# **Introduction to Molecular Dynamics**





# **PDB** Files

#### gives one the structure and starting position

- Simulations start with a crystal structure from the Protein Data Bank, in the standard PDB file format.
- PDB files contain standard records for species, tissue, authorship, citations, sequence, secondary structure, etc.
- We only care about the atom records...
  - atom name (N, C, CA)
  - residue name (ALA, HIS)
  - residue id (integer)
  - coordinates (x, y, z)
  - occupancy (0.0 to 1.0)
  - temp. factor (a.k.a. beta)
  - segment id (6PTI)
- No hydrogen atoms!

(We must add them ourselves.)



# Potential Energy Function of Biopolymer

- Simple, fixed algebraic form for every type of interaction.
- Variable parameters depend on types of atoms involved.



## Preparing Your System for MD Solvation

Biological activity is the result of interactions between molecules and occurs at the interfaces between molecules (protein-protein, protein-DNA, protein-solvent, DNA-solvent, etc).

Why model solvation?

• many biological processes occur in aqueous solution

• solvation effects play a crucial role in determining molecular conformation, electronic properties, binding energies, etc

How to model solvation?

- explicit treatment: solvent molecules are added to the molecular system
- implicit treatment: solvent is modeled as a continuum dielectric or so-called implicit force field

mitochondrial bc1 complex



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*mitochondrial bc1 complex* 



(Usually periodic! Avoids surface effects)





## Cutting Corners cutoffs, PME, rigid bonds, and multiple timesteps

- Nonbonded interactions require order N<sup>2</sup> computer time!
  - Truncating at  $R_{cutoff}$  reduces this to order N  $R_{cutoff}{}^3$
  - Particle mesh Ewald (PME) method adds long range electrostatics at order N log N, only minor cost compared to cutoff calculation.
- Can we extend the timestep, and do this work fewer times?
  - Bonds to hydrogen atoms, which require a 1fs timestep, can be held at their equilibrium lengths, allowing 2fs steps.
  - Long range electrostatics forces vary slowly, and may be evaluated less often, such as on every second or third step.
- Coarse Graining



#### **Residue-Based Coarse-Grained Model**



- Lipid model: MARTINI
- Level of coarse-graining: ~4 heavy atoms per CG bead
- Interactions parameterized based on experimental data and thermodynamic properties of small molecules

- Protein model uses two CG beads per residue
- One CG bead per side chain another for backbone



All-atom peptide

CG peptide

Peter L. Freddolino, Anton Arkhipov, Amy Y. Shih, Ying Yin, Zhongzhou Chen, and Klaus Schulten. **Application of residue-based and shape-based coarse graining to biomolecular simulations.** In Gregory A. Voth, editor, *Coarse-Graining of Condensed Phase and Biomolecular Systems*, chapter 20, pp. 299-315. Chapman and Hall/CRC Press, Taylor and Francis Group, 2008.

#### Nanodisc Assembly CG MD Simulation

- 10 µs simulation
- · Assembly proceeds in two steps:
  - Aggregation of proteins and lipids driven by the hydrophobic effect
  - Optimization of the protein structure driven by increasingly specific protein-protein interactions
- Formation of the generally accepted double-belt model for discoidal HDL





A. Shih, A. Arkhipov, P. Freddolino, and K. Schulten. J. Phys. Chem. B, 110:3674–3684, 2006; A. Shih, P. Freddolino, A. Arkhipov, and K. Schulten. J. Struct. Biol., 157:579–592,2007; A. Shih, A. Arkhipov, P. Freddolino, S. Sligar, and K. Schulten. Journal of Physical Chemistry B, 111: 11095 - 11104, 2007; A. Shih, P. Freddolino, S. Sligar, and K. Schulten. Nano Letters, 7:1692-1696, 2007.







# Example: MD Simulations of the K+ Channel Protein

Ion channels are membrane spanning proteins that form a pathway for the flux of inorganic ions across cell membranes.

Potassium channels are a particularly interesting class of ion channels, managing to distinguish with impressive fidelity between K<sup>+</sup> and Na<sup>+</sup> ions while maintaining a very high throughput of K<sup>+</sup> ions when gated.



# Setting up the system (1)



• retrieve the PDB (coordinates) file from the Protein Data Bank

• add hydrogen atoms using PSFGEN

• use psf and parameter files to set up the structure; needs better than available in Charmm to describe well the ions

• minimize the protein structure using NAMD2







# Simulating the system: Free MD

Summary of simulations:

• protein/membrane system contains 38,112 atoms, including

5117 water molecules, 100 POPE and 34 POPG lipids, plus K<sup>+</sup> counterions

- CHARMM26 forcefield
- periodic boundary conditions, PME electrostatics
- 1 ns equilibration at 310K, NpT
- 2 ns dynamics, NpT

Program: NAMD2

Platform: Cray T3E (Pittsburgh Supercomputer Center) or local computer cluster; choose ~1000 atoms per processor.



indicate that the most flexible parts of the protein are the N and C terminal ends, residues 52-60 and residues 84-90. Residues 74-80 in the selectivity filter have low temperature factors and are very stable during the simulation.



### **Equilibrium Properties of Proteins**





#### **Thermal Motion of Ubiquitin from MD** Temperature Dependence of Crystal Diffraction (Debye-Waller factor)

Bragg's law  
$$2d\sin\theta = \lambda$$

structure factor

$$f_j \exp[-i\vec{s}\cdot\vec{r_j}]$$

But the atom carries out thermal vibrations around equilibrium position  $\vec{x}_j$ 

$$\vec{r}_j(t) = \vec{x}_j + \vec{u}_j(t)$$

Accordingly:

$$\langle f_j \exp[-i\vec{s}\cdot\vec{r}_j] \rangle = f_j \exp[-i\vec{s}\cdot\vec{x}_j] \langle \exp[-i\vec{s}\cdot\vec{u}_j] \rangle$$

### **Thermal Motion of Ubiquitin from MD**

Temperature Dependence of Crystal Diffraction (Debye-Waller factor)

One can expand:

$$\langle \exp[-i\vec{s}\cdot\vec{u}_j] \rangle = 1 - i \underbrace{\langle \vec{s}\cdot\vec{u}_j \rangle}_{=0} - \frac{1}{2} \langle (\vec{s}\cdot\vec{u}_j)^2 \rangle + ..$$

Spatial average for harmonic oscillator:  $\langle (\vec{s} \cdot \vec{u}_j)^2 \rangle = \frac{1}{3} s^2 \langle u_j^2 \rangle$ 

One can carry out the expansion further and show

$$\langle \exp[-i\vec{s}\cdot\vec{u}_j] \rangle = \exp\left[-\frac{1}{6}s^2\langle\langle u_j^2\rangle\right]$$

Using for the thermal amplitude of the harmonic oscillator

$$\frac{1}{2}m\omega^2 u_j^2 = \frac{3}{2}k_B T$$

one obtains  $\langle f_j \exp \left[-i\vec{s} \cdot \vec{r}_j\right] \rangle = f_j \underbrace{\exp[-s^2 k_B T/2m\omega^2]}_{\exp[-i\vec{s} \cdot \vec{x}_j]} \exp[-i\vec{s} \cdot \vec{x}_j]$ 

# **Equilibrium Properties of Proteins**

Energies: kinetic and potential











#### **Definition of Temperature**

$$\langle \sum_j \frac{1}{2} m_j v_j^2 \rangle = \frac{3}{2} N k_B T$$

$$T = \frac{2}{3N k_B} \left\langle \sum_j \frac{1}{2} m_j v_j^2 \right\rangle$$

The atomic velocities of a protein establish a thermometer, but is it

#### **Temperatur Fluctuations**

Maxwell distribution

$$dP(v_n) = c \exp(-m v_n^2/2k_BT) dv_n \qquad (7)$$

The atomic velocity thermometer is inaccurate due to the finite size of a protein!

> 100 Temperature [K]

0.12

0.00

Individual kinetic energy  $\epsilon_n = m v_n^2/2$ 

$$dP(\epsilon_n) = (\pi T_0 \epsilon_n)^{-1/2} \exp(-\epsilon_n/k_B T_0) d\epsilon_n$$
 (8)

One can derive (temperature T\_0 in units k\_B)

$$\langle \epsilon_n \rangle = T_0/2$$
 (9)

$$\langle \epsilon_n^2 \rangle = 3T_0^2/4$$
 (10)

$$\langle \epsilon_n^2 \rangle - \langle \epsilon_n \rangle^2 = T_0^2/2$$
 (11)

The distribution of the total kinetic energy  $E_{kin} = \sum_j \frac{1}{2} m_j v_j^2$ , according to the central limit theorem, is approximately Gaussian

$$P(E_{kin}) = c \exp \left(\frac{-(E_{kin} - \langle E_{kin} \rangle)^2}{2\left(\frac{3Nk_B^2 T_0^2}{2}\right)}\right) \qquad (12)$$

The distribution function for the temperature  $(T = 2E_{kin}/3k_B)$  fluctuations  $\Delta T = T - T_0$  is then

$$P(\Delta T) = c \exp[-(\Delta T)^2/2\sigma^2], \quad \sigma^2 = 2T^2/3N$$
 (13)

For  $T_0 = 100$ K and N = 557, this gives  $\sigma = 3.6$ .



# Simulated Cooling of Ubiquitin

- Proteins function in a narrow (physiological) temperature range. What happens to them when the temperature of their surrounding changes significantly (temperature gradient) ?
- Can the heating/cooling process of a protein be simulated by molecular dynamics ? If yes, then how?



 What can we learn from the simulated cooling/heating of a protein ?

## How to simulate cooling ?

Heat transfer through mechanical coupling between atoms in the two regions Coolant layer of atoms motion of atoms is subject to stochastic Langevin dynamics  $m\ddot{r} = F_{FF} + F_H + F_f + F_L$  $F_{FF} \rightarrow$  force field  $F_H \rightarrow$  harmonic restrain  $F_f \rightarrow$  friction  $F_L \rightarrow$  Langevin force atoms in the inner region follow Newtonian dynamics  $m\ddot{r} = F_{FF}$ 

# **Simulated Cooling - Result**

t	$\langle T_{sim} \rangle$						
0.05	298.75	1.05	276.00	1.95	267.00	3.25	261.00
0.15	289.25	1.15	276.50	2.05	268.50	3.45	258.50
0.35	285.50	1.25	275.25	2.25	266.50	3.55	259.50
0.55	282.25	1.35	271.00	2.35	264.50	3.95	256,50
0.65	282.75	1.45	271.75	2.55	263.50	4.05	257.25
0.75	279.00	1.65	269.50	2.65	264.50	4.45	254.00
0.85	277.75	1.75	271.00	2.85	262.00	4.55	255.25
1.00	277.50	1.85	268.00	3.05	262.50	4.85	252.00

Result from simulation

Table 1: Mean temperature  $\langle T_{sim} \rangle$  [K] of the protein as a function of time t [ps].



#### Solution of the Heat Equation **Temperature averaged over volume** $\langle T \rangle(t) = \left(\frac{4\pi R^3}{3}\right)^{-1} \int d^3 \mathbf{r} \, T(\mathbf{r},t) = \frac{3}{R^3} \int_0^R r^2 dr \, T(r,t)$ $= T_{bath} + \sum_{n=1}^{\infty} a_n \exp\left[-\left(\frac{n\pi}{R}\right)^2 Dt\right] \frac{3}{R^3} \int_0^R r dr \sin\left(\frac{n\pi r}{R}\right)^2 dt$ $= T_{bath} + 6 \frac{\Delta T}{\pi^2} \sum_{n=1}^{\infty} \frac{1}{n^2} \exp\left[-\left(\frac{n\pi}{R}\right)^2 D t\right]$ 300 290 simulation Temperature [K] theory 280 270 $D \approx 0.38 \times 10^{-3} \mathrm{cm}^2 \mathrm{s}^{-1}$ 260 250 water $1.4 \times 10^{-3} \text{cm}^2 \text{s}^{-1}$ 1 2 3 4 5 0 Time [ps]













