TRANSITION-PATH SAMPLING AND FREE-ENERGY CALCULATIONS

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HANDS-ON WORKSHOP ON ENHANCED SAMPLING AND FREE-ENERGY CALCULATIONS NIH CENTER FOR MACROMOLECULAR MODELING & BIOINFORMATICS, URBANA, ILLINOIS, SEPTEMBER 2017

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WHAT IS A GOOD REACTION-COORDINATE MODEL ?

- Reaction coordinate versus order parameter
- Committor distributions

THE STRING METHOD

- The basic string method
- The string method with swarms of trajectories

EXTRACTING THE FREE ENERGY FROM THE MINIMUM-ACTION PATH

- Path-collective variables
- Perturbative approach

APPLICATIONS

- Transition path of activation loop in c-Src kinase
- Chemomechanical coupling in V₁-ATPase

RECONCILING THERMODYNAMICS AND KINETICS









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WHY DO WE NEED REACTION COORDINATES?

Capture both the kinetics and the thermodynamics of complex chemical and biological processes

Intricate transitions between metastable states often mirror substantial collectivity and rugged free-energy landscapes



De Donder, T. L'affinité. Gauthier-Villars: Paris, 1927 Kirkwood, J. G. J. Chem. Phys. 1935, 3, 300–313



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INTRODUCTION









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The *true* reaction coordinate generally refers to a unique mathematical object on \mathbb{R}^{3N} .

It defines the minimum free-energy pathway connecting the reference state to the target state of the transformation.

In practice, we coarse-grain the atomic detail: $\{x_1, x_2, \dots, x_N\} \longrightarrow \{z_1, z_2, \dots, z_n\}$, where $n \ll N$



Bolhuis, P. G.; Dellago, C.; Chandler, D. Proc. Natl. Acad. Sci. U. S. A. 2000, 97, 5877-5882 Bolhuis, P. G.; Chandler, D.; Dellago, C.; Geissler, P. Ann. Rev. Phys. Chem. 2002, 59, 291-318



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collective variables





1.1









Bolhuis, P. G.; Dellago, C.; Chandler, D. Proc. Natl. Acad. Sci. U. S. A. 2000, 97, 5877-5882 Bolhuis, P. G.; Chandler, D.; Dellago, C.; Geissler, P. Ann. Rev. Phys. Chem. 2002, 59, 291-318









Including relevant collective variables is absolutely crucial for finding true dynamical pathways.

Committor — The probability to reach the target state before returning to the reference state.



Bolhuis, P. G.; Dellago, C.; Chandler, D. Proc. Natl. Acad. Sci. U. S. A. 2000, 97, 5877-5882 Bolhuis, P. G.; Chandler, D.; Dellago, C.; Geissler, P. Ann. Rev. Phys. Chem. 2002, 59, 291-318



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The transition state surface for A \rightarrow B transitions is formed of configurations where $p_{\rm B} = p_{\rm A} = 1/2$.











While the choice of the reaction-coordinate model does not impact the thermodynamics of the process at hand, it modulates its kinetics.



Option: Determine $N(p_A)$, the distribution of the committor probability, p_A , for the model of the reaction coordinate, ξ .



Run a series of molecular dynamics simulations from the putative maximum of the freeenergy barrier and infer $N(p_A)$.



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Possible linear combination of variables



Degenerate variable

From normal mode or principal component analysis

Hénin, J.; Forin, G.; Chipot, C.; Klein, M. L. J. Chem. Theor. Comput. 2010, 6, 35-47

Fiorin, G.; Klein, M. L.; Hénin, J. Mol. Phys. 2013, 111, 3345-3362









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Movements in molecular objects can be more complex than suggested by

Liu, P.; Shao, X.; Chipot, C.; Cai, W. Chem. Sci. 2015







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THE BASIC STRING METHOD



Let us assume some minimum-action (most probable) transition path connecting two basins of a free-energy landscape defined by a set of collective variables z. Let us also consider the potential of mean force along this path,

$$\exp[-\beta w(z)] = \frac{\int d\mathbf{x} \, \delta[\mathbf{z} - \mathbf{z}'(\mathbf{x})] \, \exp[-\beta w(z)]}{\int d\mathbf{x} \, \exp[-\beta U(\mathbf{x})]}$$

 $\mathbf{F} = -\boldsymbol{\nabla}_z w(z)$

Evolve the string of images until:

$$\left(\mathbf{DF}\right)^{\perp} = 0$$

Maragliano, L.; Fischer, A.; Vanden-Eijnden, E.; Ciccotti, G. J. Chem. Phys. 2006, 125, 024106



$$U(\mathbf{x})]$$











Let us assume that the collective variables evolve on the free-energy landscape according to non-inertial Brownian dynamics, $| + \nabla_{z_i} D_{ij} [\mathbf{z}(0)] \} \delta t + R_i(0)$ where $F_i = -\nabla_i w(\mathbf{z})$ \mathcal{Z} $\langle R_i(0)R_i(\delta t)\rangle = 2D_{ij}\delta t$

$$z_i(\delta t) = z_i(0) + \sum_j \left\{ \beta D_{ij}[\mathbf{z}(0)] F_j[\mathbf{z}(0)] \right\}$$

Let us consider a path $\mathbf{z}(\alpha)$ connecting the two basins, such that α varies between 0 and 1, $z_i(\alpha) = z_i(\alpha') + \sum_{i} \left\{ \beta D_{ij}[\mathbf{z}(0)] F_j[\mathbf{z}(0)] + \nabla_{z_j} D_{ij}[\mathbf{z}(0)] \right\} \delta t$

Average drift from an ensemble of unbiased trajectories of length δt , $\overline{z_i(\delta t)} = \overline{z_i(\delta t) - z_i(0)} = \sum_i \left\{ \beta D_{ij}[\mathbf{z}(0)] F_j[\mathbf{z}(0)] + \nabla_{z_j} D_{ij}[\mathbf{z}(0)] \right\} \delta t$

Pan, A. C.; Sezer, D.; Roux, B. J. Phys. Chem. B 2008, 112, 3432-3440 Ren, W.; Vanden-Eijnden, E.; Maragakis, P.; E, W. J. Chem. Phys. 2005, 123, 134109















Prepare a configuration for each one of the P images of the string, the corresponding collective variables of which are close to \mathbf{z}_i , for i = 1, ..., P.

Generate an equilibrium trajectory for each image with z restrained around z_i .

From the equilibrium trajectory, generate a large number of short, unbiased trajectories

Use the resulting average displacement, $\overline{\Delta z_i}$, to determine the position of the P images.

Parameterize the string to ensure that the images are equidistant in collective-variable

Pan, A. C.; Sezer, D.; Roux, B. J. Phys. Chem. B 2008, 112, 3432-3440







Reparametrization prevents the path from becoming under-resolved, especially near free energy barriers

E, W.; Ren, W.; Vanden-Eijnden, E. Phys. Rev. B 2002, 66, 052301 Maragliano, L.; Fischer, A.; Vanden-Eijnden, E.; Ciccotti, G. J. Chem. Phys. 2006, 125, 024106











Pan, A. C.; Sezer, D.; Roux, B. J. Phys. Chem. B 2008, 112, 3432-3440 Ren, W.; Vanden-Eijnden, E.; Maragakis, P.; E, W. J. Chem. Phys. 2005, 123, 134109





















DERIVING THE ONE-DIMENSIONAL FREE-ENERGY PROFILE



Singharoy, A.; Chipot, C. J. Phys. Chem. B 2017, 121, 3502-3514









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Path collective variables:

$$s(\mathbf{z}) = \lim_{\lambda \to \infty} \frac{\int_0^1 \mathrm{d}t \ t \ \exp\{-\lambda[\mathbf{z} - \mathbf{z}(t)]^2\}}{\int_0^1 \mathrm{d}t \ \exp\{-\lambda[\mathbf{z} - \mathbf{z}(t)]^2\}}$$
Free-energy

$$\zeta(\mathbf{z}) = \lim_{\lambda \to \infty} -\frac{1}{\lambda} \int_0^1 dt \, \exp\{-\lambda [\mathbf{z} - \mathbf{z}(t)]^2\} \qquad F(s, \zeta) =$$

Branduardi, D.; Gervasio, F. L.; Parrinello, M. J. Chem. Phys. 2007, 126, 054103



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gy surface:

Closest minimum free-energy path:

$$-\frac{1}{\beta}\ln\left\langle\delta\left[s-s(\mathbf{z})\right]\delta\left[\zeta-\zeta(\mathbf{z})\right]\right\rangle$$

$$\mathscr{T}(\mathbf{z}) = \int_0^1 \mathrm{d}s \ F(s,\zeta)$$

Variational principle:
$$\frac{\delta \mathscr{T}}{\delta}$$











DERIVING THE ONE-DIMENSIONAL FREE-ENERGY PROFILE

Discrete case: The putative reaction coordinate, z(t), is discretized in a collection of images.

$$s(\mathbf{z}) = \frac{1}{N-1} \frac{\sum_{i} (i-1) \exp\left\{-\lambda \left[\mathbf{z} - \mathbf{z}(i)\right]^{2}\right\}}{\sum_{i} \exp\left\{-\lambda \left[\mathbf{z} - \mathbf{z}(i)\right]^{2}\right\}}$$
$$\zeta(\mathbf{z}) = -\frac{1}{\lambda} \ln\left(\sum_{i} \exp\left\{-\lambda \left[\mathbf{z} - \mathbf{z}(i)\right]^{2}\right\}\right)$$

Closest minimum free-energy path:

$$\mathscr{T}(\mathbf{z}) = \frac{1}{N} \left\{ \frac{1}{2} \left[F\left(s(\mathbf{z}(1)), 0\right) + F\left(s(\mathbf{z}(N)), 0\right) \right] \sum_{i=-2}^{N-1} F\left(s(\mathbf{z}(i))\right) \right\}$$

Branduardi, D.; Gervasio, F. L.; Parrinello, M. J. Chem. Phys. 2007, 126, 054103













TRANSITION-PATH SAMPLING AND FREE-ENERGY CALCULATIONS

DERIVING THE ONE-DIMENSIONAL FREE-ENERGY PROFILE



From the integration of C_{7eq} and C_{7ax} basins, $\Delta\Delta A = 2.5$ kcal/mol. From the difference of RMSD's, $\Delta A^{\ddagger} = 5.6$ kcal/mol.

Hénin, J.; Fiorin, G.; Chipot, C.; Klein, M. L. J. Chem. Theor. Comput. 2010, 6, 35-47



TRANSITION-PATH SAMPLING AND FREE-ENERGY CALCULATIONS

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Free energy of the biased system, or perturbed free energy, $\mathrm{e}^{-\beta F(\boldsymbol{\eta})} = \int \mathrm{d}\boldsymbol{\zeta} \mathrm{e}^{-\beta \left(\mathrm{G}(\boldsymbol{\zeta}) + \mathrm{U}_{\boldsymbol{\eta}}(\boldsymbol{\zeta})\right)}$

 $G(\boldsymbol{\zeta}) \approx F(\boldsymbol{\zeta}) + \frac{1}{2\beta k} \Big(\beta \boldsymbol{\nabla}_{\boldsymbol{\zeta}} F(\boldsymbol{\zeta}) \cdot \boldsymbol{\nabla}_{\boldsymbol{\zeta}} F(\boldsymbol{\zeta}) - \boldsymbol{\nabla}_{\boldsymbol{\zeta}}^2 F(\boldsymbol{\zeta}) \Big)$

Assuming that $\zeta(s)$ approximately represents the minimum free-energy path, and s is its arc-length,

$$G\left(\boldsymbol{\zeta}(s)\right) \approx F(s) + \frac{1}{2\beta k} \left\{ \beta \left(\frac{\mathrm{d}}{\mathrm{ds}}F(s)\right)^2 - \frac{\mathrm{d}^2}{\mathrm{ds}^2}F(s) \right\}$$

The perturbed free energies, $F_i = F(\zeta(s_i))$, can be estimated by solving self-consistently,

$$e^{-\beta F_i} = \sum_t \frac{e^{-\beta U_i(\boldsymbol{\zeta}^t)}}{\sum_j T_j e^{-\beta \left(U_j(\boldsymbol{\zeta}^t) - F_j\right)}} \qquad w^t = \left(\sum_i T_i e^{-\beta \left(U_i(\boldsymbol{\zeta}^t) - F_j\right)}\right)$$

$$G(\boldsymbol{\xi}) = -\beta^{-1} \ln \left\{ \sum_{t} w^{t} \mathbf{K} \left(\boldsymbol{\xi}(\mathbf{x}^{t}) - \boldsymbol{\xi} \right) \right\}$$



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$$e^{-\beta G(\boldsymbol{\zeta})} \propto e^{-\frac{1}{2\beta k} \boldsymbol{\nabla}_{\boldsymbol{\zeta}}^2} e^{-\beta F(\boldsymbol{\zeta})}$$

For large k, in the stiff-spring approximation, one may expand the PMF to extract the first two terms in 1/k,

 $\left(F_{i} \right) - F_{i} \right) \right)^{-1}$

Alternatively,
$$\begin{cases} w^{t} = \left(\sum_{i} T_{i} e^{-\beta \left(U_{i}(\boldsymbol{\zeta}^{t}) - F_{i}\right)}\right) \\ e^{-\beta F_{i}} = \sum_{i} w^{t} e^{-\beta U_{i}(\boldsymbol{\zeta}^{t})} \end{cases}$$

Hummer, G.; Szabo, A. Proc. Natl. Acad. Sci. U. S. A. 2010, 107, 21441–21446 Bartels, C. Chem. Phys. Lett. 2000, 331, 446-454 Shirts, M. R.; Chodera, J. D. J. Chem. Phys. 2008, 129, 124105 Moradi, M.; Tajkhorshid, E. J. Chem. Theory Comput. 2014, 10, 2866–2880.





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APPLICATIONS



TRANSITION PATH OF ACTIVATION LOOP IN C-SRC KINASE



Meng, Y.; Pond, M. P.; Roux, B. Acc. Chem. Res. 2017, 50, 1193-1201



- String has 51 images
- Each swarm contains 100 trajectories, 1 ps each
- 50 ps for constrained equilibration
- Pathways is evolved for 100 iterations
- MD simulations are done with NAMD
- All simulations with explicit solvent









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BEYOND THERMODYNAMICS

$$\Delta Z = \beta D(Z_t) F(Z_t, t) \Delta t + \nabla D(Z_t) \Delta t + \sqrt{2D(Z_t) \Delta t} g_t$$

Let:
$$\begin{cases} \mu = \beta D(Z_t) F(Z_t, t) \Delta t + \nabla D(Z_t) \Delta t \\ \sigma^2 = 2D(Z_t) \Delta t \end{cases}$$

Then: $\Delta Z = \mu + \sigma g_t$

$$P[\Delta Z|w(z), D(z)] = \frac{1}{\sigma\sqrt{2\pi}} \exp\left(-\frac{(\Delta Z - \mu)^2}{2\sigma^2}\right)$$

Probability over the entire trajectory, given the parameters:

$$P[Z(t)|w(z), D(z)] = \prod_{i} \frac{1}{\sigma_i \sqrt{2\pi}} \exp\left(-\frac{(\Delta Z_i - \mu_i)^2}{2\sigma_i^2}\right)$$

Comer, J. R.; Chipot, C. J.; González-Nilo, F. D. J. Chem. Theory Comput. 2013, 9, 876-882 Hummer, G. New J. Phys. 2005, 7, 34



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(1) The molecular dynamics supplies the trajectory of the collective variable, Z(t).

(2) The molecular dynamics supplies also, $f_{\text{bias}}(t)$.

(3) Pick trial parameters, w(z) and D(z).

(4) Assume a propagator, e.g., Brownian dynamics.

(5) Calculate the probability of the trajectory given the parameters.

(6) Bayes's theorem: Get the probability of the parameters given the trajectory.

(7) Optimize the parameters to yield the greatest probability.

Ermak, D.; McCammon, J. J. Chem. Phys. 1978, 69, 1352-1360 Türkcan, S.; Alexandrou, A.; Masson, J. Biophys. J. 2012, 102, 2288-2298





BEYOND THERMODYNAMICS

Mean first passage time:

$$\tau = \int_{a}^{b} \mathrm{d}\xi \exp\left[\beta \Delta A(\xi)\right] D^{-1}(\xi) \int_{a}^{\xi} \mathrm{d}\xi' \exp\left[-\beta \Delta A(\xi')\right]$$



Szabo, A.; Schulten, K.; Schulten, Z. J. Chem. Phys. 1980, 72, 4350-4357 Singharoy, A.; Chipot, C.; Moradi, M.; Schulten, K. J. Am. Chem. Soc. 2017, 139, 293-310



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$$k = 1/\tau$$





Hénin, J.; Tajkhorshid, E.; Schulten, K.; Chipot, C. Biophys. J. 2008, 94, 832-839





MEMBRANE PERMEABILITY TO SMALL ALCOHOLS

Naive model of the reaction coordinate: Euclidian distance separating the center of mass of the alcohol to that of the bilayer, projected onto z.

Permeation does not obey a random walk: $\langle z(t)^2 \rangle \sim K_{\alpha} t^{\alpha}$

Under these premises, one needs to turn to an alternate theoretical framework, based on fractional diffusion:

$$\partial_t^{\alpha(z)} c(z,t) = \partial_z \left[K_\alpha(z) \partial_z - \beta K_\alpha(z) F(z,t) \right] c(z,t)$$

Comer, J.; Chipot, C.. J. Chem. Theory Comput. 2017, 13, 2523-2532









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path sampling advanced tutorial

Department of Biochemistry and Molecular Biology Gordon Center for Integrative Science The University of Chicago

Centre National de la Recherche Scientifique Laboratoire International Associé CNRS-UIUC Université de Lorraine

University of Illinois at Urbana-Champaign Beckman Institute for Advanced Science and Technology Theoretical and Computational Biophysics Group

> String method with swarms of trajectories: A tutorial for free-energy calculations along a minimum-action path



Please visit www.ks.uiuc.edu/Training/Tutorials/ to get the latest version of this tutorial, to obtain more tutorials like this one, or to join the tutorial-l@ks.uiuc.edu mailing list for additional help.

Contributors: Gumbart, J. C.; Hénin, J.; Fajer, M.; Roux, B.; Chipot, C.



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string method with swarm of trajectories, free-energy calculations along a path-collective variable

Mikolai Fajer Jérôme Hénin Benoît Roux **Christophe Chipot** August 19, 2015





TUTORIALS