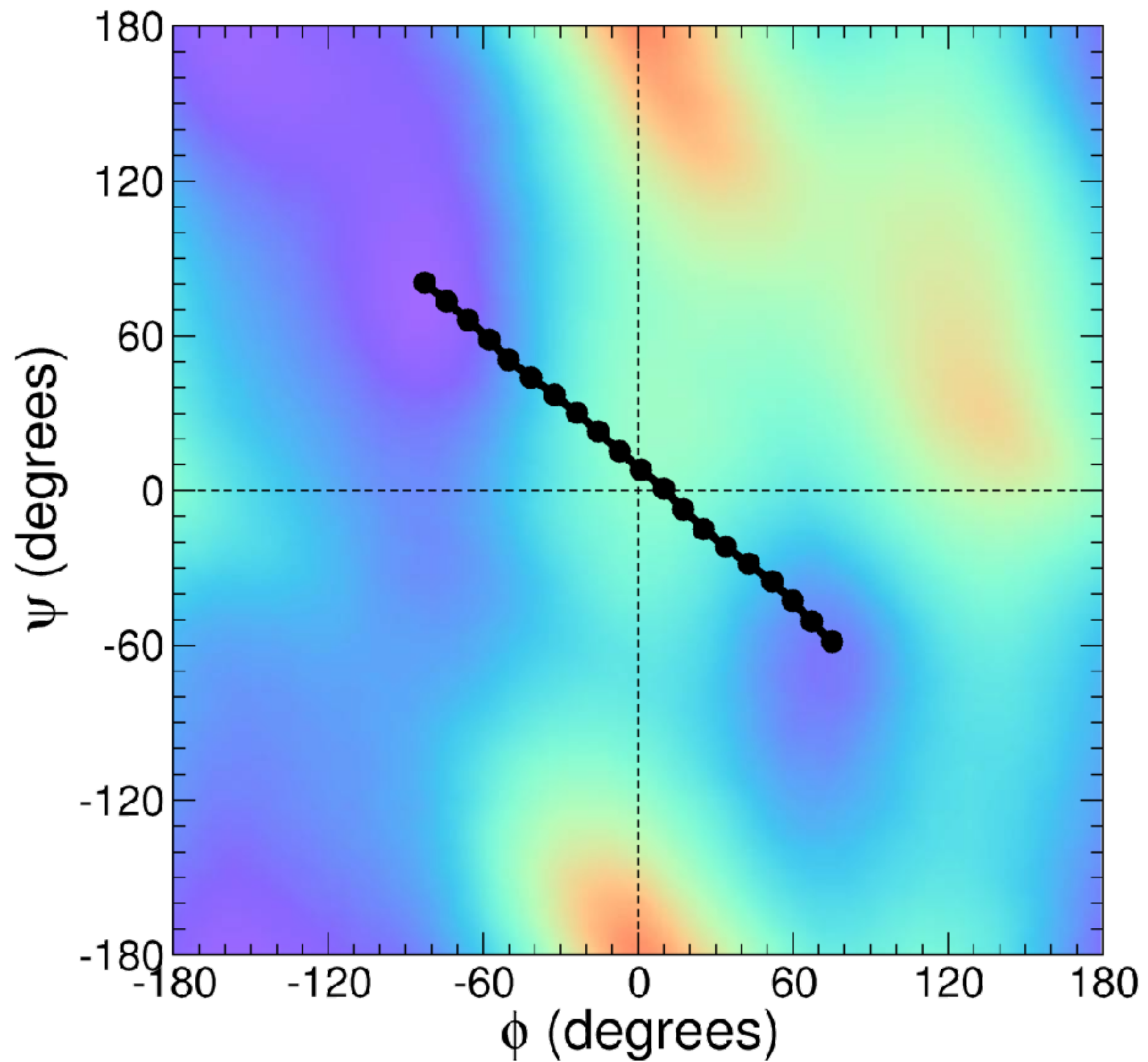


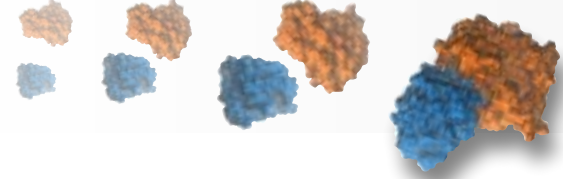
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THE STRING METHOD WITH SWARMS OF TRAJECTORIES





DERIVING THE ONE-DIMENSIONAL FREE-ENERGY PROFILE



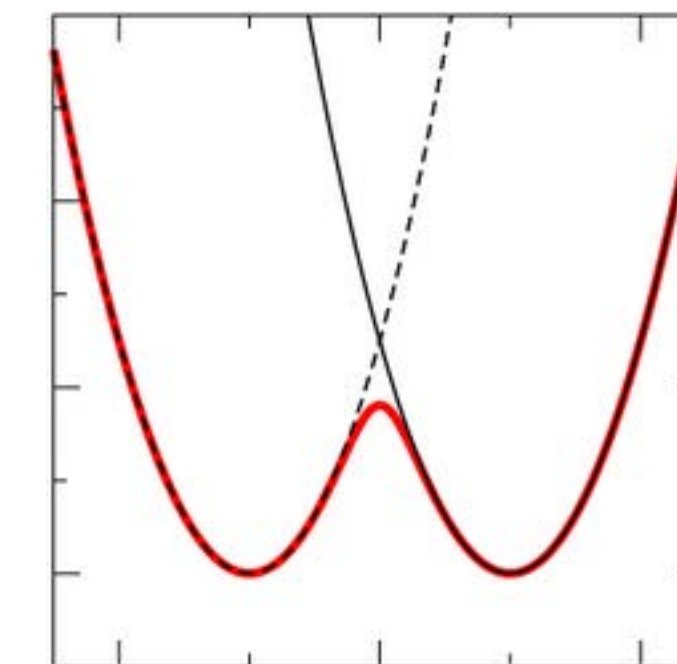
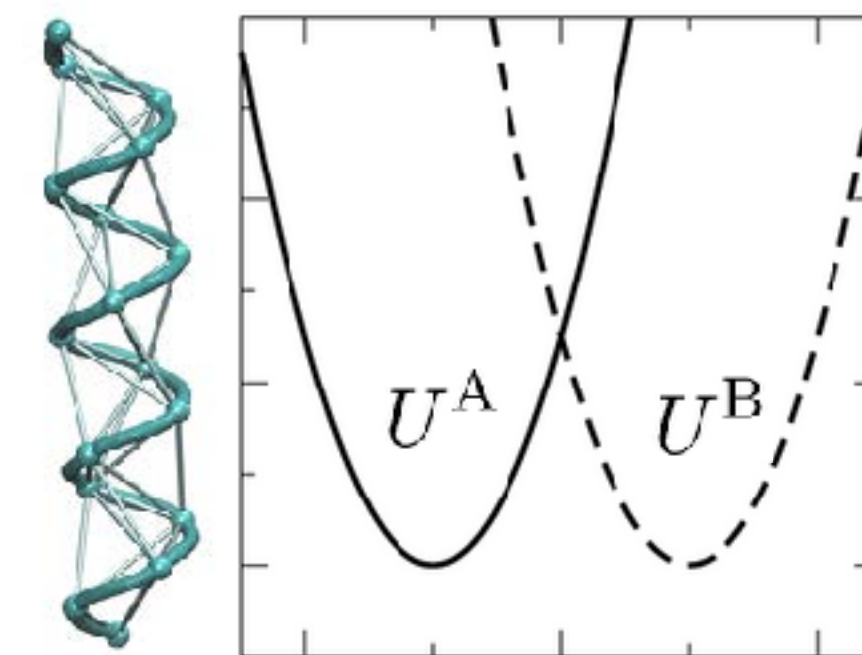
Definition of the end states, A and B, of the transformation

Two-state elastic network model of the transition

Targeted molecular dynamics of the transition

Elaborate transition-path sampling scheme

Free-energy change along minimum-action path



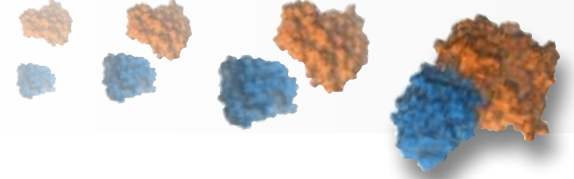
$$\exp(-\beta_m U) = \exp(-\beta_m U^A) + \exp(-\beta_m U^B)$$

Tirion, M. M. *Phys. Rev. Lett.* **1996**, 77, 1905

Maragakis, P.; Karplus, M. *J. Mol. Biol.* **2005**, 352, 807-822

Schlitter, J.; Engels, M.; Krüger, P. *J. Mol. Graph.* **1994**, 12, 84-89

Singharoy, A.; Chipot, C. *J. Phys. Chem. B* **2017**, 121, 3502–3514



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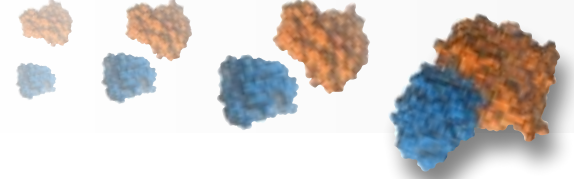
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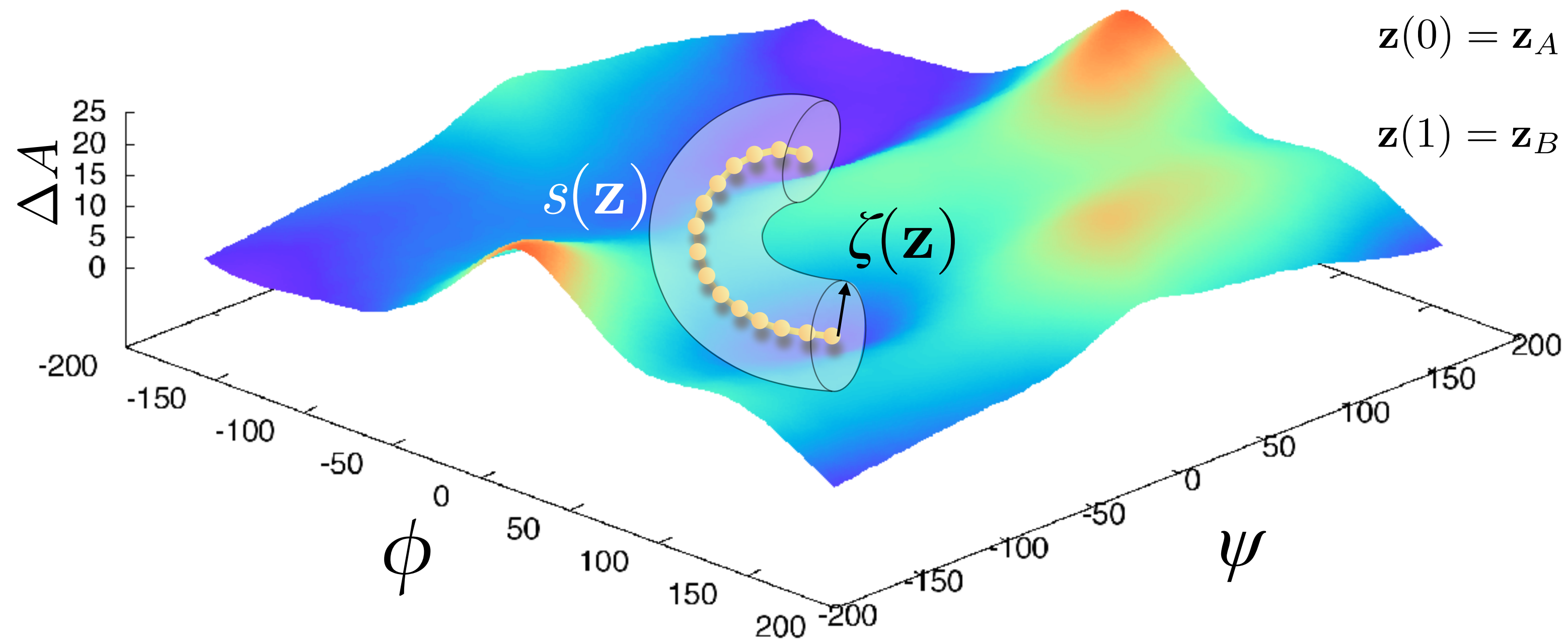
DERIVING THE ONE-DIMENSIONAL FREE-ENERGY PROFILE



Path collective variables:

$$s(\mathbf{z}) = \lim_{\lambda \rightarrow \infty} \frac{\int_0^1 dt t \exp\{-\lambda[\mathbf{z} - \mathbf{z}(t)]^2\}}{\int_0^1 dt \exp\{-\lambda[\mathbf{z} - \mathbf{z}(t)]^2\}}$$

$$\zeta(\mathbf{z}) = \lim_{\lambda \rightarrow \infty} -\frac{1}{\lambda} \int_0^1 dt \exp\{-\lambda[\mathbf{z} - \mathbf{z}(t)]^2\}$$



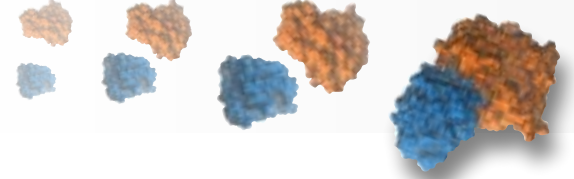
Free-energy surface:

$$F(s, \zeta) = -\frac{1}{\beta} \ln \langle \delta[s - s(\mathbf{z})] \delta[\zeta - \zeta(\mathbf{z})] \rangle$$

Closest minimum free-energy path:

$$\mathcal{F}(\mathbf{z}) = \int_0^1 ds F(s, \zeta)$$

$$\text{Variational principle: } \frac{\delta \mathcal{F}(\mathbf{z})}{\delta \mathbf{z}} = 0$$



DERIVING THE ONE-DIMENSIONAL FREE-ENERGY PROFILE

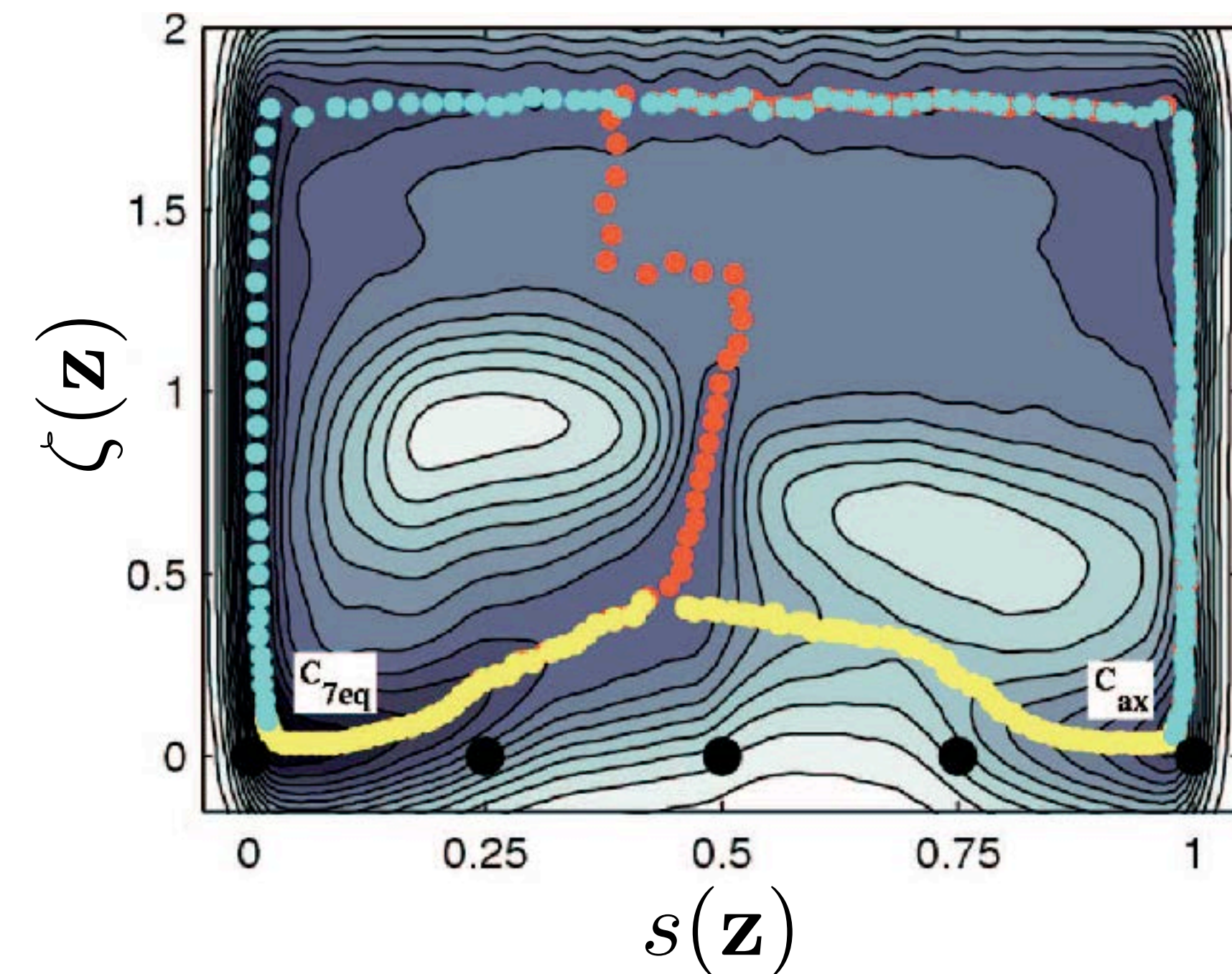
Discrete case: The putative reaction coordinate, $\mathbf{z}(t)$, is discretized in a collection of images.

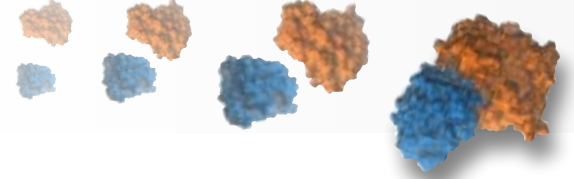
$$s(\mathbf{z}) = \frac{1}{N-1} \frac{\sum_i (i-1) \exp \{ -\lambda [\mathbf{z} - \mathbf{z}(i)]^2 \}}{\sum_i \exp \{ -\lambda [\mathbf{z} - \mathbf{z}(i)]^2 \}}$$

$$\zeta(\mathbf{z}) = -\frac{1}{\lambda} \ln \left(\sum_i \exp \{ -\lambda [\mathbf{z} - \mathbf{z}(i)]^2 \} \right)$$

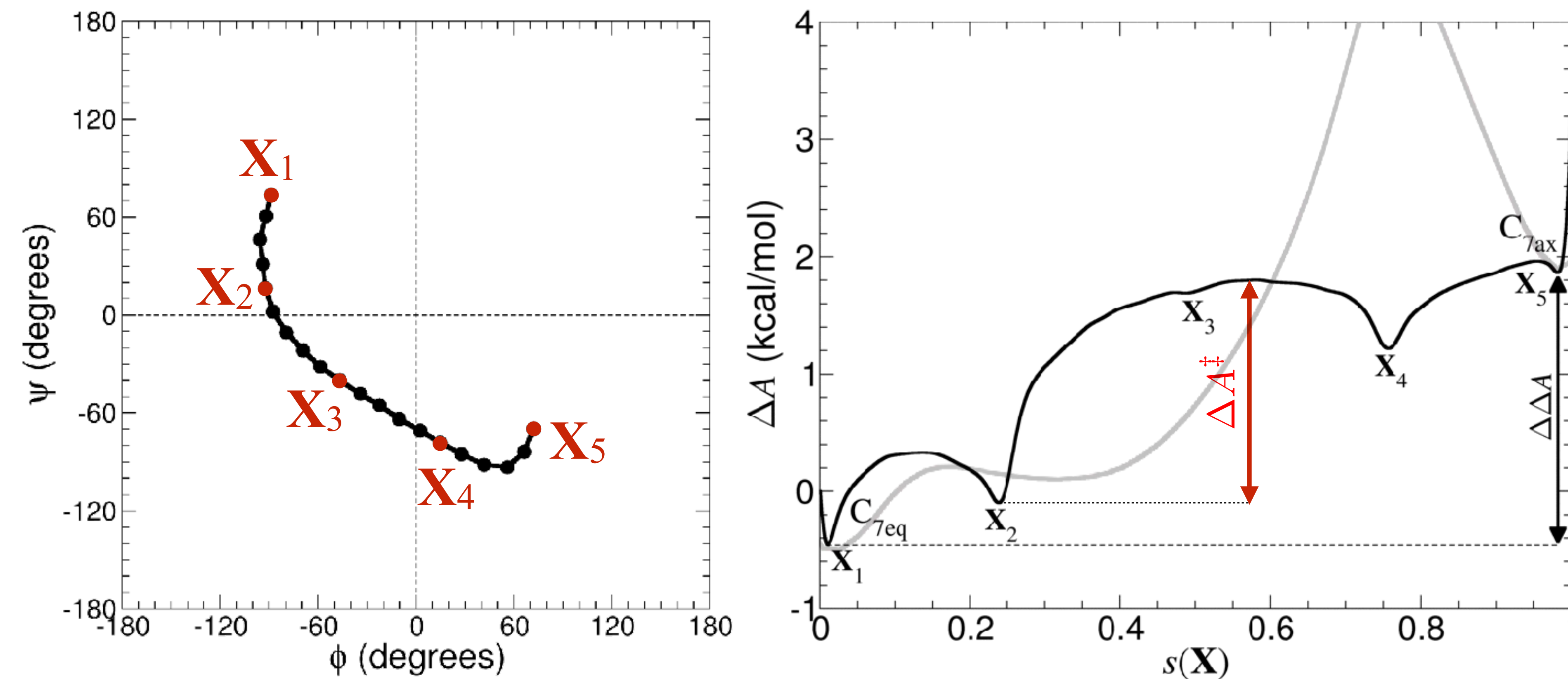
Closest minimum free-energy path:

$$\mathcal{F}(\mathbf{z}) = \frac{1}{N} \left\{ \frac{1}{2} \left[F(s(\mathbf{z}(1)), 0) + F(s(\mathbf{z}(N)), 0) \right] + \sum_{i=2}^{N-1} F(s(\mathbf{z}(i)), 0) \right\}$$



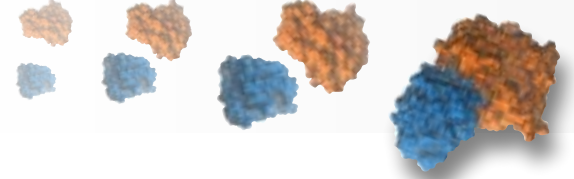


DERIVING THE ONE-DIMENSIONAL FREE-ENERGY PROFILE



From the integration of C_{7eq} and C_{7ax} basins, $\Delta\Delta A = 2.5$ kcal/mol.

From the difference of RMSD's, $\Delta A^\ddagger = 5.6$ kcal/mol.



DERIVING THE ONE-DIMENSIONAL FREE-ENERGY PROFILE



Free energy of the biased system, or perturbed free energy,

$$e^{-\beta F(\boldsymbol{\eta})} = \int d\boldsymbol{\zeta} e^{-\beta(G(\boldsymbol{\zeta}) + U_{\boldsymbol{\eta}}(\boldsymbol{\zeta}))}$$

$$e^{-\beta G(\boldsymbol{\zeta})} \propto e^{-\frac{1}{2\beta k} \nabla_{\boldsymbol{\zeta}}^2} e^{-\beta F(\boldsymbol{\zeta})}$$

For large k , in the stiff-spring approximation, one may expand the PMF to extract the first two terms in $1/k$,

$$G(\boldsymbol{\zeta}) \approx F(\boldsymbol{\zeta}) + \frac{1}{2\beta k} \left(\beta \nabla_{\boldsymbol{\zeta}} F(\boldsymbol{\zeta}) \cdot \nabla_{\boldsymbol{\zeta}} F(\boldsymbol{\zeta}) - \nabla_{\boldsymbol{\zeta}}^2 F(\boldsymbol{\zeta}) \right)$$

Assuming that $\zeta(s)$ approximately represents the minimum free-energy path, and s is its arc-length,

$$G(\zeta(s)) \approx F(s) + \frac{1}{2\beta k} \left\{ \beta \left(\frac{d}{ds} F(s) \right)^2 - \frac{d^2}{ds^2} F(s) \right\}$$

The perturbed free energies, $F_i = F(\zeta(s_i))$, can be estimated by solving self-consistently,

$$e^{-\beta F_i} = \sum_t \frac{e^{-\beta U_i(\boldsymbol{\zeta}^t)}}{\sum_j T_j e^{-\beta(U_j(\boldsymbol{\zeta}^t) - F_j)}} \quad w^t = \left(\sum_i T_i e^{-\beta(U_i(\boldsymbol{\zeta}^t) - F_i)} \right)^{-1}$$

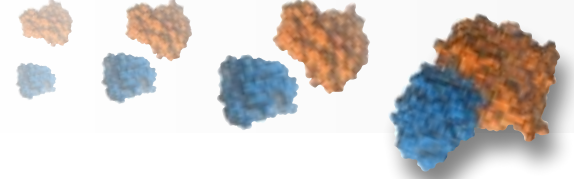
Alternatively,
$$\begin{cases} w^t = \left(\sum_i T_i e^{-\beta(U_i(\boldsymbol{\zeta}^t) - F_i)} \right)^{-1} \\ e^{-\beta F_i} = \sum_t w^t e^{-\beta U_i(\boldsymbol{\zeta}^t)} \end{cases}$$

Hummer, G.; Szabo, A. Proc. Natl. Acad. Sci. U. S. A. 2010, 107, 21441–21446

Bartels, C. Chem. Phys. Lett. 2000, 331, 446–454

Shirts, M. R.; Chodera, J. D. J. Chem. Phys. 2008, 129, 124105

Moradi, M.; Tajkhorshid, E. J. Chem. Theory Comput. 2014, 10, 2866–2880.



INTRODUCTION

Why do we need reaction coordinates?

WHAT IS A GOOD REACTION-COORDINATE MODEL ?

- Reaction coordinate versus order parameter
- Committor distributions

THE STRING METHOD

- The basic string method
- The string method with swarms of trajectories

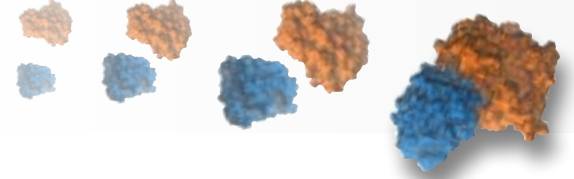
EXTRACTING THE FREE ENERGY FROM THE MINIMUM-ACTION PATH

- Path-collective variables
- Perturbative approach

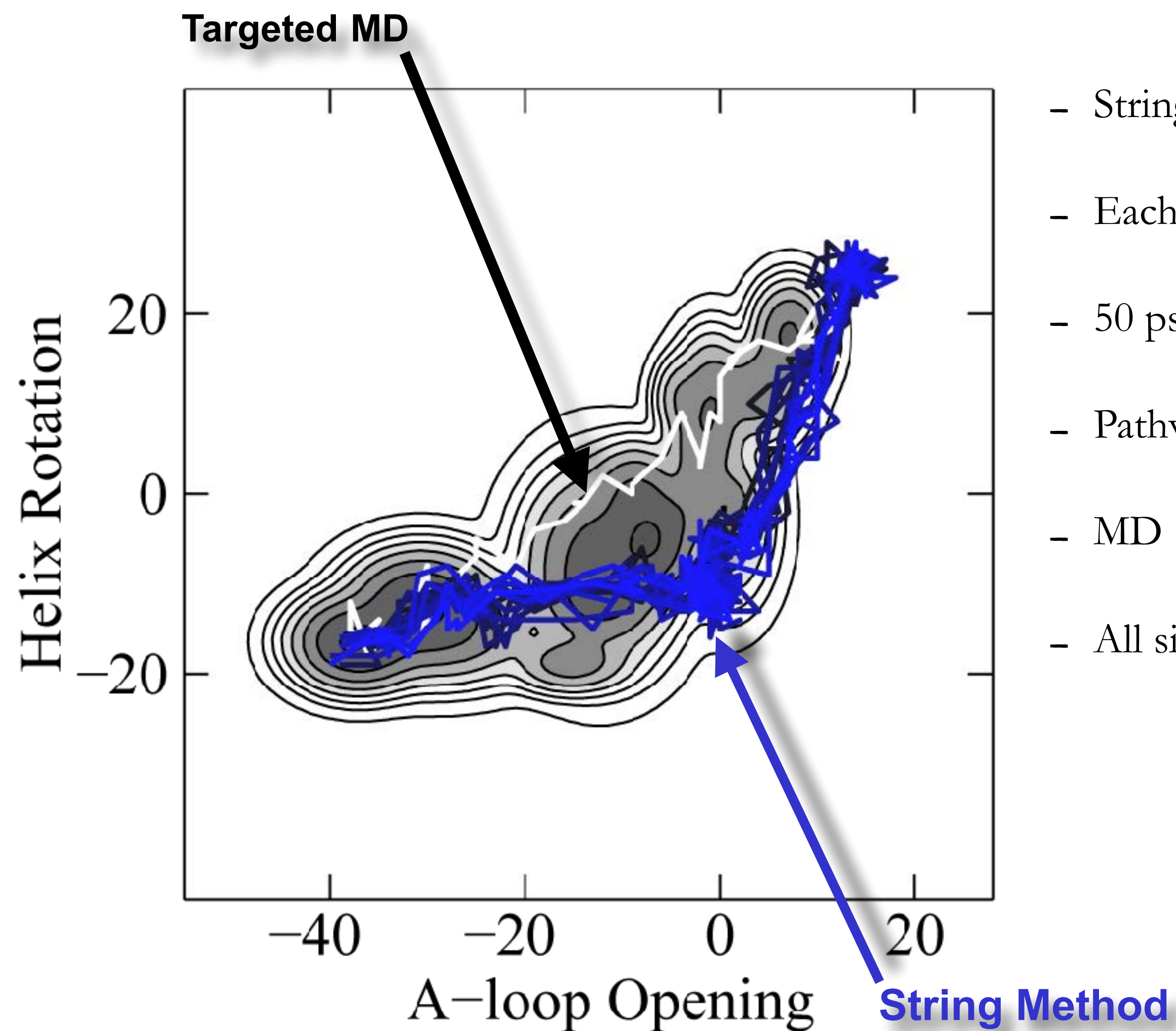
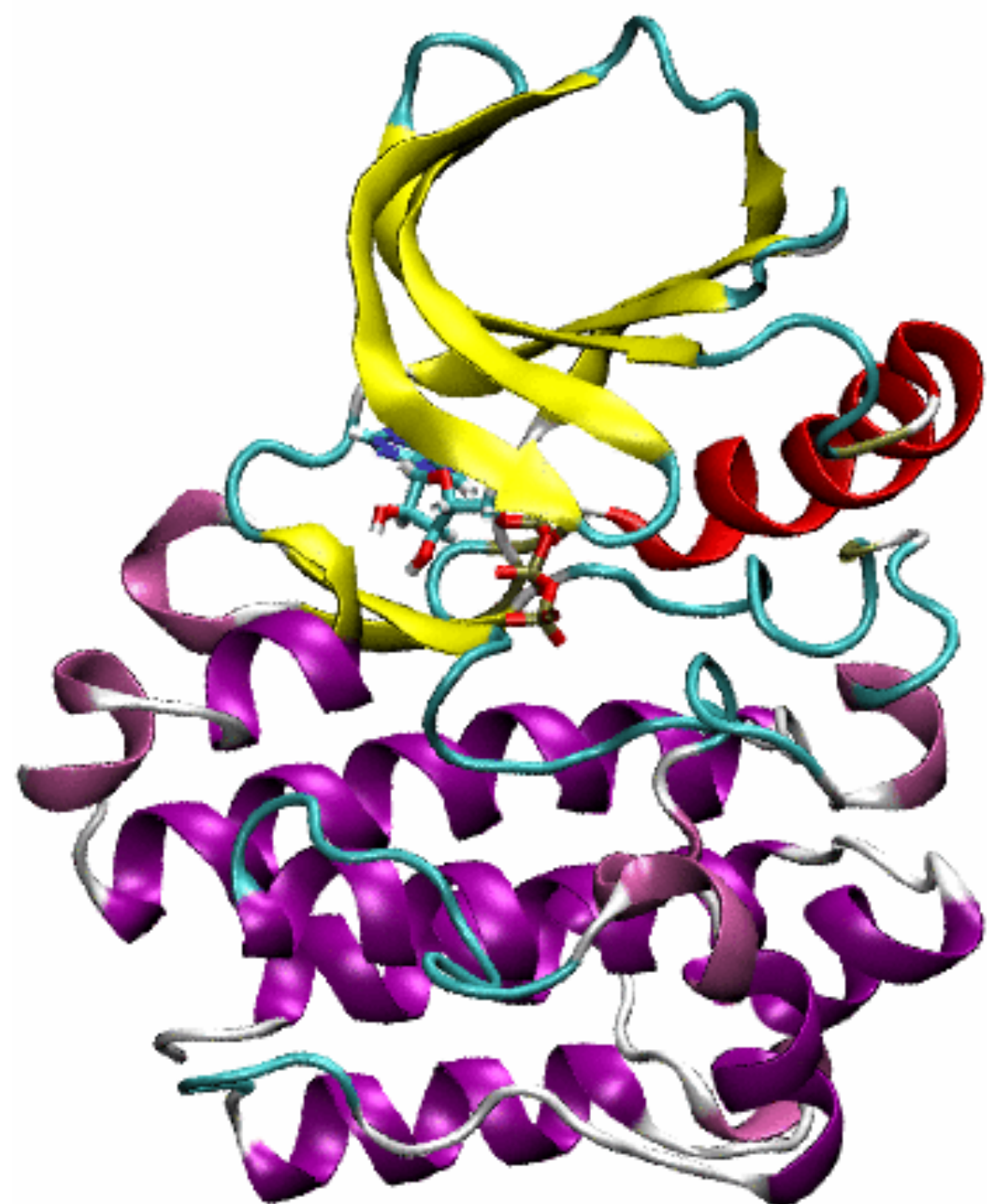
APPLICATIONS

- Transition path of activation loop in c-Src kinase
- Chemomechanical coupling in V_1 -ATPase

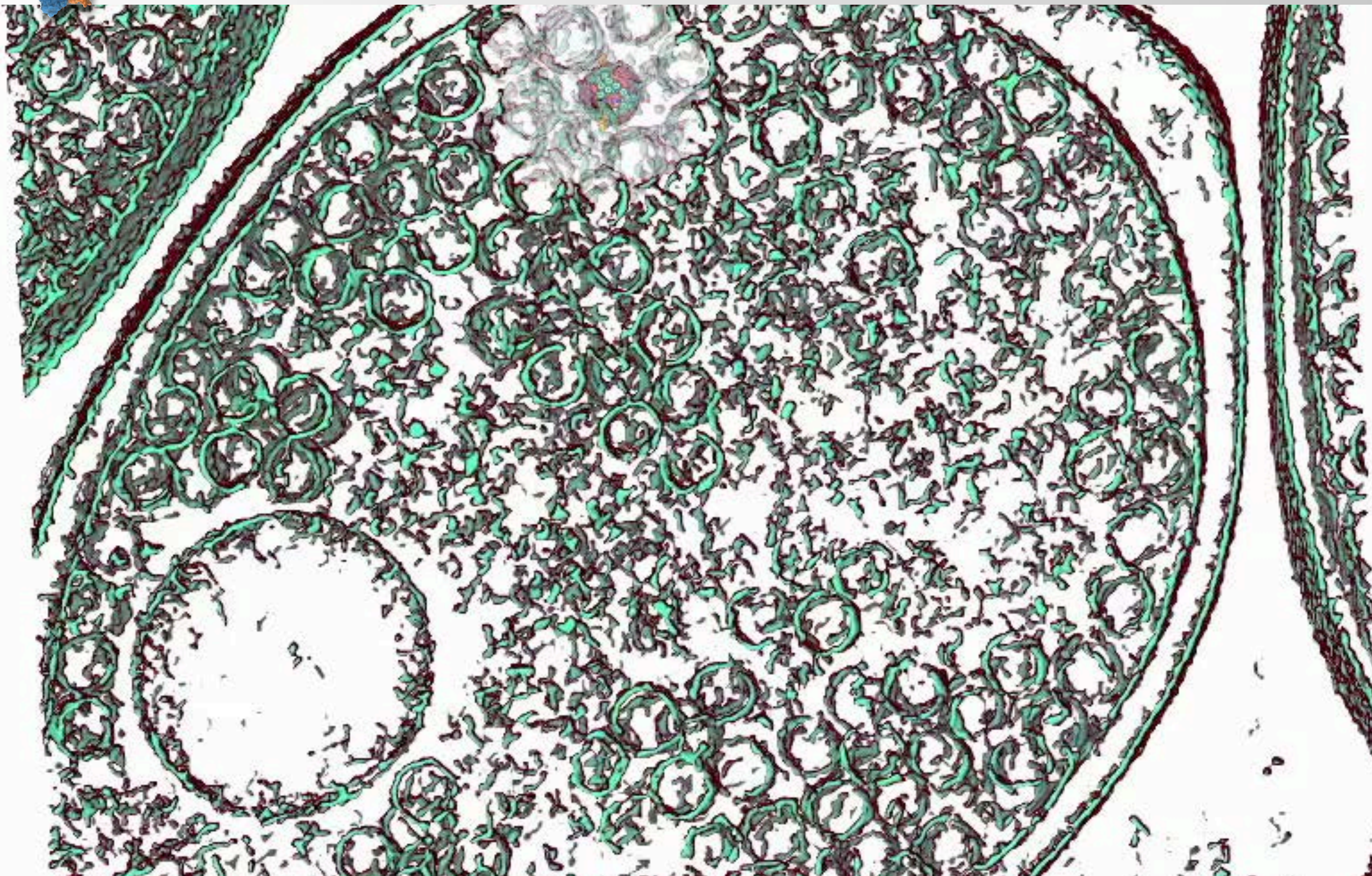
RECONCILING THERMODYNAMICS AND KINETICS

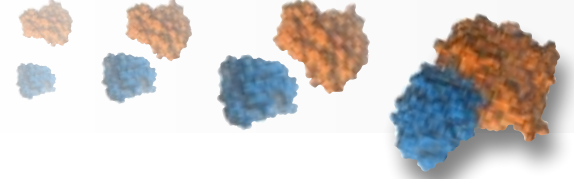


TRANSITION PATH OF ACTIVATION LOOP IN C-SRC KINASE

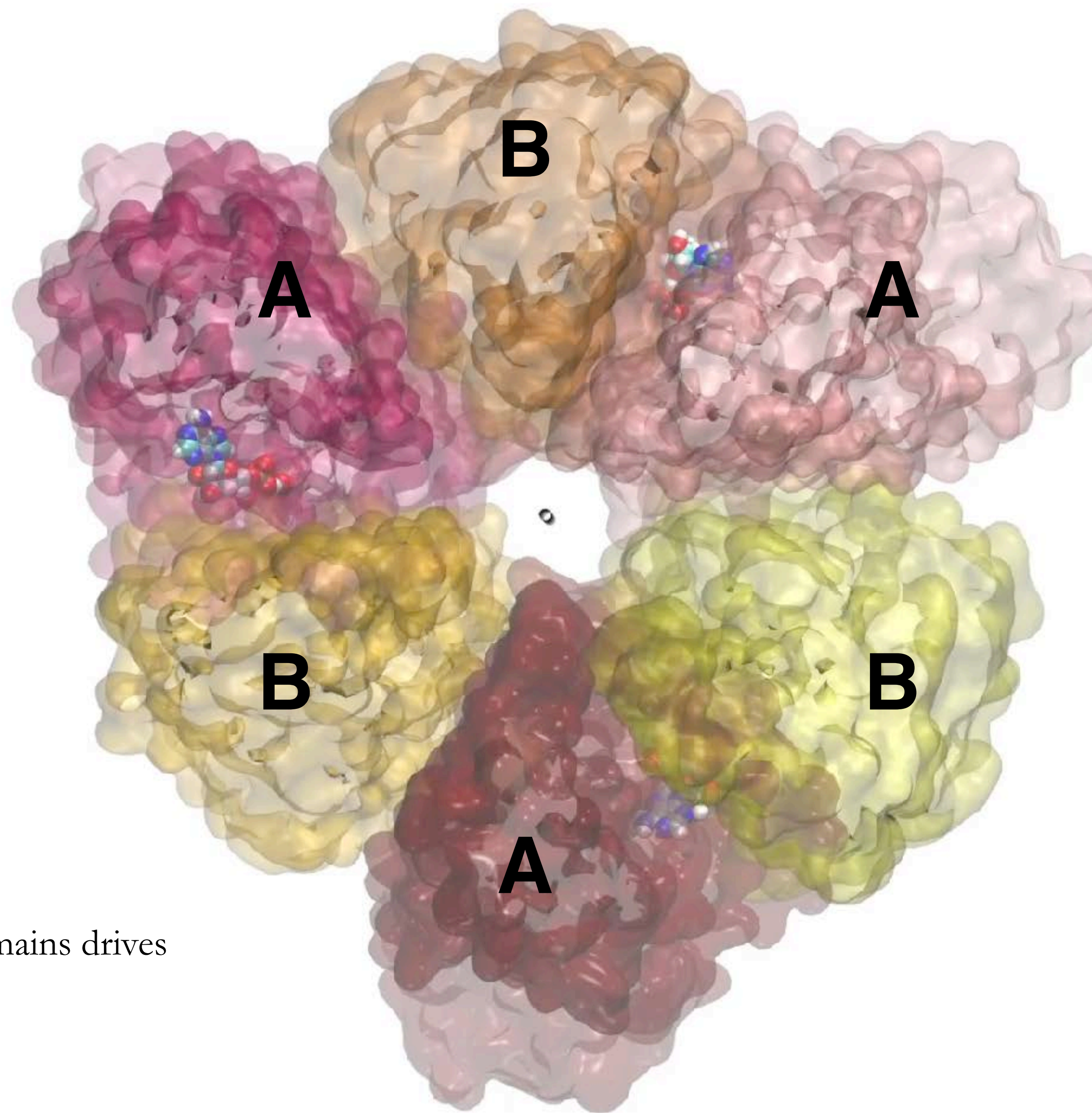
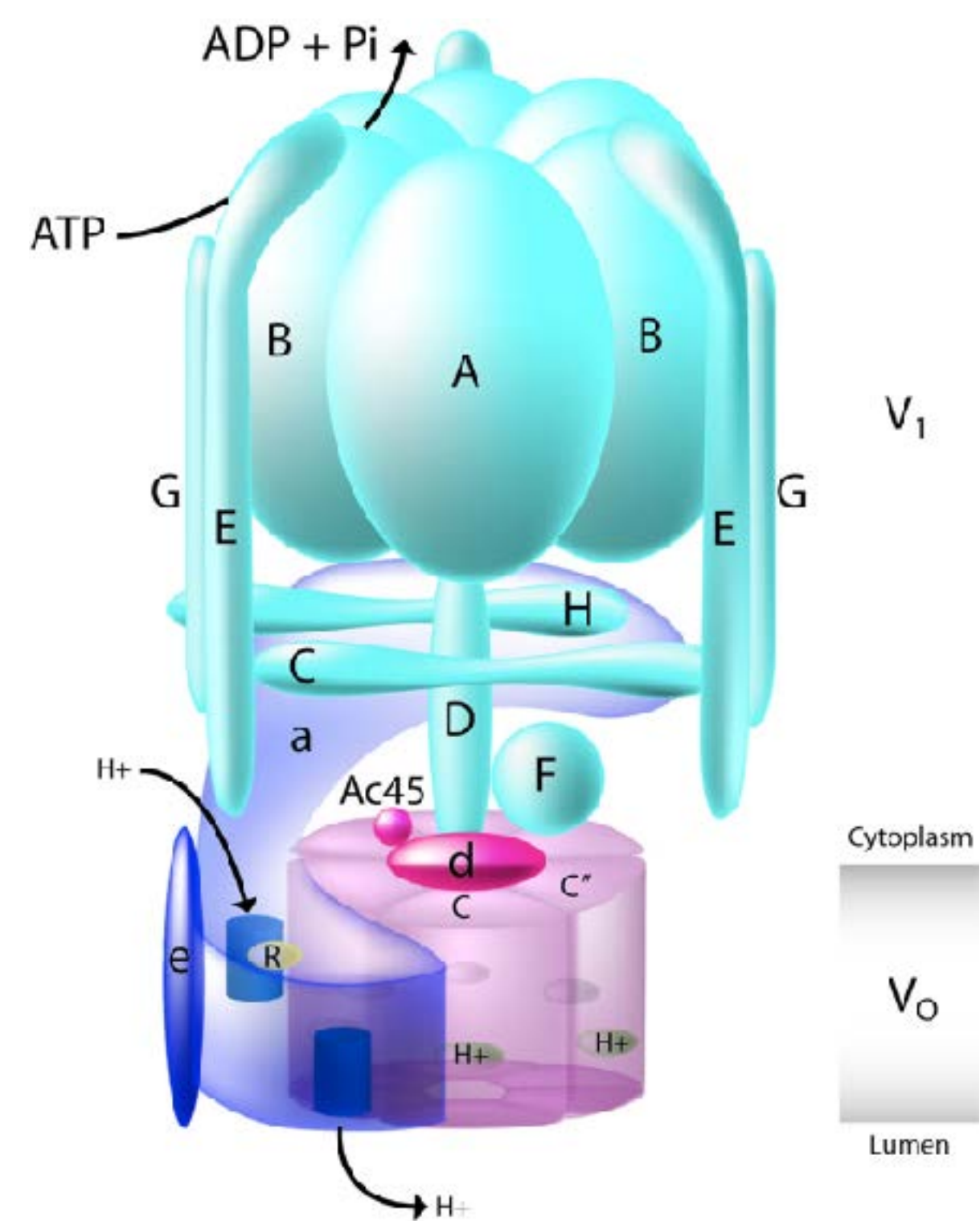


- String has 51 images
- Each swarm contains 100 trajectories, 1 ps each
- 50 ps for constrained equilibration
- Pathways is evolved for 100 iterations
- MD simulations are done with NAMD
- All simulations with explicit solvent



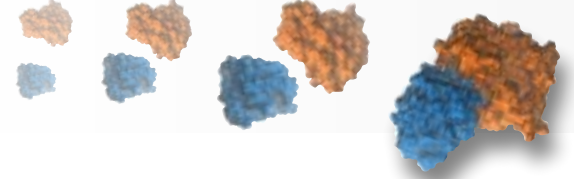


CHEMOMECHANICAL COUPLING IN V1-ATPASE



ATP hydrolysis-driven conformational transitions in the A₃B₃ domains drives rotation of the central stalk.

Singharoy, A.; Chipot, C.; Moradi, M.; Schulten, K. *J. Am. Chem. Soc.* **2017**, *139*, 293-310



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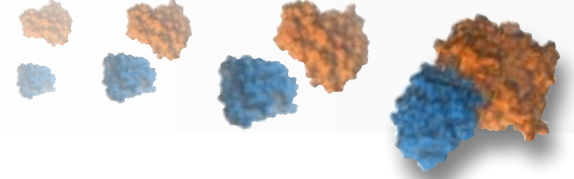
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RECONCILING THERMODYNAMICS AND KINETICS



BEYOND THERMODYNAMICS

$$\Delta Z = \beta D(Z_t) F(Z_t, t) \Delta t + \nabla D(Z_t) \Delta t + \sqrt{2D(Z_t) \Delta t} g_t$$

$$\text{Let: } \begin{cases} \mu = \beta D(Z_t) F(Z_t, t) \Delta t + \nabla D(Z_t) \Delta t \\ \sigma^2 = 2D(Z_t) \Delta t \end{cases}$$

$$\text{Then: } \Delta Z = \mu + \sigma g_t$$

$$P[\Delta Z | w(z), D(z)] = \frac{1}{\sigma \sqrt{2\pi}} \exp\left(-\frac{(\Delta Z - \mu)^2}{2\sigma^2}\right)$$

Probability over the entire trajectory, given the parameters:

$$P[Z(t) | w(z), D(z)] = \prod_i \frac{1}{\sigma_i \sqrt{2\pi}} \exp\left(-\frac{(\Delta Z_i - \mu_i)^2}{2\sigma_i^2}\right)$$

(1) The molecular dynamics supplies the trajectory of the collective variable, $Z(t)$.

(2) The molecular dynamics supplies also, $f_{\text{bias}}(t)$.

(3) Pick trial parameters, $w(z)$ and $D(z)$.

(4) Assume a propagator, e.g., Brownian dynamics.

(5) Calculate the probability of the trajectory given the parameters.

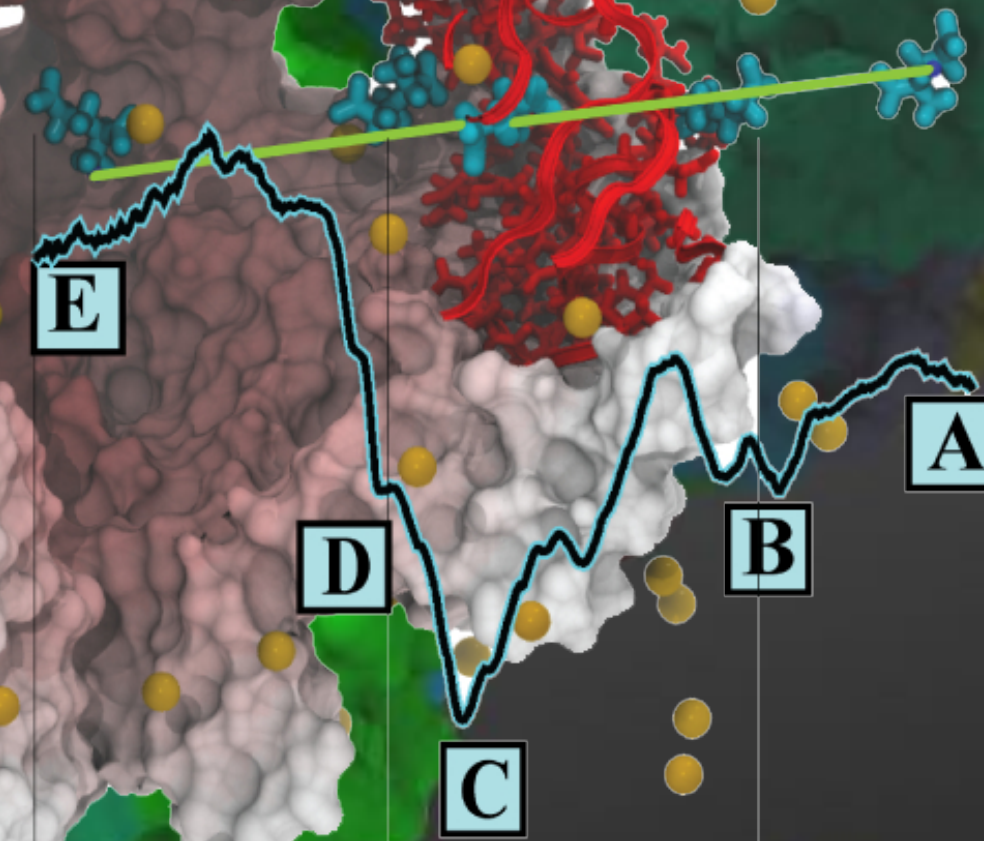
(6) Bayes's theorem: Get the probability of the parameters given the trajectory.

(7) Optimize the parameters to yield the greatest probability.

Kinetic models

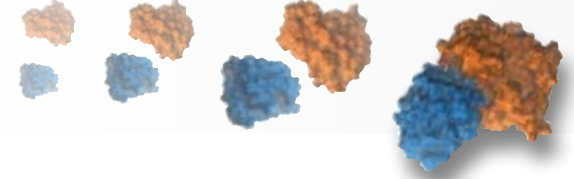
Bacteria resist osmotic stress by means of MscS, but why does the latter include a balloon-like, water filled cytoplasmic domain?

CD acts as an entropic filter that prevents Glu^- from reaching the pores, but once Glu^- enters a pore, it passes unhindered by virtue of a reduced enthalpic barrier preventing clogging.



Martinac, B. et al. *Proc. Natl. Acad. Sci. USA* 1987, 84, 2297-2301

Sotomayor et al. *Biophys. J.* 2006, 90, 3496-3510; *Biophys. J.* 2007, 92, 886-902



Kinetic models

Bacteria resist osmotic stress by means of MscS, but why does the latter include a balloon-like, water filled cytoplasmic domain?

CD acts as an entropic filter that prevents Glu⁻ from reaching the pores, but once Glu⁻ enters a pore, it passes unhindered by virtue of a reduced enthalpic barrier preventing clogging.

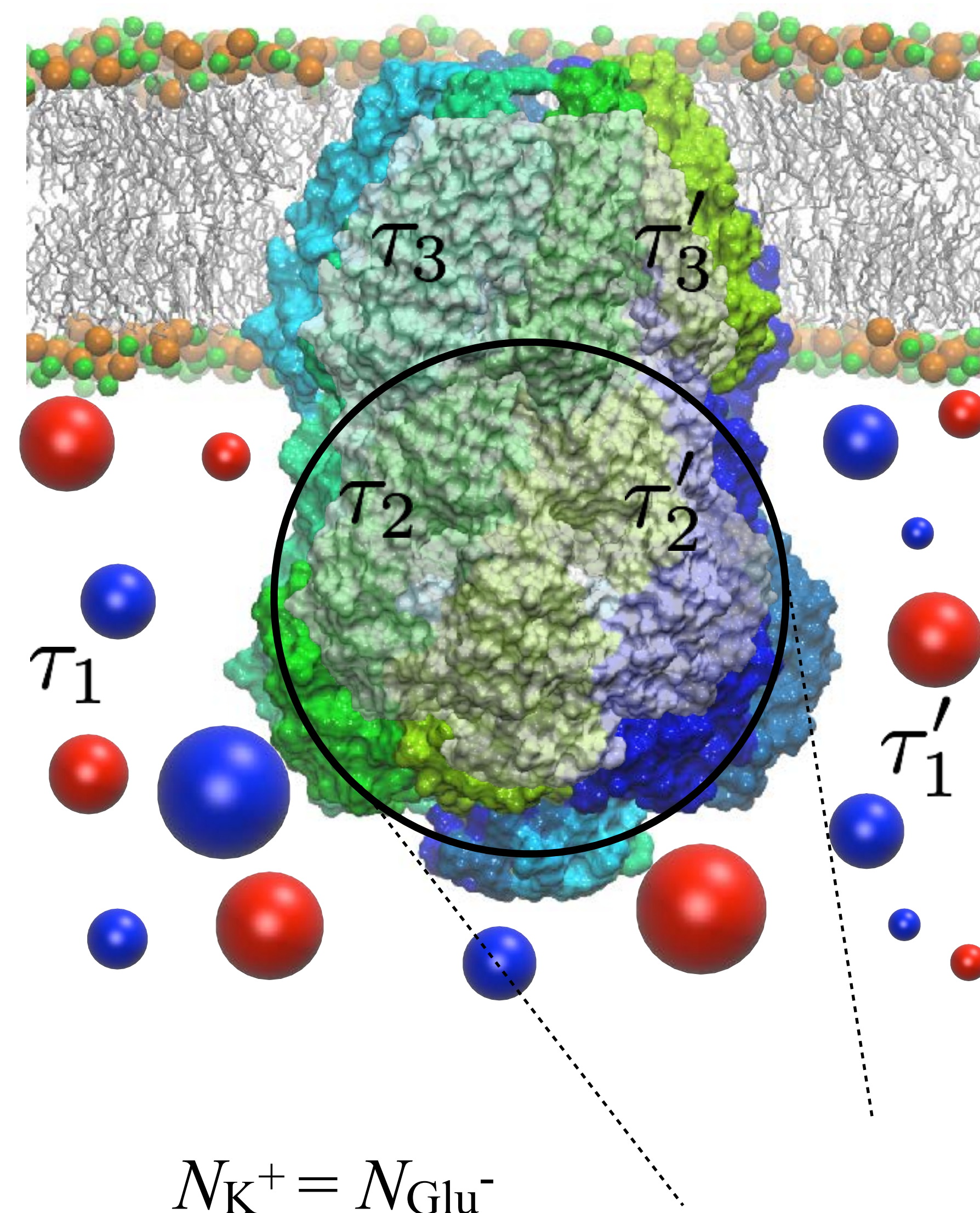
Determine τ_1 , τ_2 and τ_3 for Glu⁻ and K⁺ :

$$\tau = \int_a^b dz \exp[\beta\Delta G(z)] D^{-1}(z) \int_a^z d\zeta \exp[-\beta\Delta G(\zeta)]$$

CD maintains an overall balance of electrolytes to preclude collapse of the transmembrane potential whilst over-coming osmotic shock.

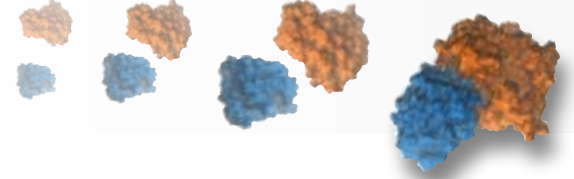
Martinac, B. et al. *Proc. Natl. Acad. Sci. USA* **1987**, *84*, 2297-2301

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Szabo, A.; Schulten, K.; Schulten, Z. *J. Chem. Phys.* **1980**, *72*, 4350-4357

Gamini, R. et al. *Biophys. J.*, **2011**, *101*, 80-89



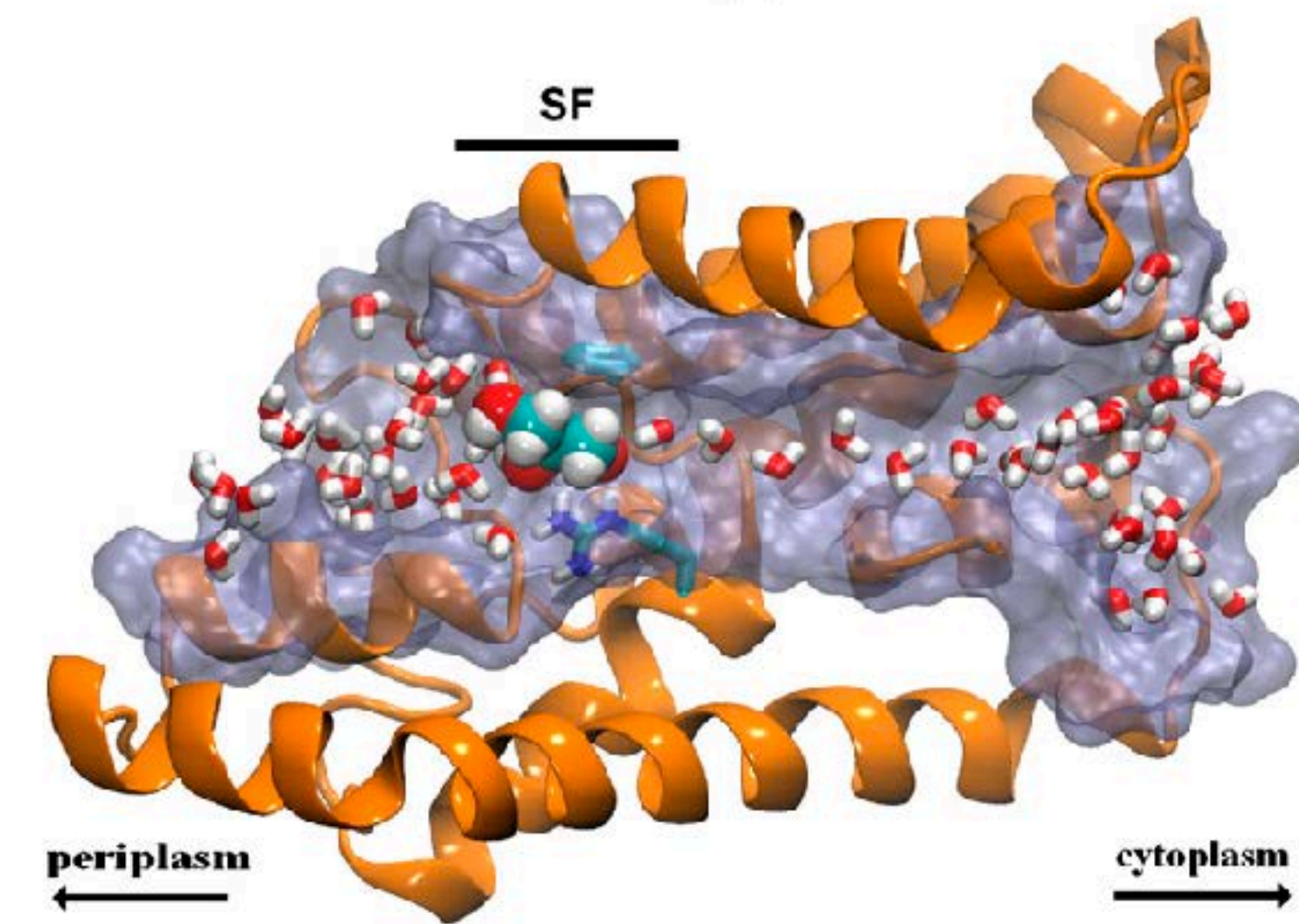
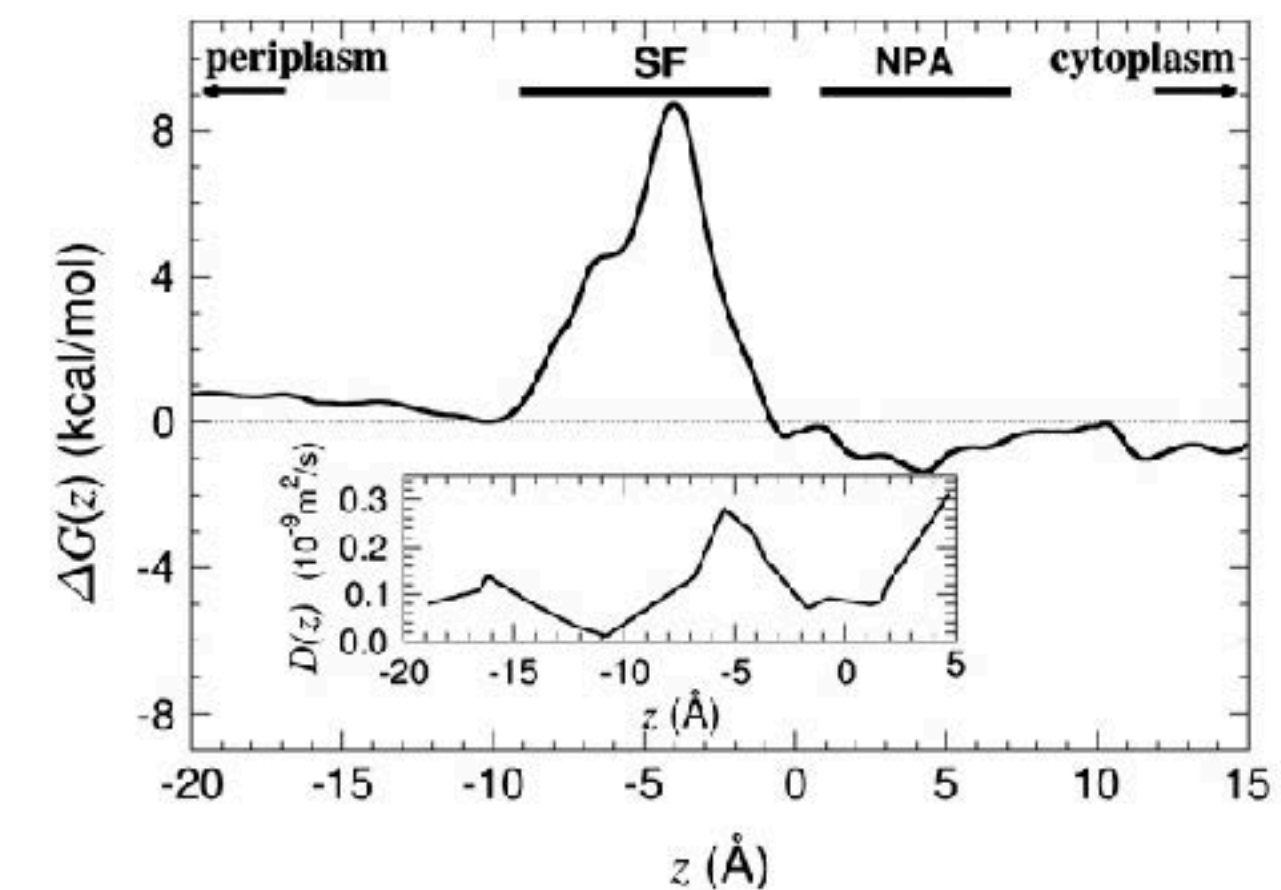
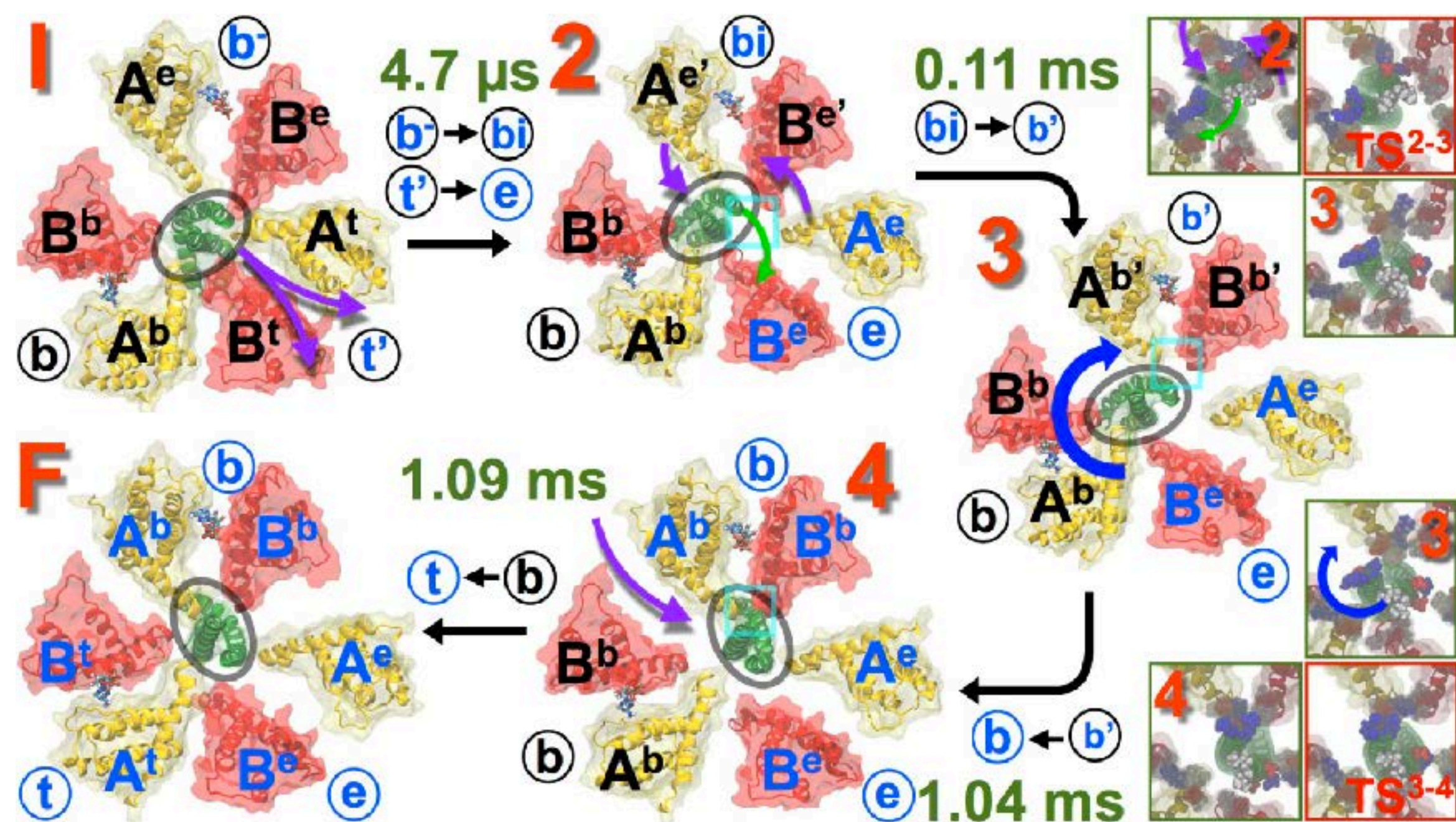
BEYOND THERMODYNAMICS

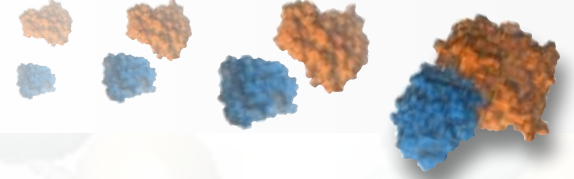
Mean first passage time:

$$\tau = \int_a^b d\xi \exp[\beta\Delta A(\xi)] D^{-1}(\xi) \int_a^\xi d\xi' \exp[-\beta\Delta A(\xi')]$$

Rate constant:

$$k = 1/\tau$$

Szabo, A.; Schulten, K.; Schulten, Z. *J. Chem. Phys.* **1980**, *72*, 4350-4357Singharoy, A.; Chipot, C.; Moradi, M.; Schulten, K. *J. Am. Chem. Soc.* **2017**, *139*, 293-310Hénin, J.; Tajkhorshid, E.; Schulten, K.; Chipot, C. *Biophys. J.* **2008**, *94*, 832-839



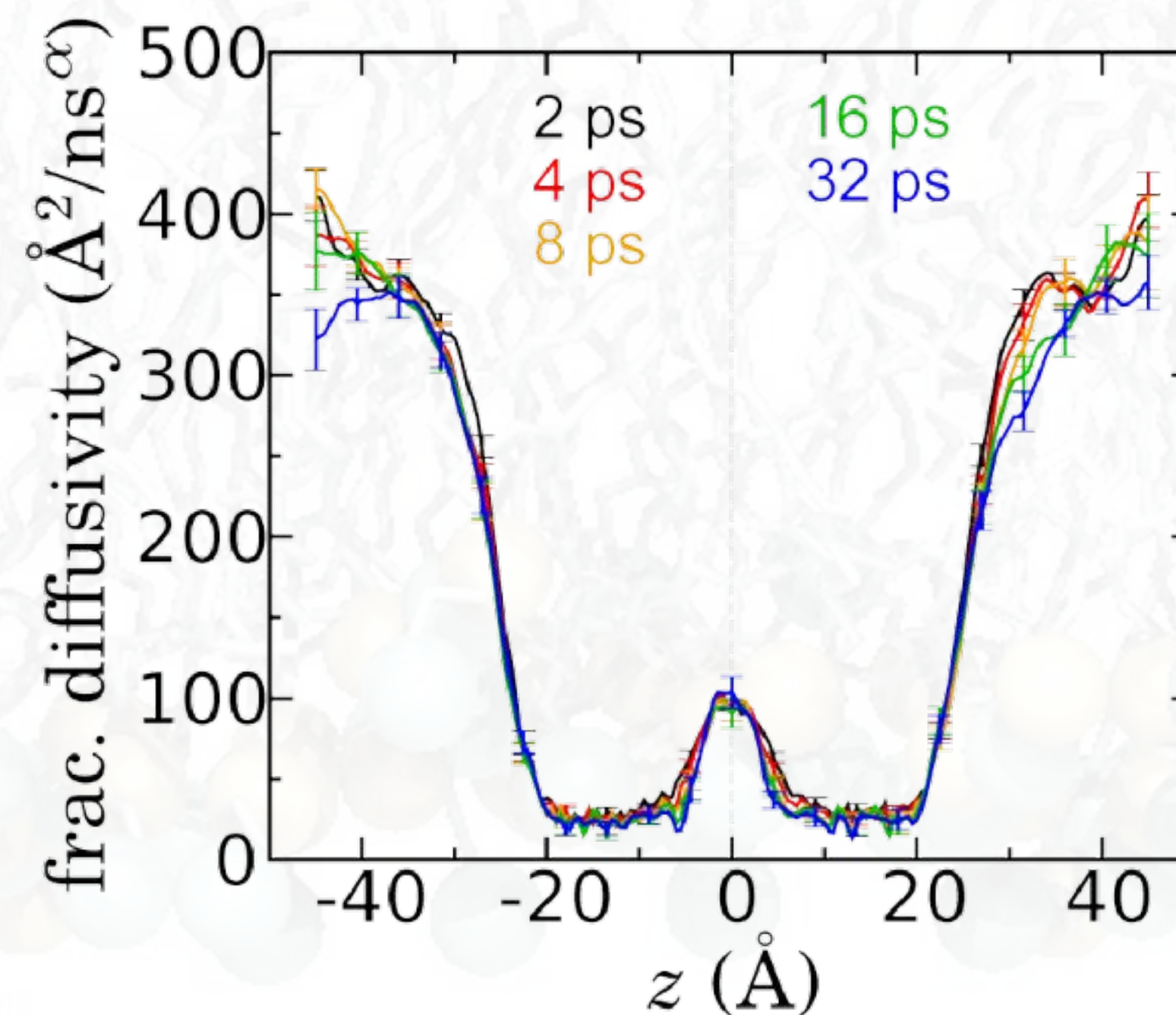
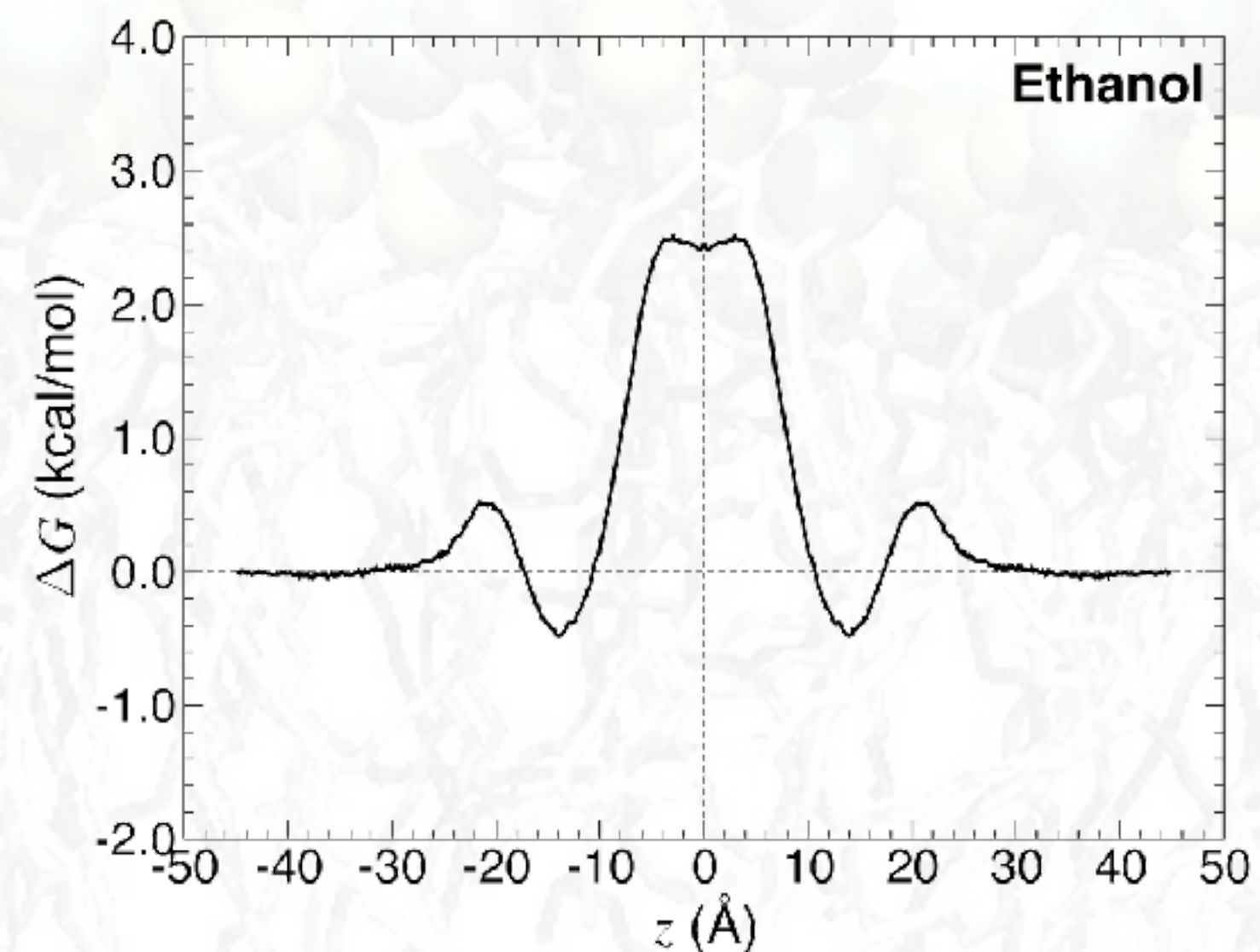
MEMBRANE PERMEABILITY TO SMALL ALCOHOLS

Naive model of the reaction coordinate: Euclidian distance separating the center of mass of the alcohol to that of the bilayer, projected onto z .

Permeation does not obey a random walk: $\langle z(t)^2 \rangle \sim K_\alpha t^\alpha$

Under these premises, one needs to turn to an alternate theoretical framework, based on fractional diffusion:

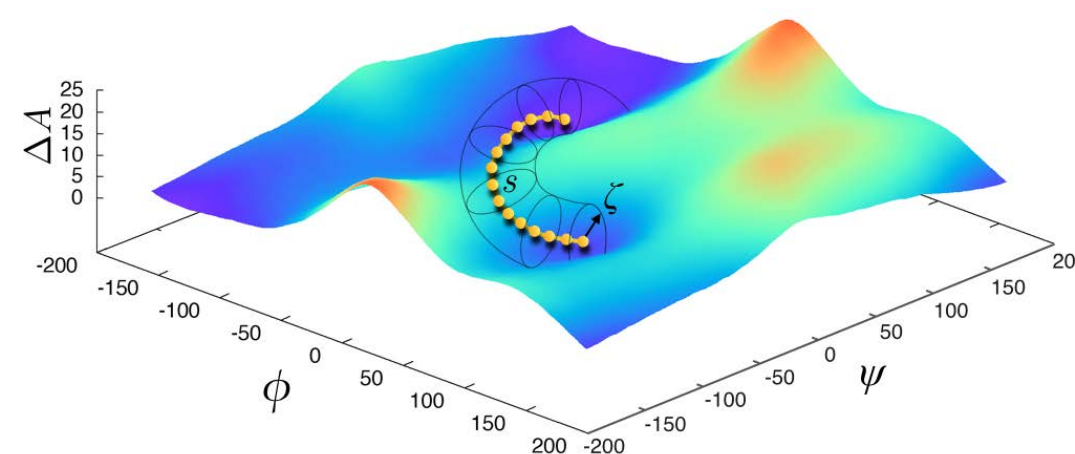
$$\partial_t^{\alpha(z)} c(z, t) = \partial_z [K_\alpha(z) \partial_z - \beta K_\alpha(z) F(z, t)] c(z, t)$$



path sampling
 string method with swarm of trajectories, free-energy calculations along a path-collective variable
advanced tutorial

Department of Biochemistry and Molecular Biology
 Gordon Center for Integrative Science
 The University of Chicago
 Centre National de la Recherche Scientifique
 Laboratoire International Associé CNRS-UIUC
 Université de Lorraine
 University of Illinois at Urbana-Champaign
 Beckman Institute for Advanced Science and Technology
 Theoretical and Computational Biophysics Group

**String method with swarms of trajectories:
 A tutorial for free-energy calculations along a
 minimum-action path**



**Mikolai Fajer
 Jérôme Héning
 Benoît Roux
 Christophe Chipot**

August 19, 2015

Please visit www.ks.uiuc.edu/Training/Tutorials/ to get the latest version of this tutorial, to obtain more tutorials like this one, or to join the tutorial-1@ks.uiuc.edu mailing list for additional help.

Contributors: Gumbart, J. C.; Héning, J.; Fajer, M.; Roux, B.; Chipot, C.