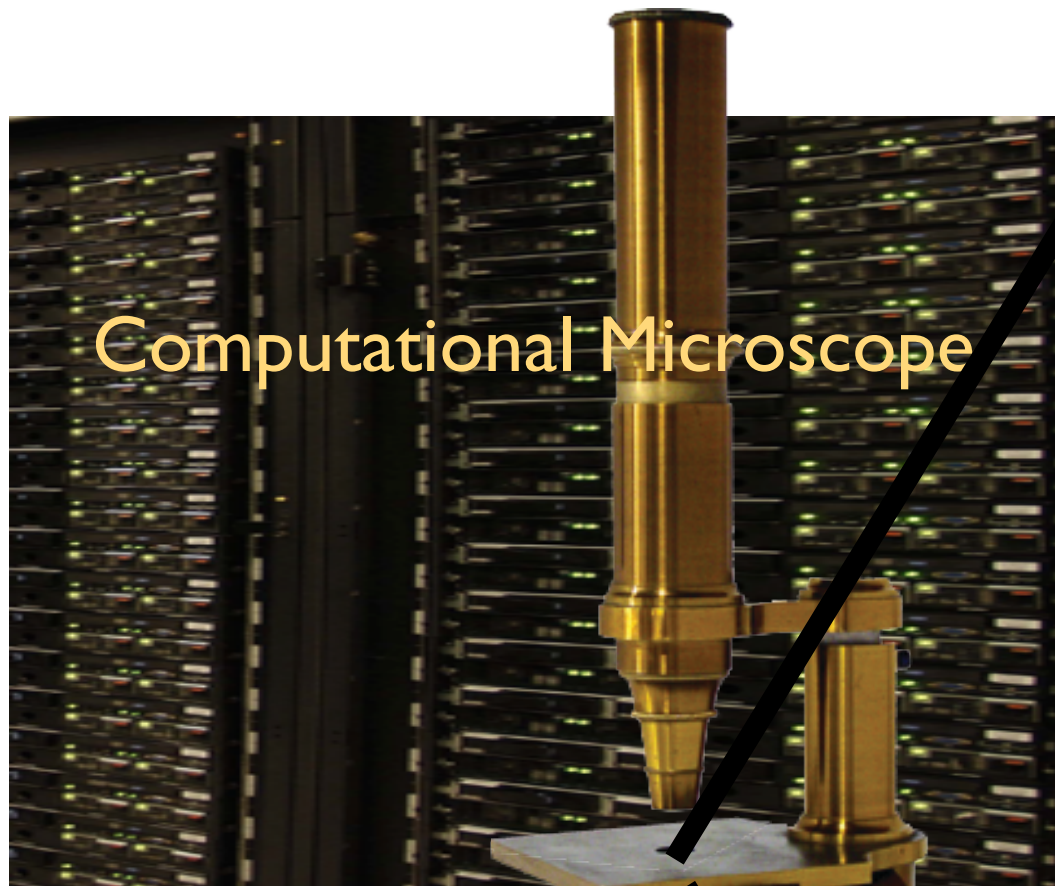


Principles and Applications of Molecular Dynamics Simulations with NAMD

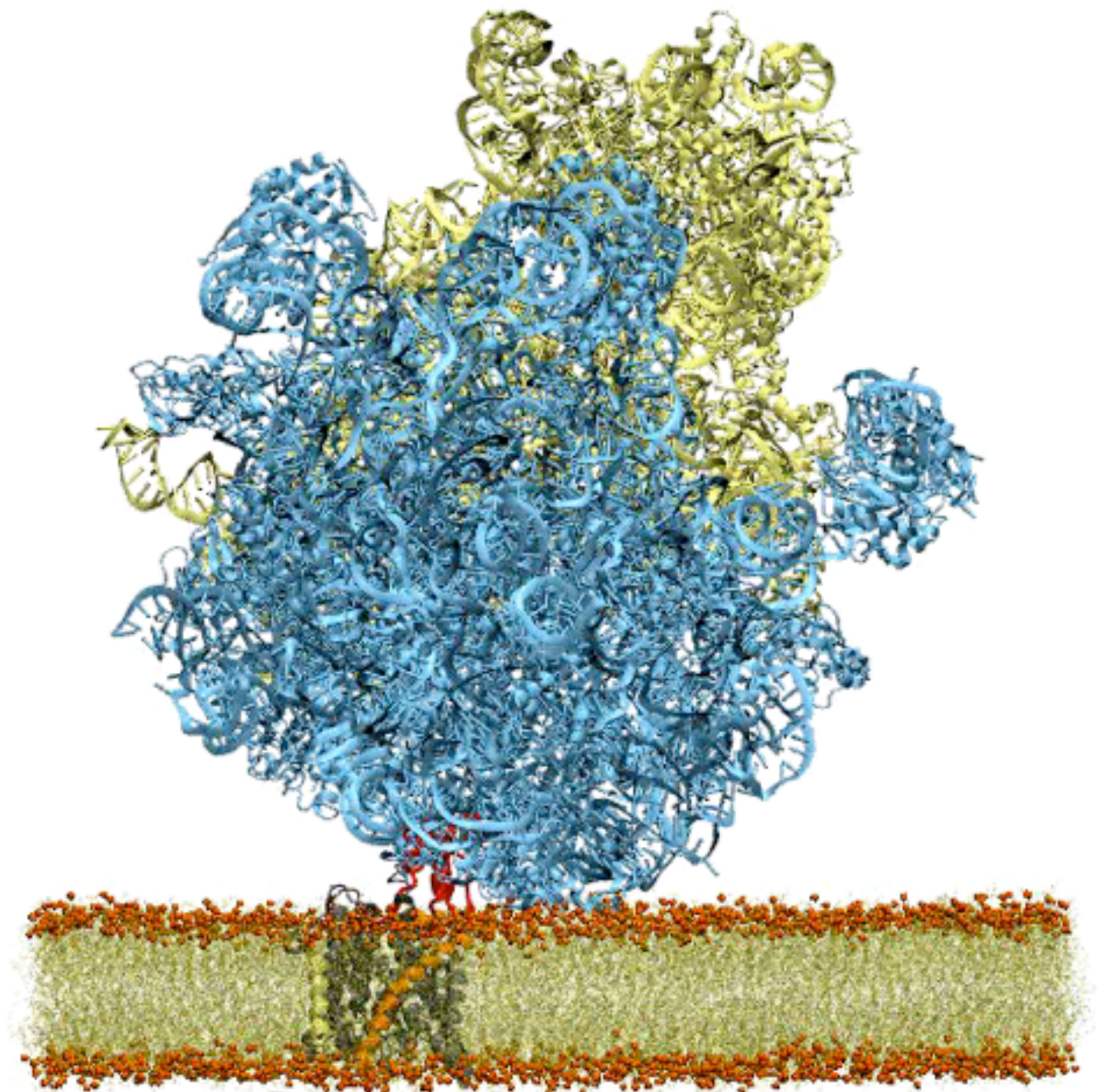
Nov. 14, 2016



NCSA
supercomputer

JC Gumbart

Assistant Professor of Physics
Georgia Institute of Technology
gumbart@physics.gatech.edu
simbac.gatech.edu



These workshops started as one person's dream



Klaus Schulten 1947-2016

"We really feel that you cannot teach just by lecturing; you have to teach through hands-on examples."

This is workshop **#45**
I've taught at workshops 4, 5, 8,
9, 10, 12, 14, 15, 20, 21, 34, 35



June 2003, Urbana



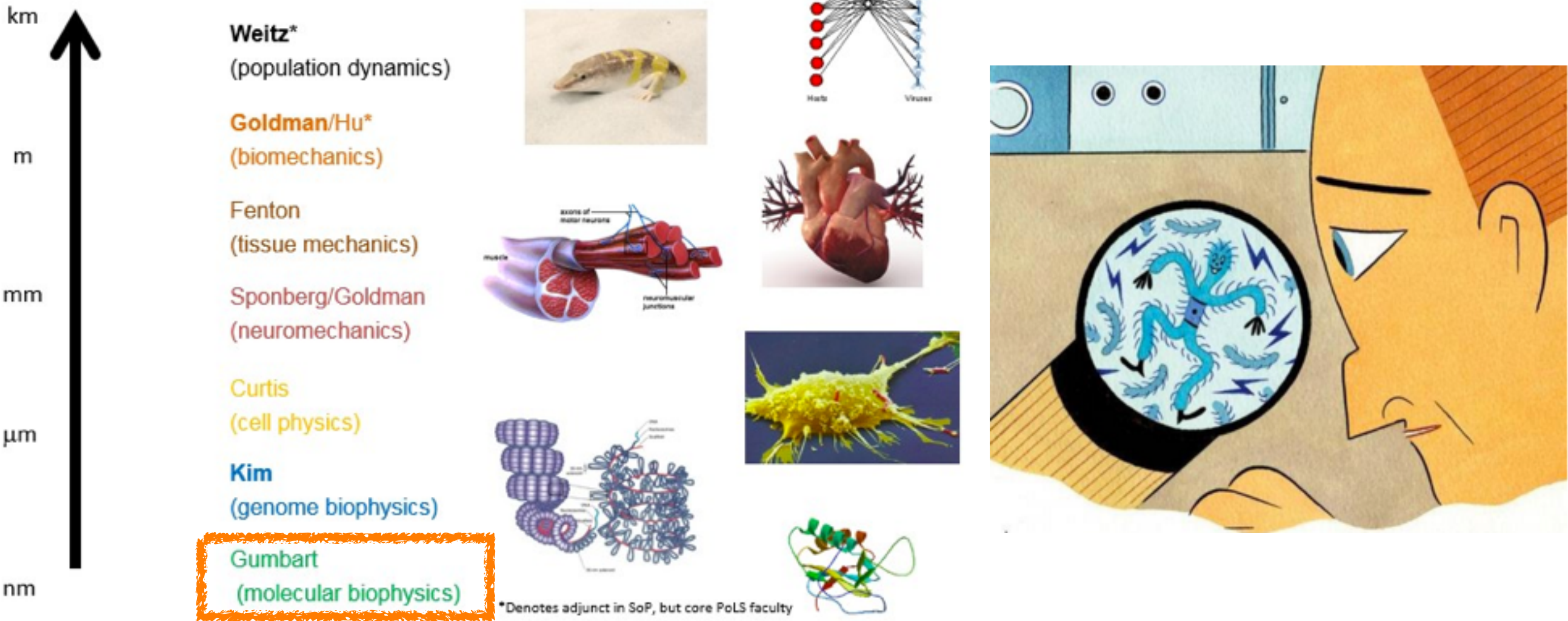
March 2011, Atlanta



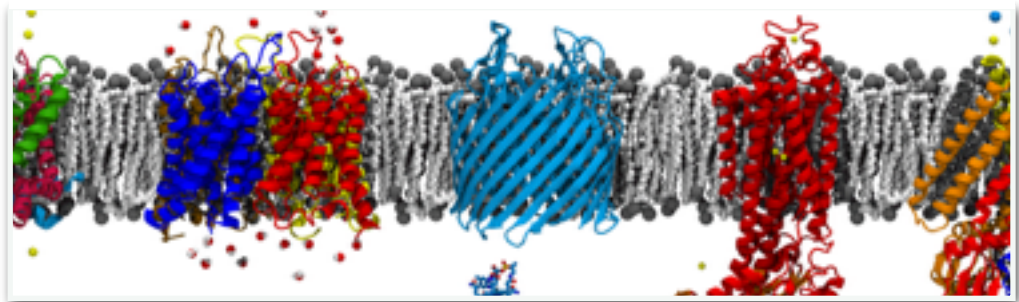
Nov. 2014, Atlanta

Biophysics Research at Georgia Tech

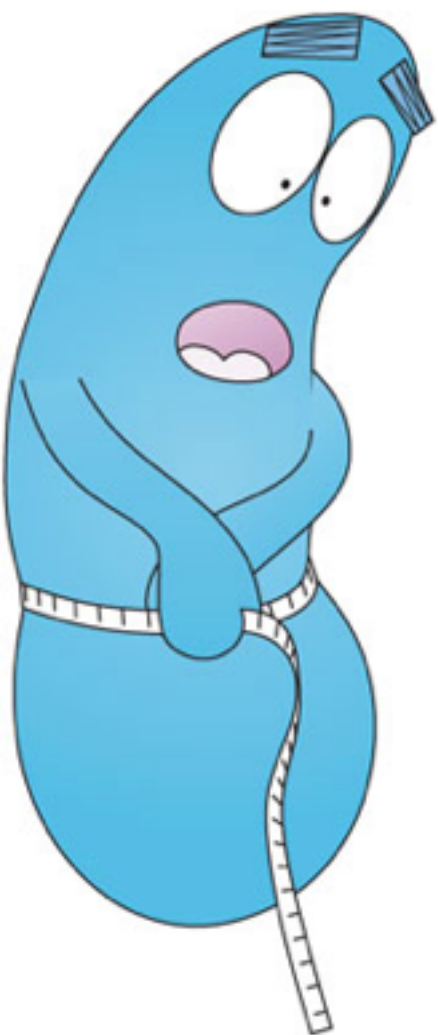
Physics of Living Systems @ GT



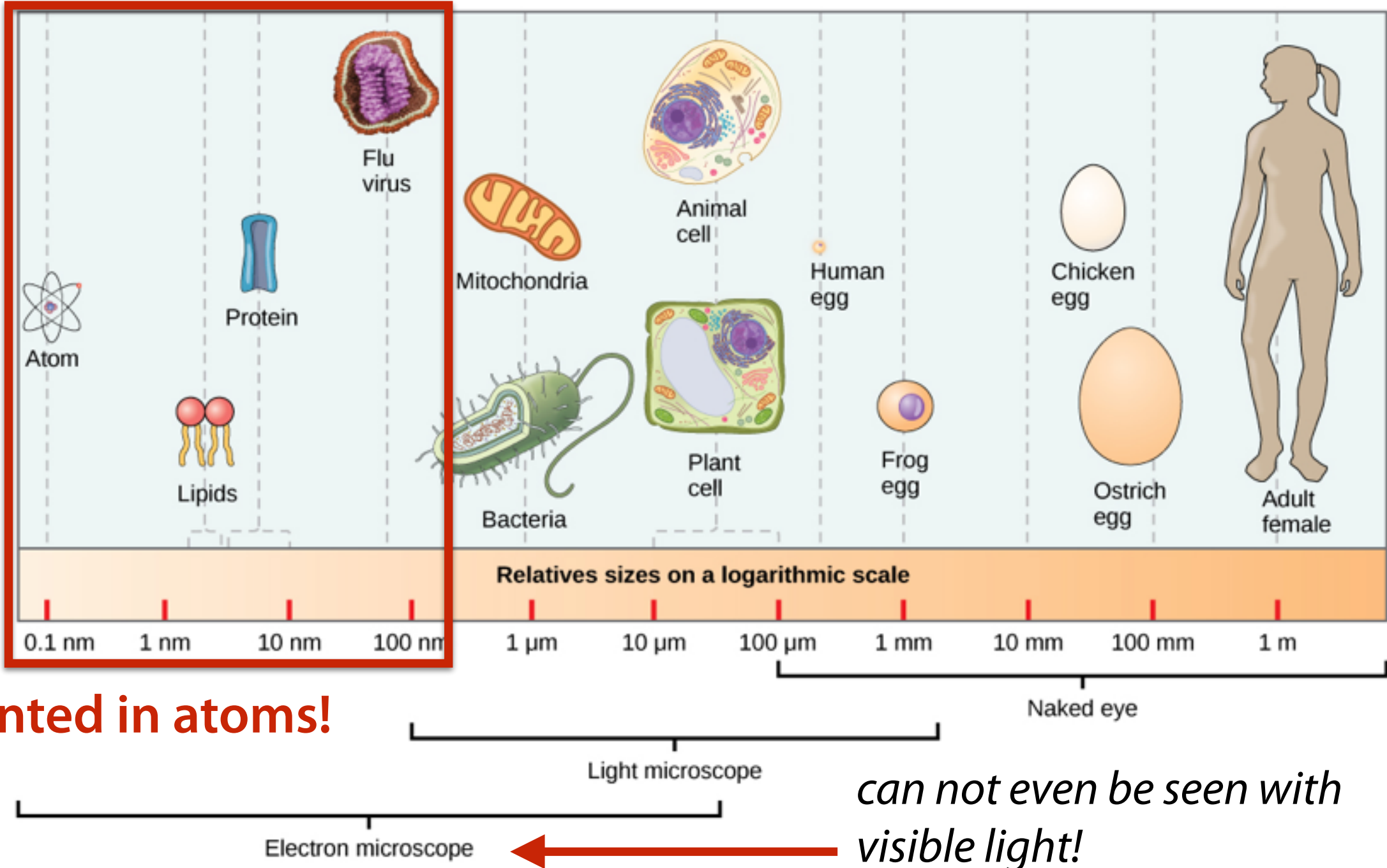
SIMBAC(simulations of bacterial systems)
Lab at Georgia Tech



How small are the things we simulate?



VERY, VERY SMALL!!!!



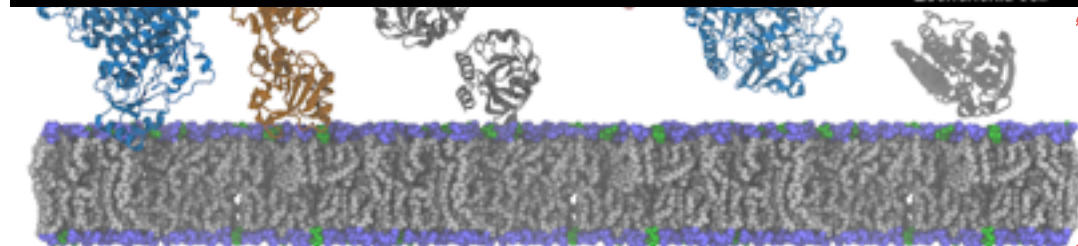
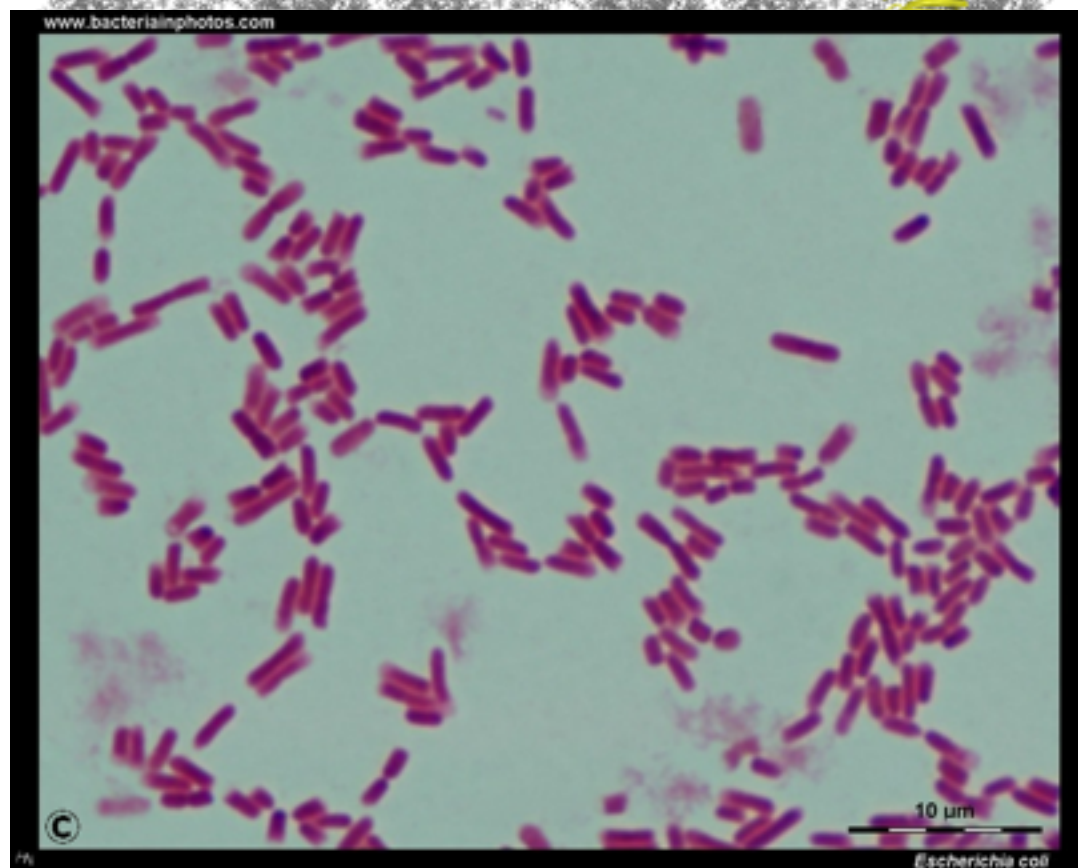
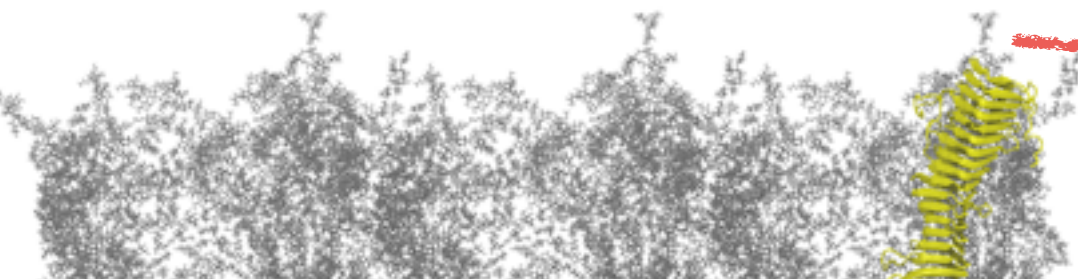
MD simulations as a computational microscope

the traditional light microscope
can **barely** see bacteria

01010100001110111101
01010101010100111011
11010101101010111101

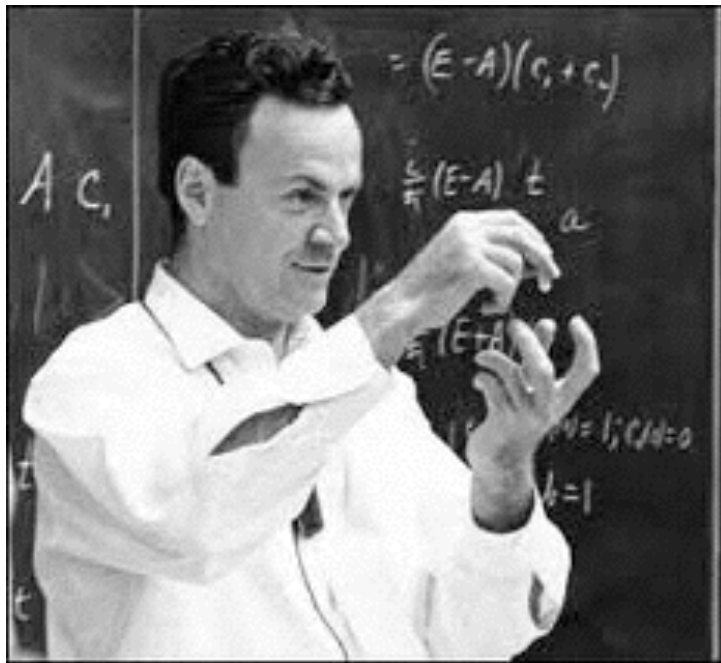
Scalable molecular
dynamics with NAMD.

J. C. Phillips, R. Braun,
W. Wang, J. Gumbart et
al. *J. Comp. Chem.*,
26:1781-1802, 2005.



the computational microscope can
“see” **everything** with infinite resolution

“Everything that living things do can be reduced to wiggling and jiggling of atoms.”



Richard Feynman (1963)

The Nobel Prize in Chemistry 2013



Photo: A. Mahmoud
Martin Karplus
Prize share: 1/3



Photo: A. Mahmoud
Michael Levitt
Prize share: 1/3



Photo: A. Mahmoud
Arieh Warshel
Prize share: 1/3

The Nobel Prize in Chemistry 2013 was awarded jointly to Martin Karplus, Michael Levitt and Arieh Warshel *"for the development of multiscale models for complex chemical systems"*.

50 years later...

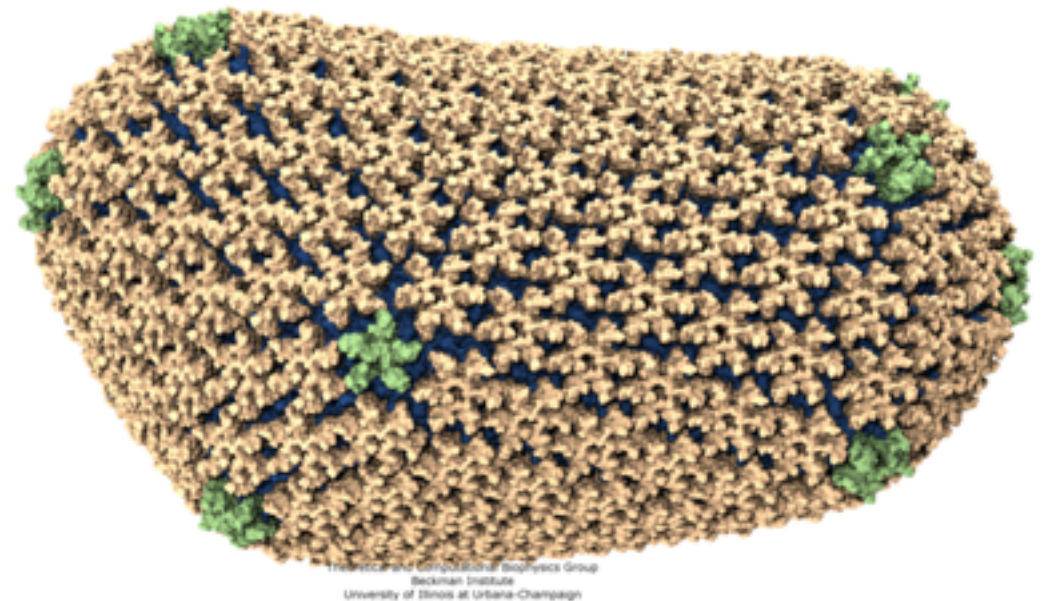
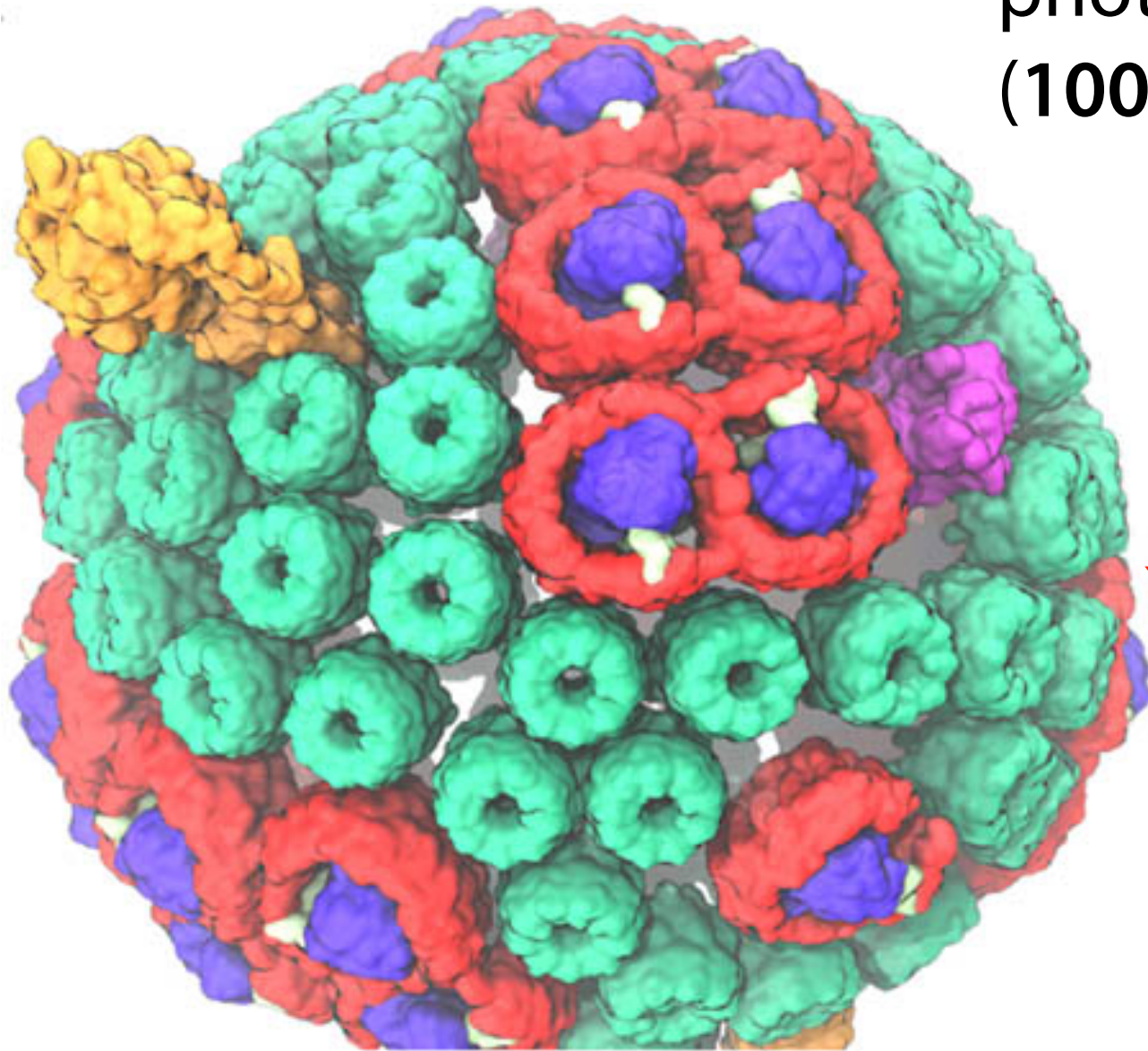
But we are far from done!

State of the art

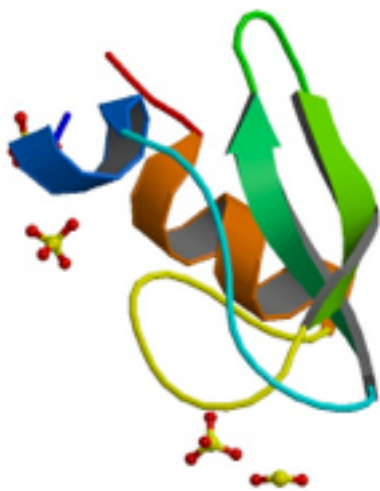
photosynthetic chromatophore
(100 million atoms, 2016)

Sener, Strumpfer, Singharoy, Hunter, Schulten. Overall energy conversion efficiency of a photosynthetic vesicle. *eLife*, 5:e09541, 2016.

HIV virus capsid (64 million
atoms, 2013)



Zhao, Perilla, Yufenyuy, Meng, Chen, Ning, Ahn, Gronenborn, Schulten, Aiken, Zhang. Mature HIV-1 capsid structure by cryo-electron microscopy and all-atom molecular dynamics. *Nature*, 497:643-646, 2013.



individual protein (500 atoms,
1977)

McCammon, Gelin, Karplus. Dynamics of folded proteins. *Nature*, 267:585-590, 1977.

Why do we use MD?

“It is nice to know that the computer understands the problem. But I would like to understand it, too.”

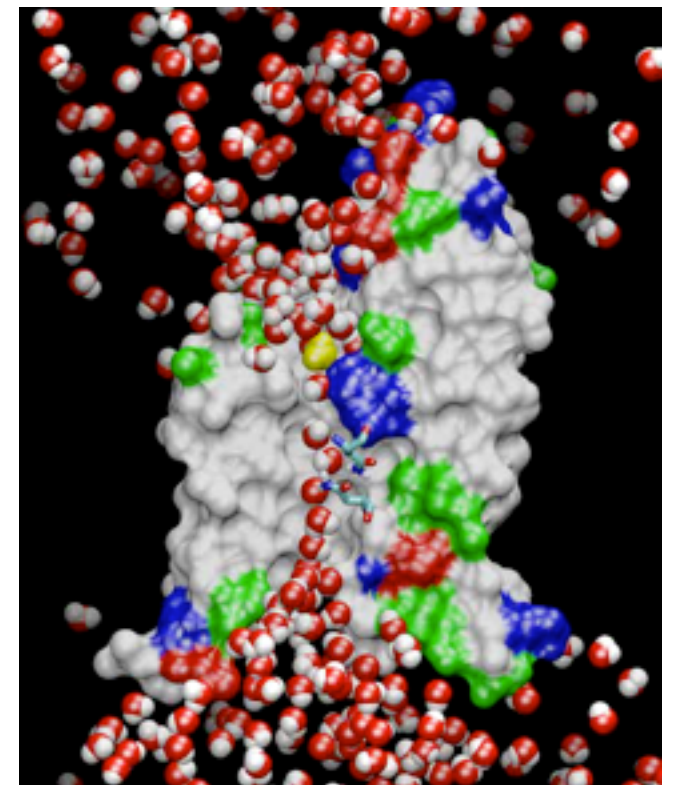
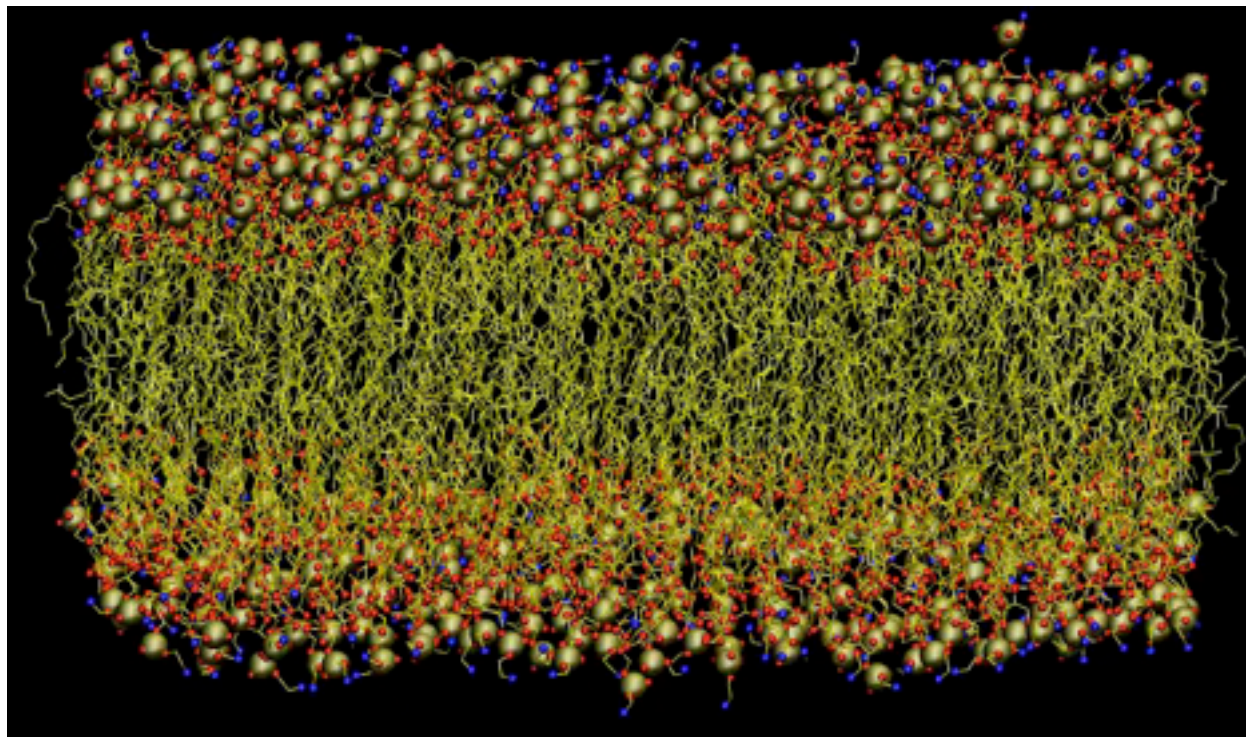


Eugene Wigner

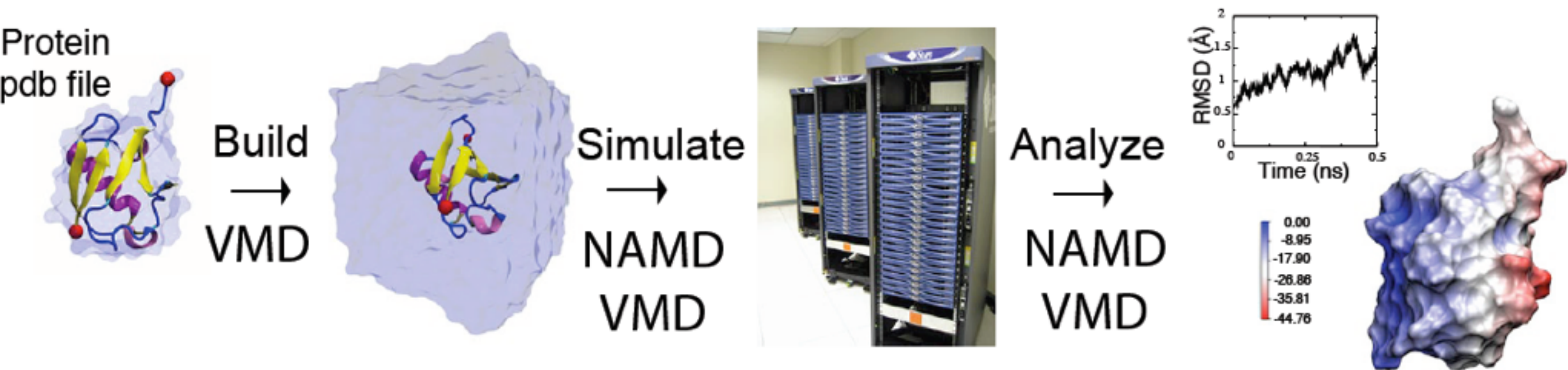
We don't simulate just to watch! But also to measure, analyze, and understand

Why do we use MD?

- Generating a thermodynamic ensemble (sampling / statistics)
- Taking into account fluctuations/dynamics in interpretation of experimental observables
- Describing molecular processes + free energy
- Help with molecular modeling



The Molecular Dynamics Simulation Process



For textbooks see:

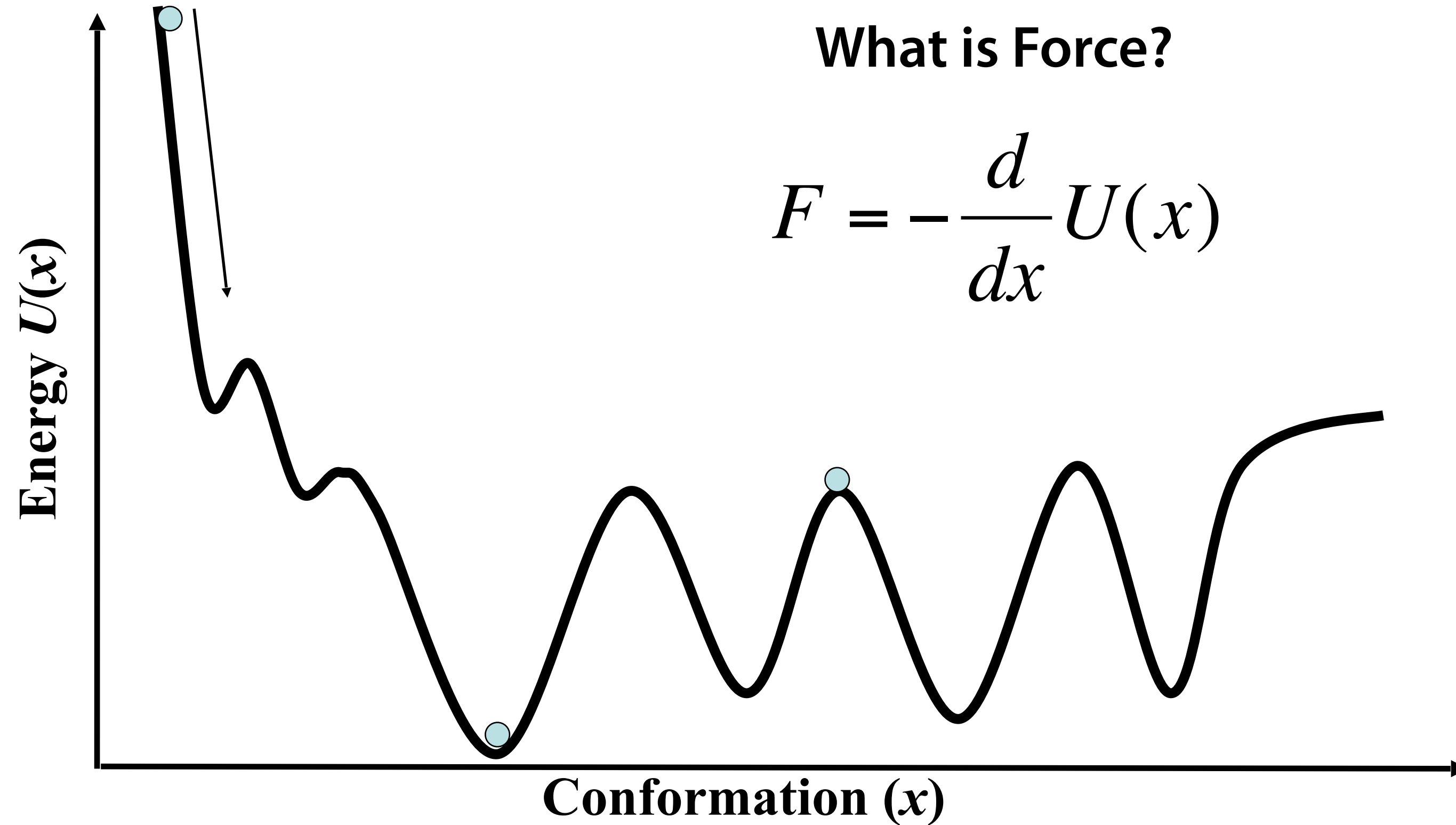
M.P. Allen and D.J. Tildesley. *Computer Simulation of Liquids*. Oxford University Press, New York, 1987.

D. Frenkel and B. Smit. *Understanding Molecular Simulations. From Algorithms to Applications*. Academic Press, San Diego, California, 1996.

A. R. Leach. *Molecular Modelling. Principles and Applications*. Addison Wesley Longman, Essex, England, 1996.

More at <http://www.biomath.nyu.edu/index/course/99/textbooks.html>

Potential Energy (hyper)Surface



Classical Molecular Dynamics *at 300 K*

Energy function: $U(\vec{r}_1, \vec{r}_2, \dots \vec{r}_N) = U(\vec{R})$

used to determine the force on each atom:

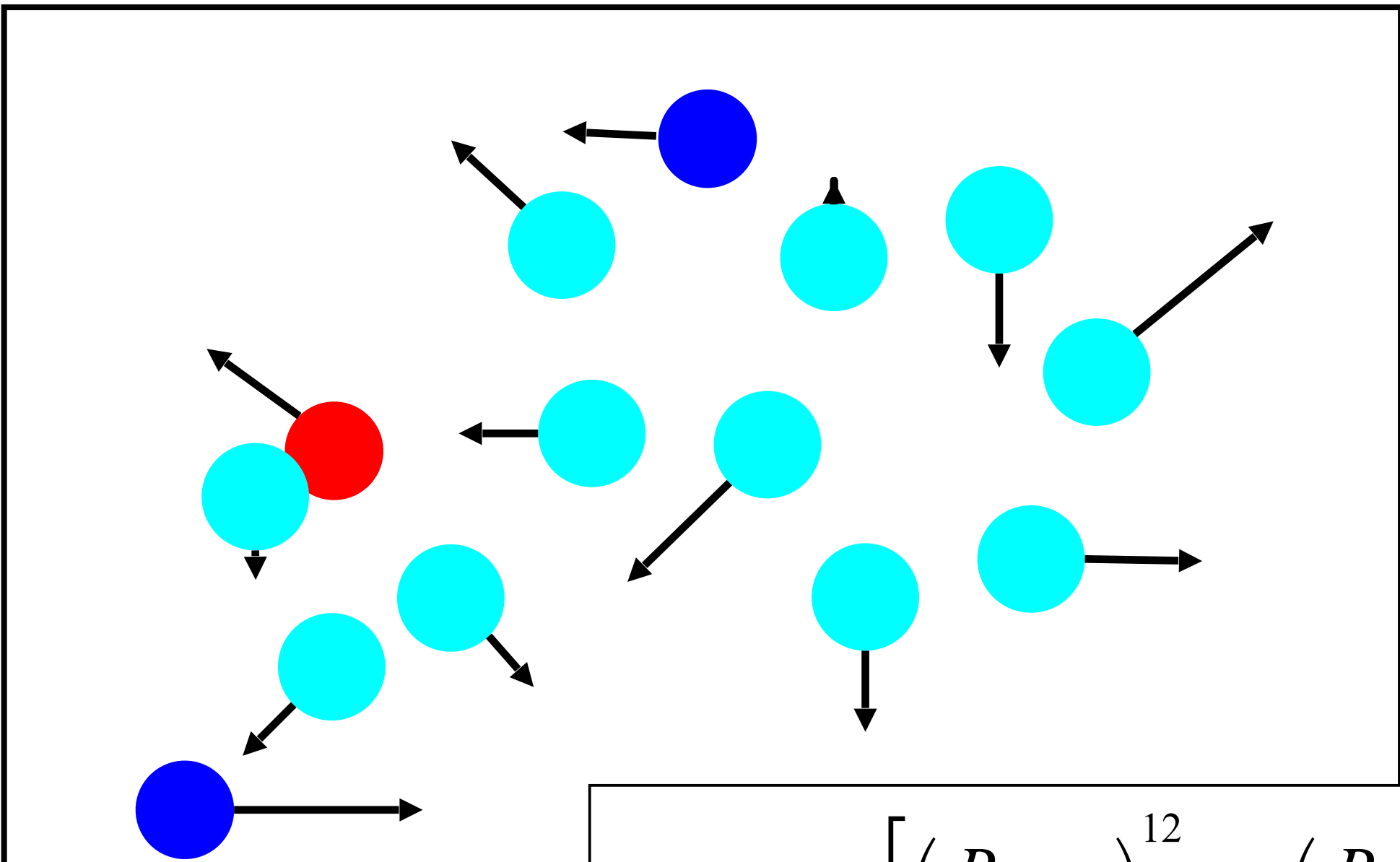
$$m_i \frac{d^2 \vec{r}_i}{dt^2} = \vec{F}_i = -\vec{\nabla} U(\vec{R})$$

yields a set of $3N$ coupled 2nd-order differential equations that can be propagated forward (or backward) in time.

Initial coordinates obtained from crystal structure,
velocities taken at random from Boltzmann distribution.

Classical Molecular Dynamics

tryptophan



$$U(r) = \frac{1}{4\pi\epsilon_0} \frac{q_i q_j}{r_{ij}}$$

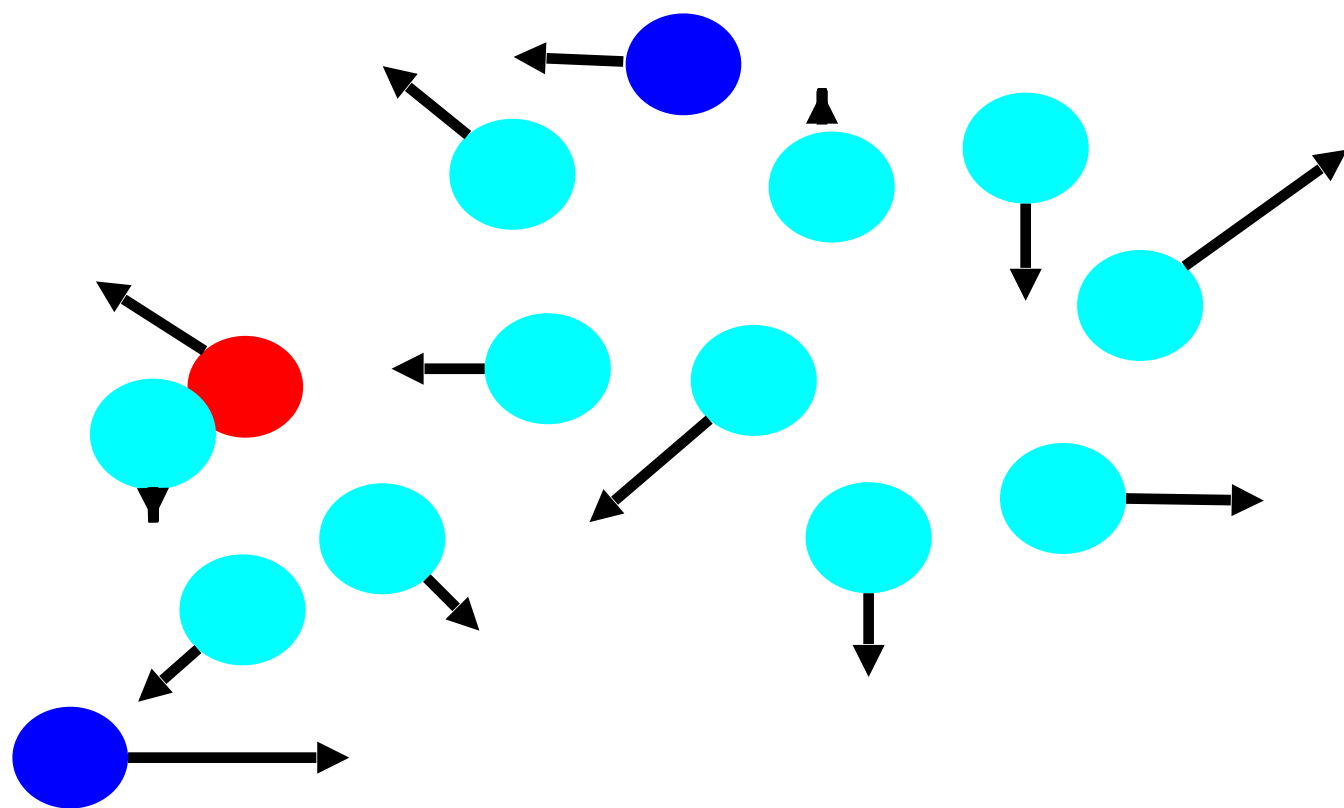
Coulomb interaction

$$U(r) = \epsilon_{ij} \left[\left(\frac{R_{\min,ij}}{r_{ij}} \right)^{12} - 2 \left(\frac{R_{\min,ij}}{r_{ij}} \right)^6 \right]$$

van der Waals interaction

Classical Molecular Dynamics

tryptophan

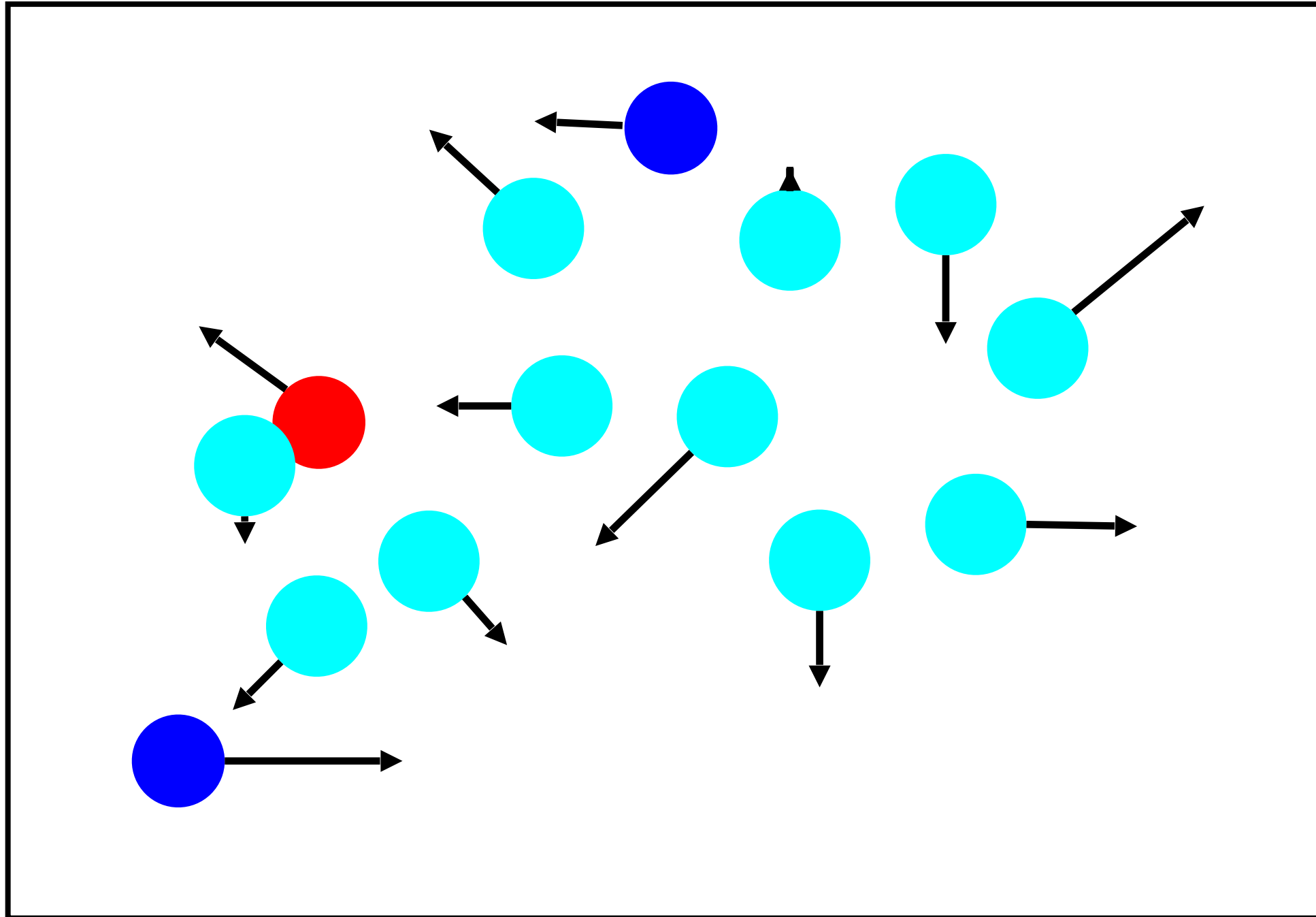


$$U(r) = \frac{1}{4\pi\epsilon_0} \frac{q_i q_j}{r_{ij}} + \epsilon_{ij} \left[\left(\frac{R_{\min,ij}}{r_{ij}} \right)^{12} - 2 \left(\frac{R_{\min,ij}}{r_{ij}} \right)^6 \right]$$

$$\mathbf{F}(\mathbf{r}) = \left(-\frac{1}{4\pi\epsilon_0} \frac{q_i q_j}{r_{ij}^2} - 12 \frac{\epsilon_{ij}}{|r_{ij}|} \left[\left(\frac{R_{\min,ij}}{r_{ij}} \right)^{12} - \left(\frac{R_{\min,ij}}{r_{ij}} \right)^6 \right] \right) \hat{\mathbf{r}}_{ij}$$

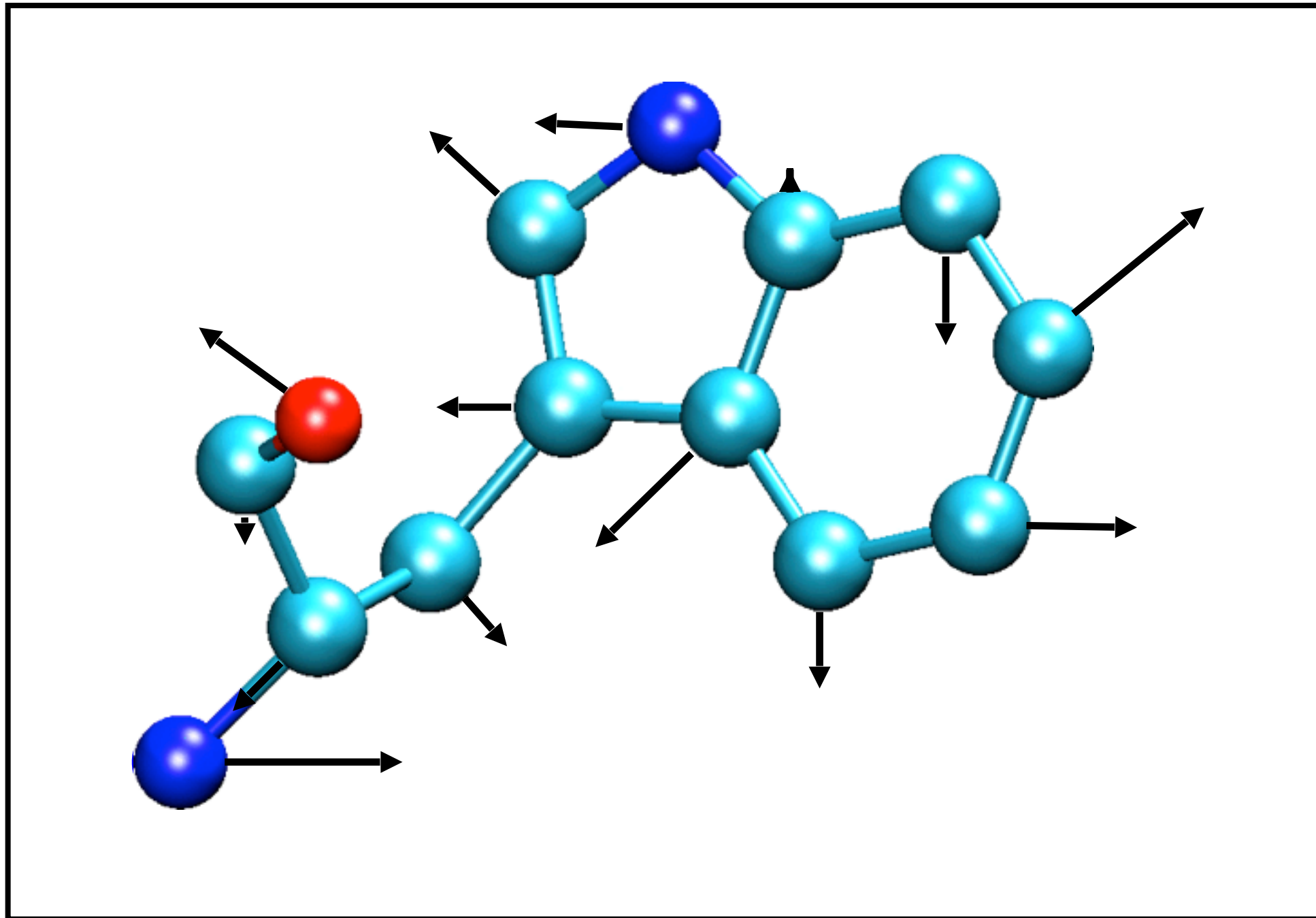
Classical Molecular Dynamics

tryptophan



Classical Molecular Dynamics

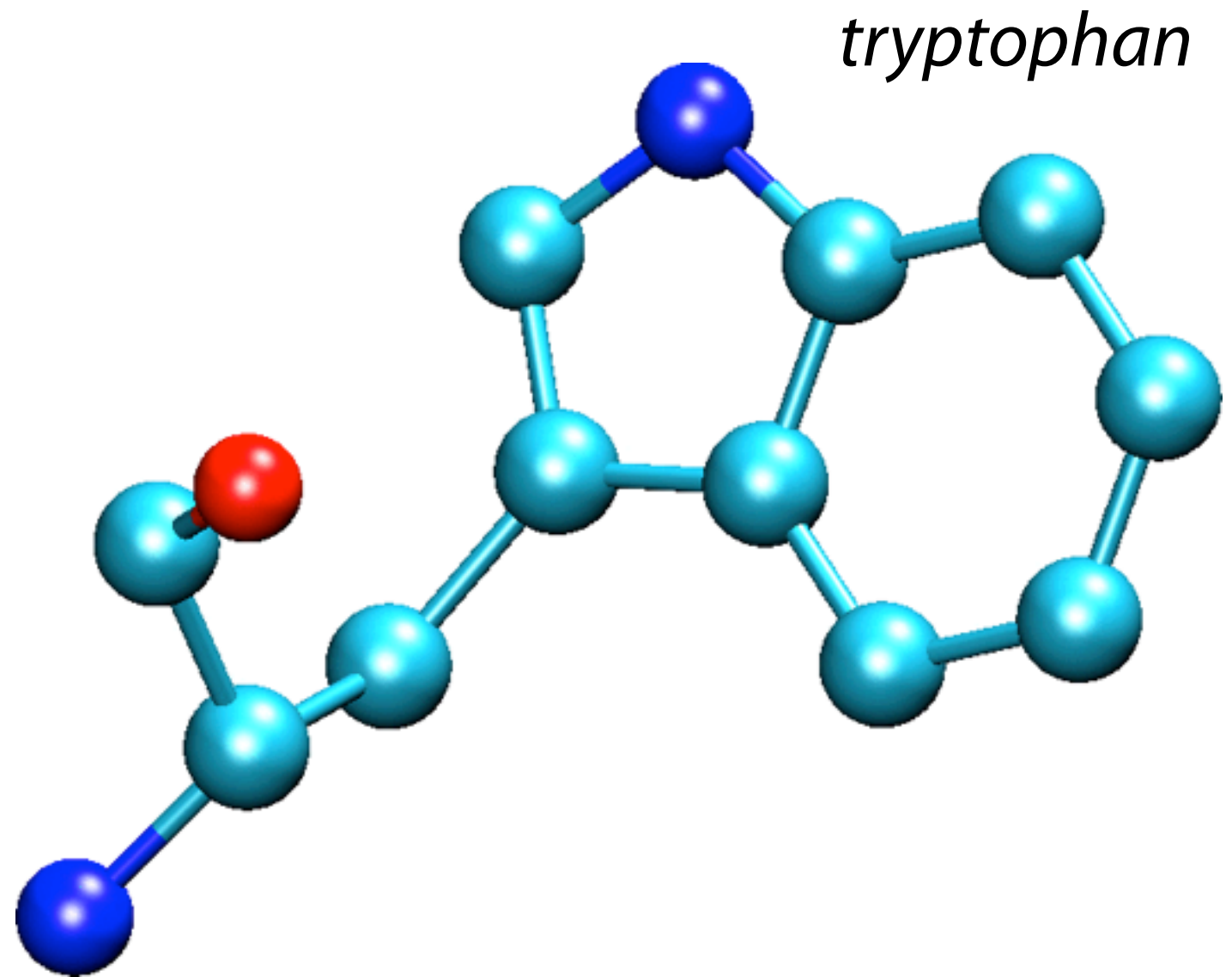
tryptophan



Bond definitions, atom types, atom names, parameters,

What is a Force Field?

In molecular dynamics a molecule is described as a series of **charged points** (atoms) **linked by springs** (bonds).

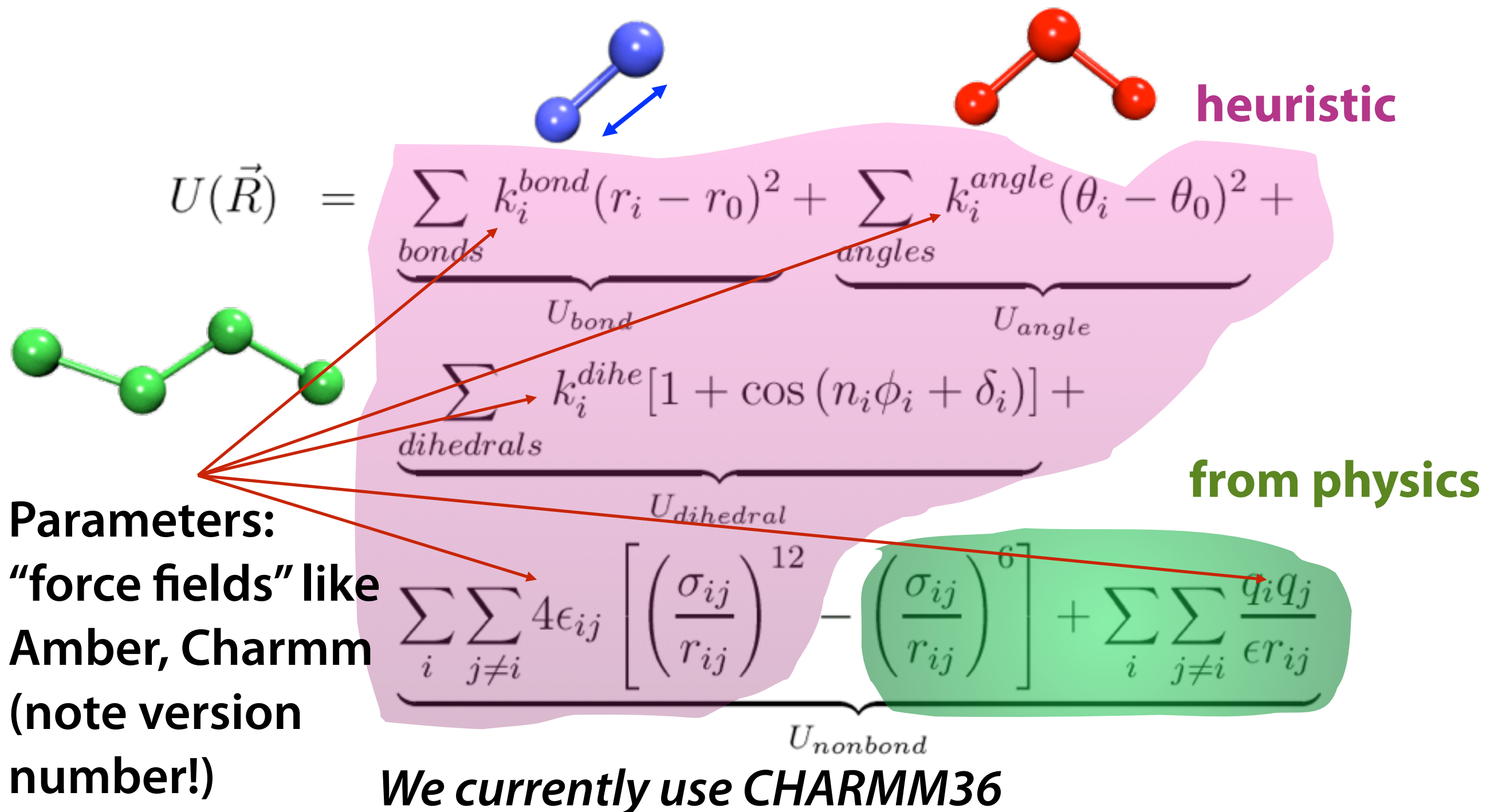


To describe the time evolution of bond lengths, bond angles and torsions, also the non-bonding van der Waals and electrostatic interactions between atoms, one uses a **force field**.

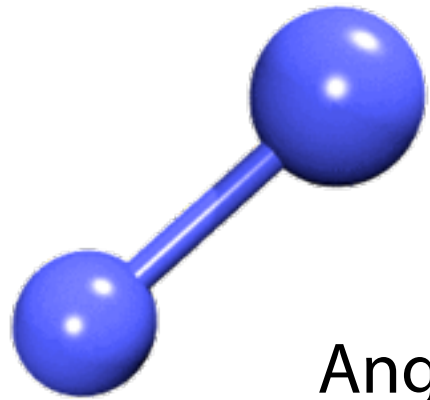
The **force field** is a collection of equations and associated constants designed to reproduce molecular geometry and selected properties of tested structures.

Potential Energy Function of Biopolymers

- Simple, fixed algebraic form for every type of interaction.
- Variable parameters depend on types of atoms involved.

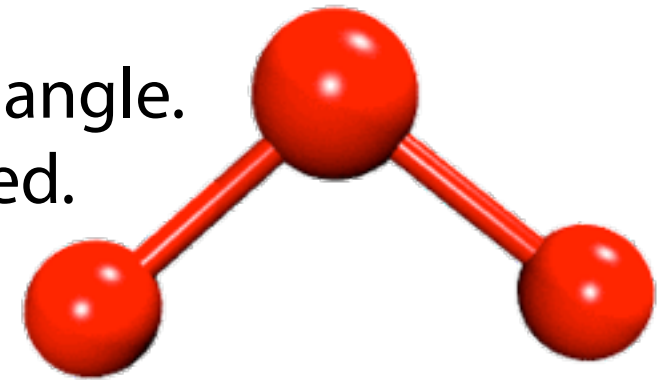


Potential Energy Function of Biopolymers

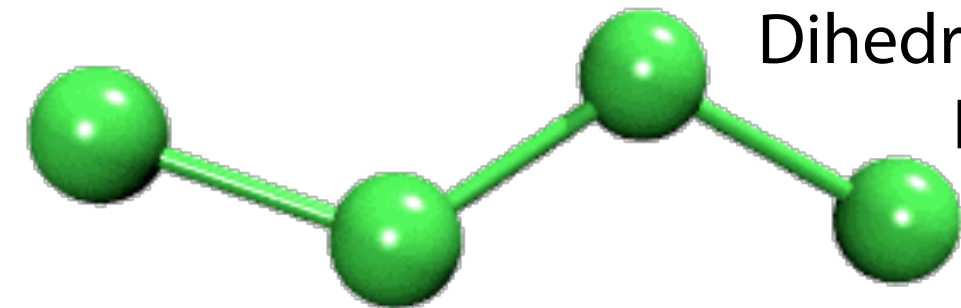


Bonds: Every pair of covalently bonded atoms is listed.

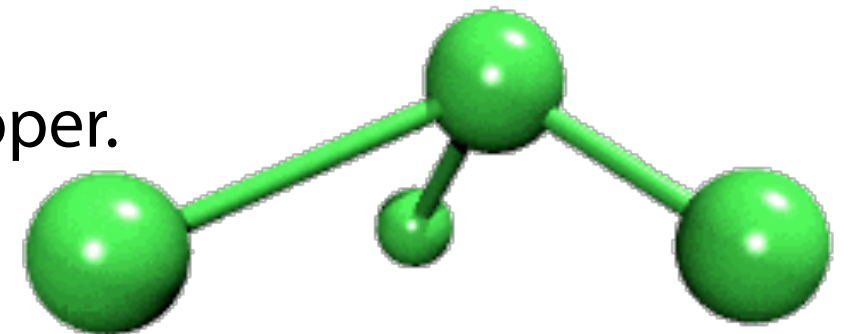
Angles: Two bonds that share a common atom form an angle.
Every such set of three atoms in the molecule is listed.



Dihedrals: Two angles that share a common bond form a dihedral.
Every such set of four atoms in the molecule is listed.



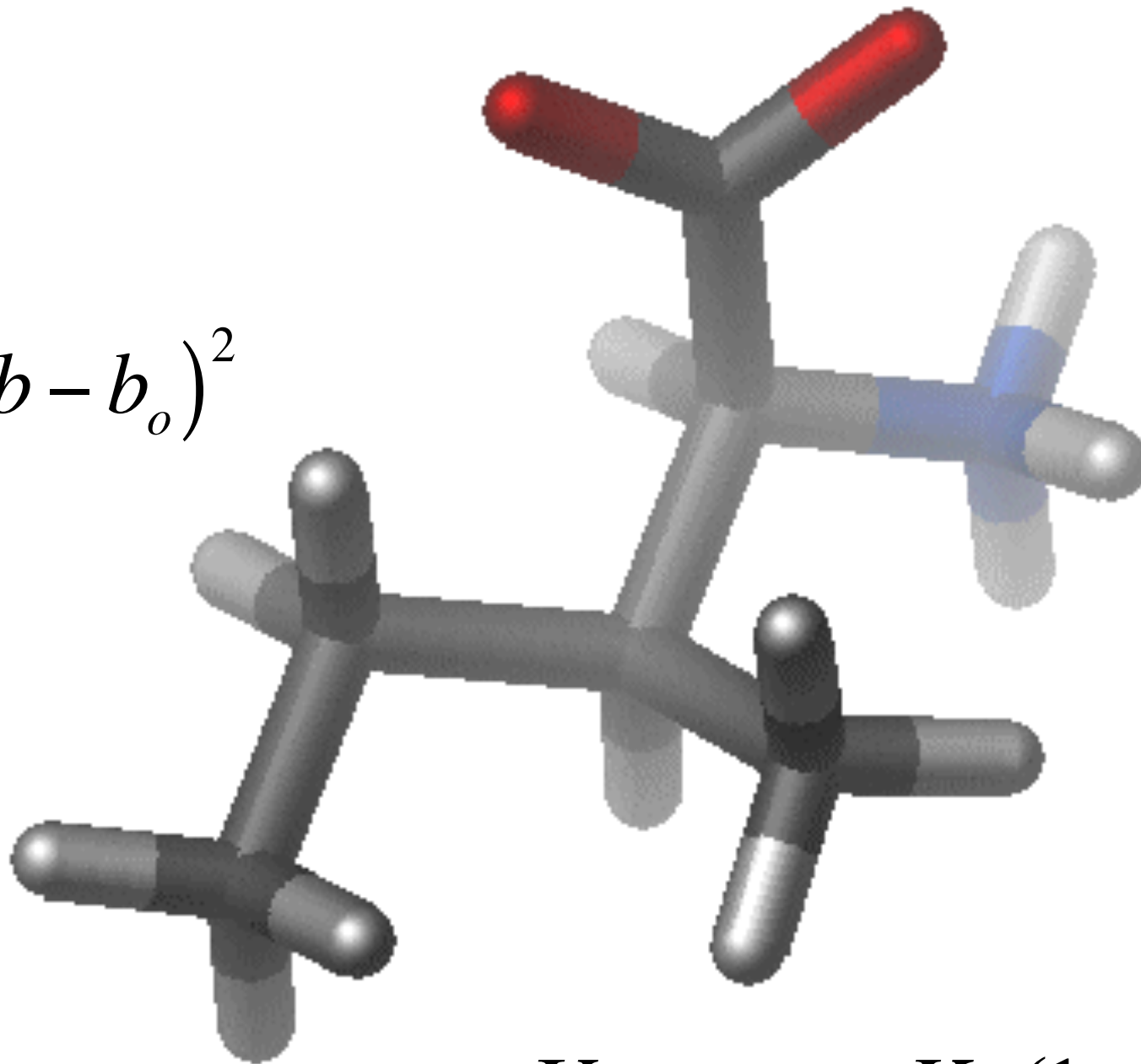
Impropers: Any *planar* group of four atoms forms an improper.
Specific sets of four atoms in the molecule are listed.



Interactions between bonded atoms

$$V_{angle} = K_{\theta} (\theta - \theta_o)^2$$

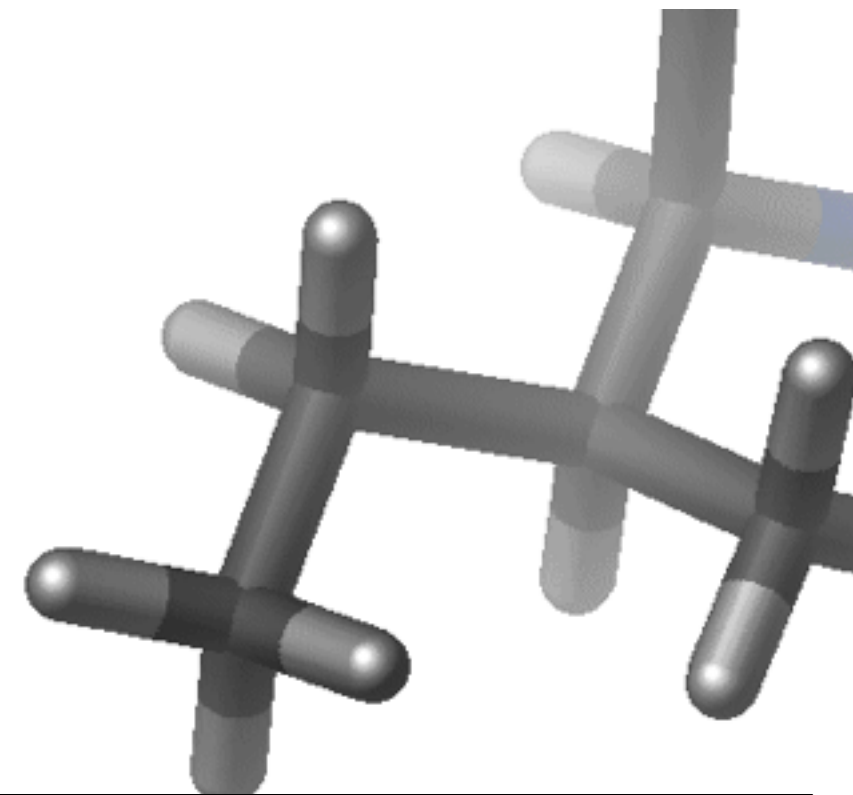
$$V_{bond} = K_b (b - b_o)^2$$



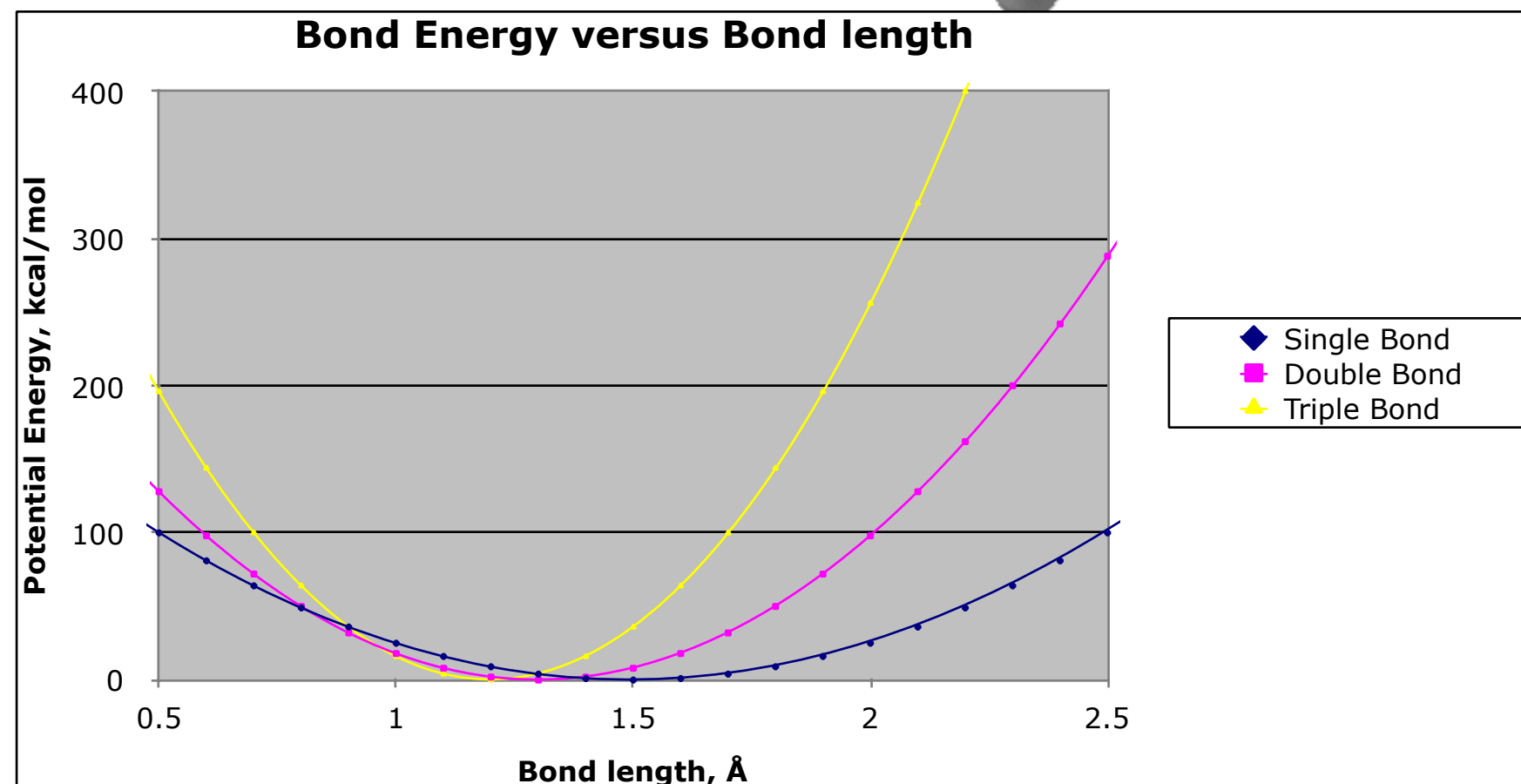
$$V_{dihedral} = K_{\phi} (1 + \cos(n\phi - \delta))$$

Bond potential

Chemical type	K_{bond}	b_0
C-C	100 kcal/mole/Å ²	1.5 Å
C=C	200 kcal/mole/Å ²	1.3 Å
C≡C	400 kcal/mole/Å ²	1.2 Å



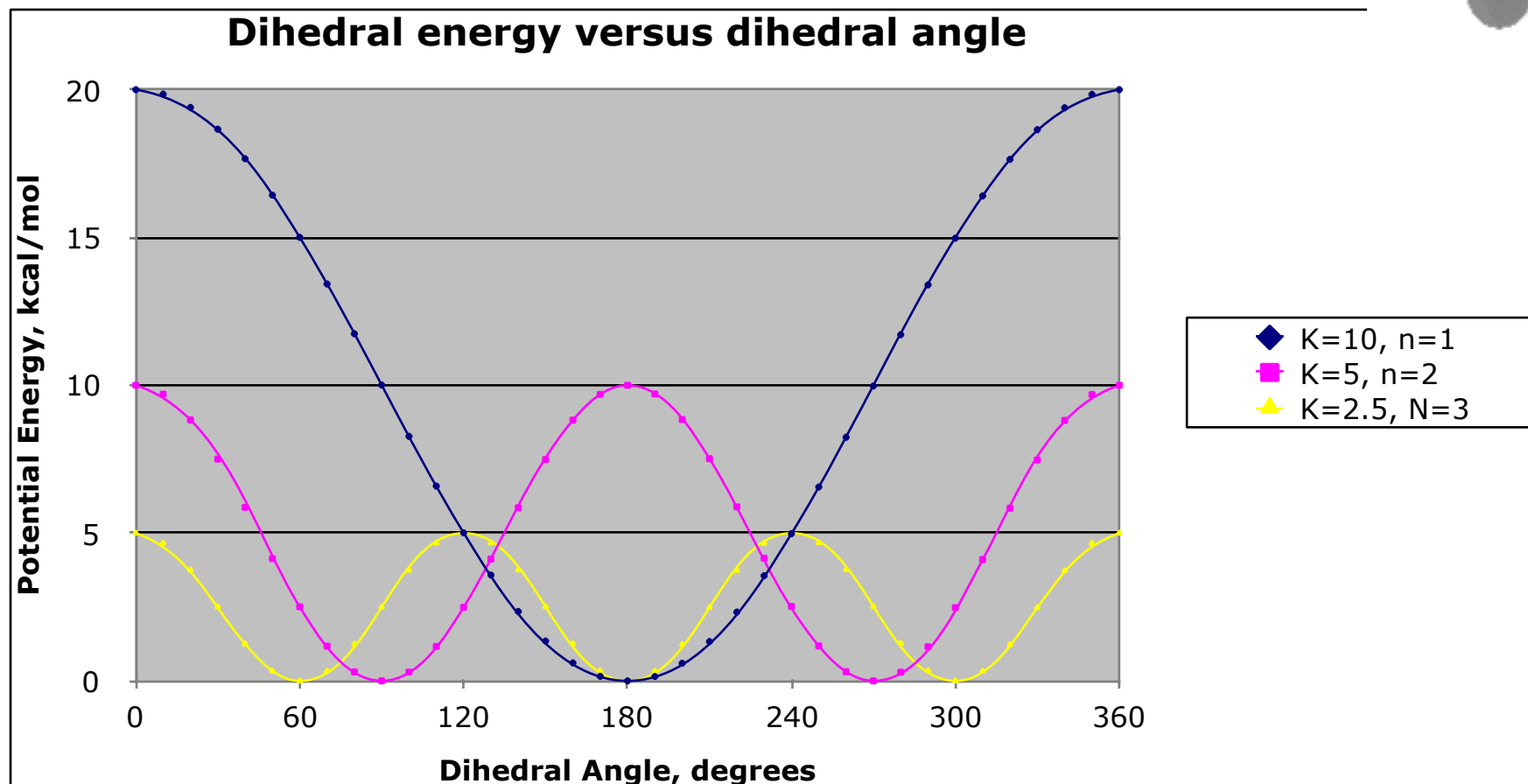
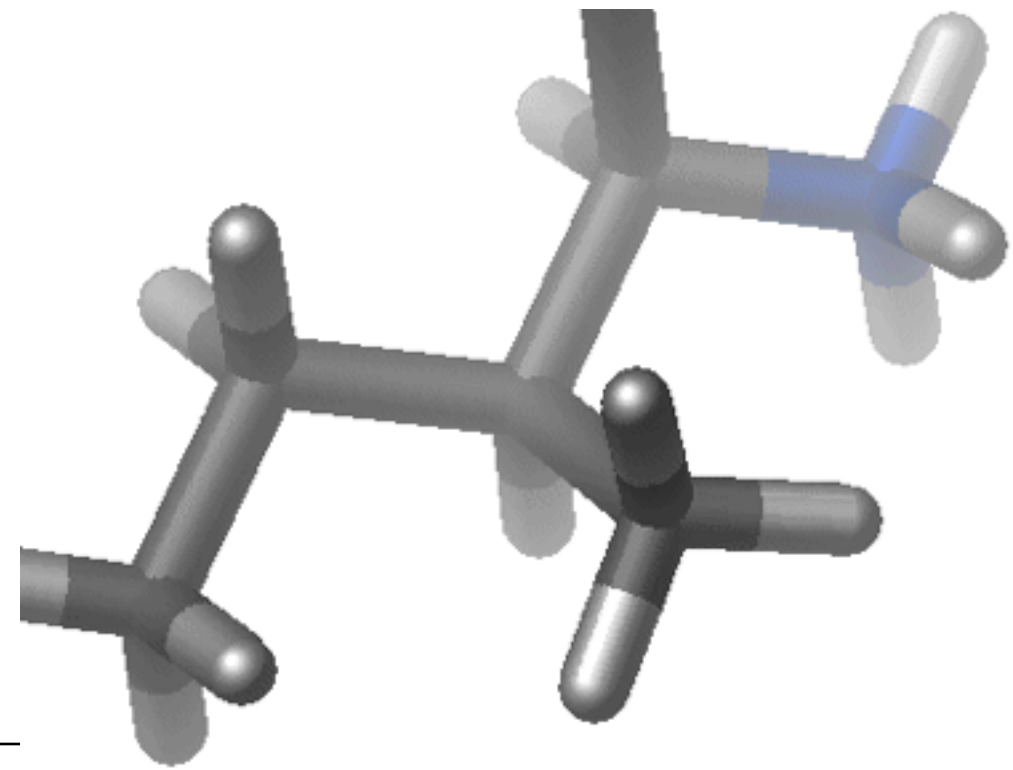
$$V_{\text{bond}} = K_b (b - b_0)^2$$



Bond-angle (3-body) and improper (4-body about a center) terms have similar quadratic forms, but with softer spring constants. The force constants can be obtained from vibrational analysis of the molecule (experimentally or theoretically).

Dihedral potential

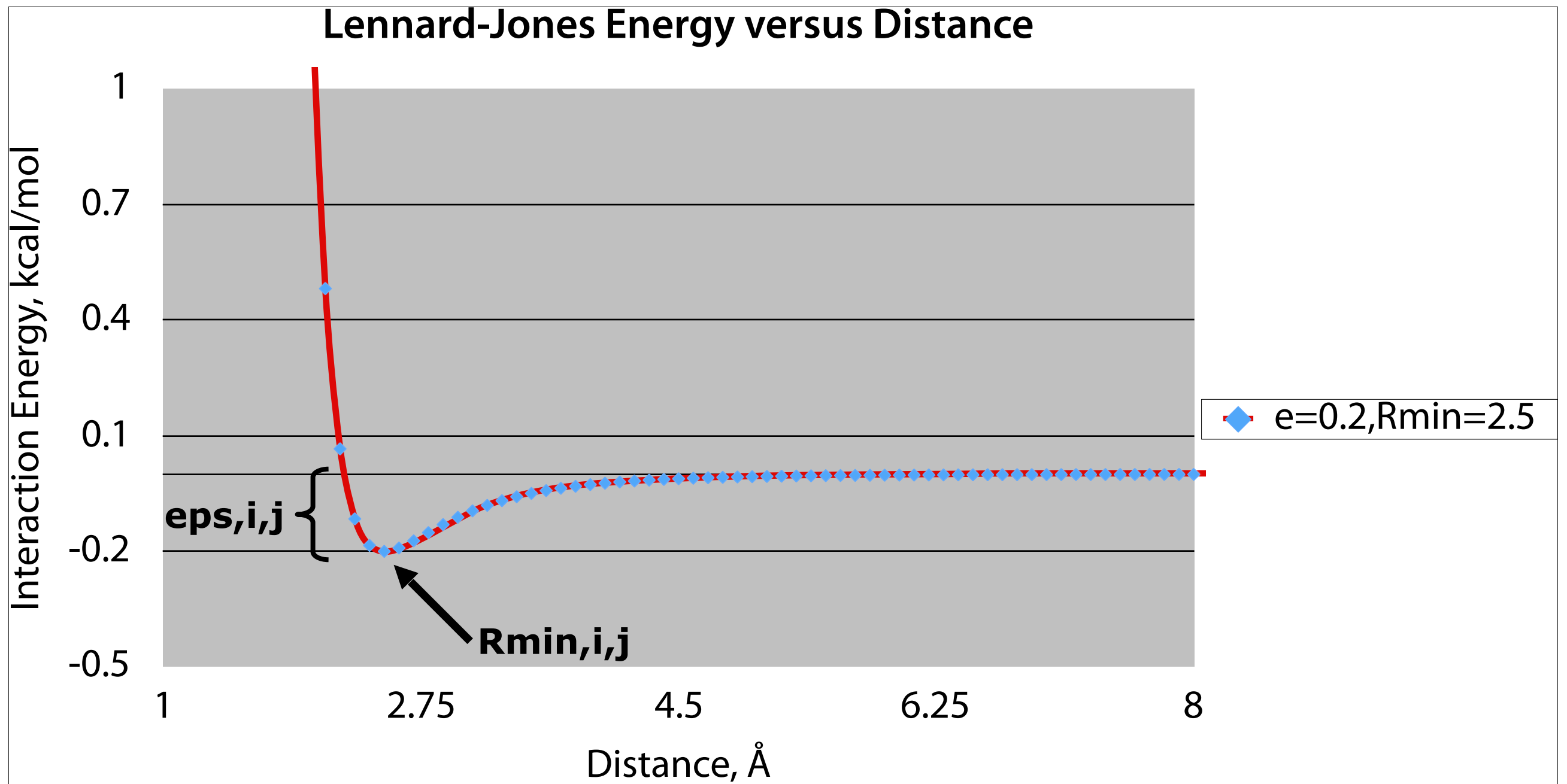
$$V_{dihedral} = K_{\phi} (1 + \cos(n\phi - \delta))$$



$\delta = 0$ for all three

dihedral-angle (4-body) terms come from symmetry in the electronic structure. Cross-term map (CMAP) terms in CHARMM force field are a refinement to this part of the potential

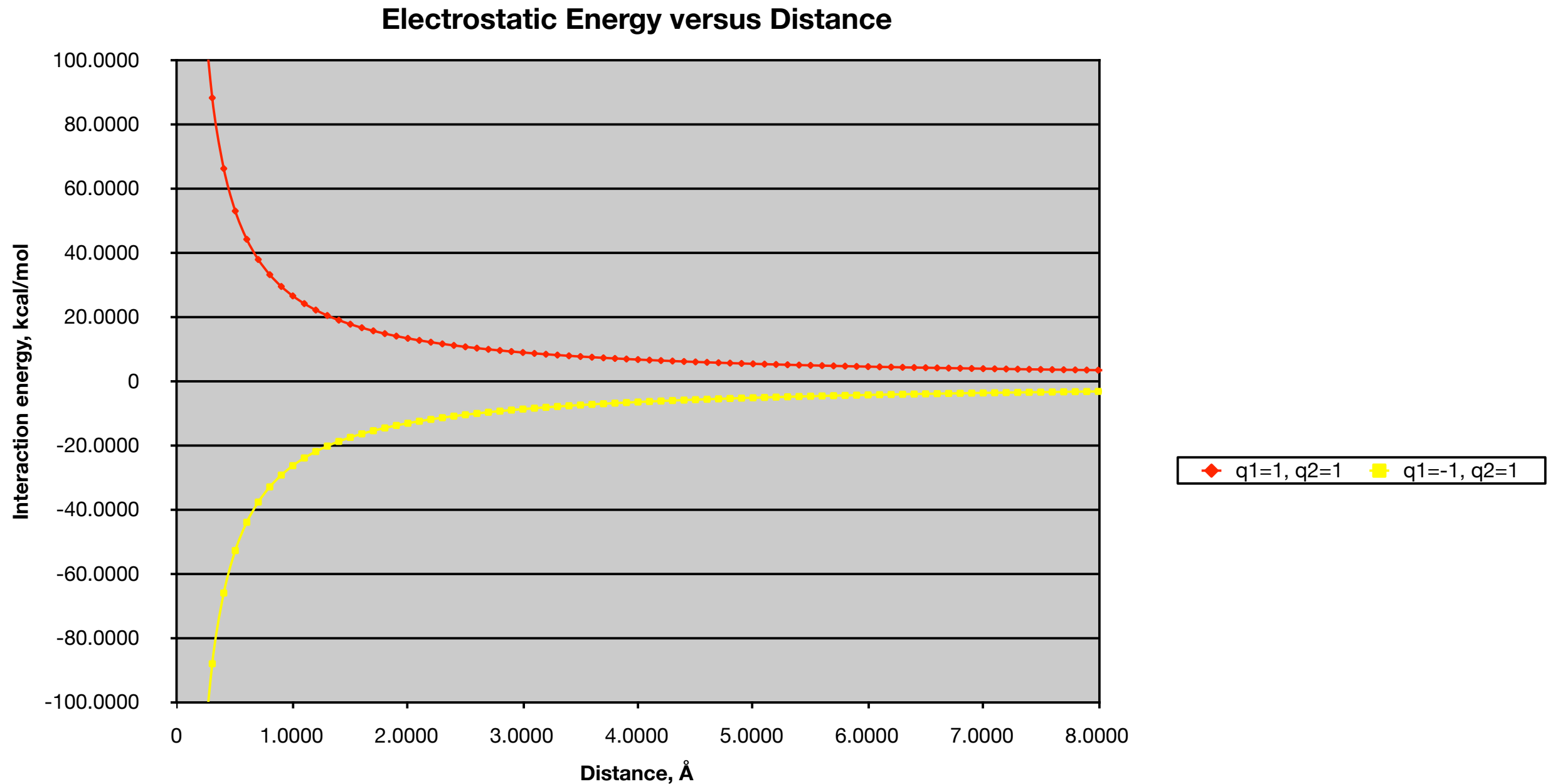
van der Waals potential



$$\epsilon_{ij} \left[\left(\frac{R_{min,ij}}{r_{ij}} \right)^{12} - 2 \left(\frac{R_{min,ij}}{r_{ij}} \right)^6 \right]$$

Short range; in CHARMM, cutoff at 12 Å
other FFs will cutoff at lower values (even
8 Å!) but beware the easy “**shortcut**”

Coulomb potential



Note that the effect is long range.

Equilibrium Properties of Proteins

Energies: kinetic and potential

?

temperature
dependence

$$\left\langle \sum_j \frac{1}{2} m_j v_j^2 \right\rangle = \frac{3}{2} N k_B T$$

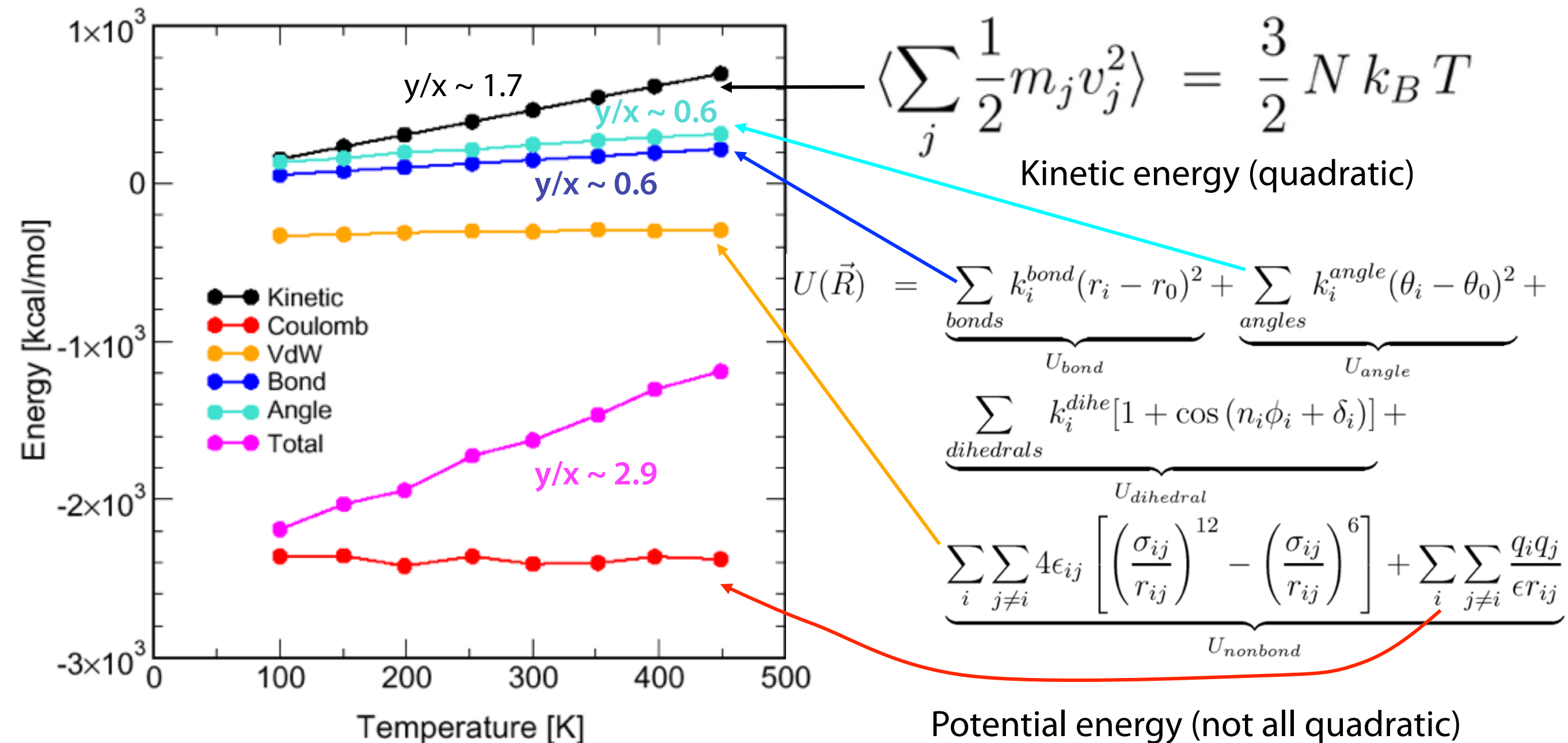
Kinetic energy (quadratic)

$$U(\vec{R}) = \underbrace{\sum_{bonds} k_i^{bond} (r_i - r_0)^2}_{U_{bond}} + \underbrace{\sum_{angles} k_i^{angle} (\theta_i - \theta_0)^2}_{U_{angle}} + \underbrace{\sum_{dihedrals} k_i^{dihe} [1 + \cos(n_i \phi_i + \delta_i)]}_{U_{dihedral}} + \underbrace{\sum_i \sum_{j \neq i} 4\epsilon_{ij} \left[\left(\frac{\sigma_{ij}}{r_{ij}} \right)^{12} - \left(\frac{\sigma_{ij}}{r_{ij}} \right)^6 \right] + \sum_i \sum_{j \neq i} \frac{q_i q_j}{\epsilon r_{ij}}}_{U_{nonbond}}$$

Potential energy (not all quadratic)

Equilibrium Properties of Proteins

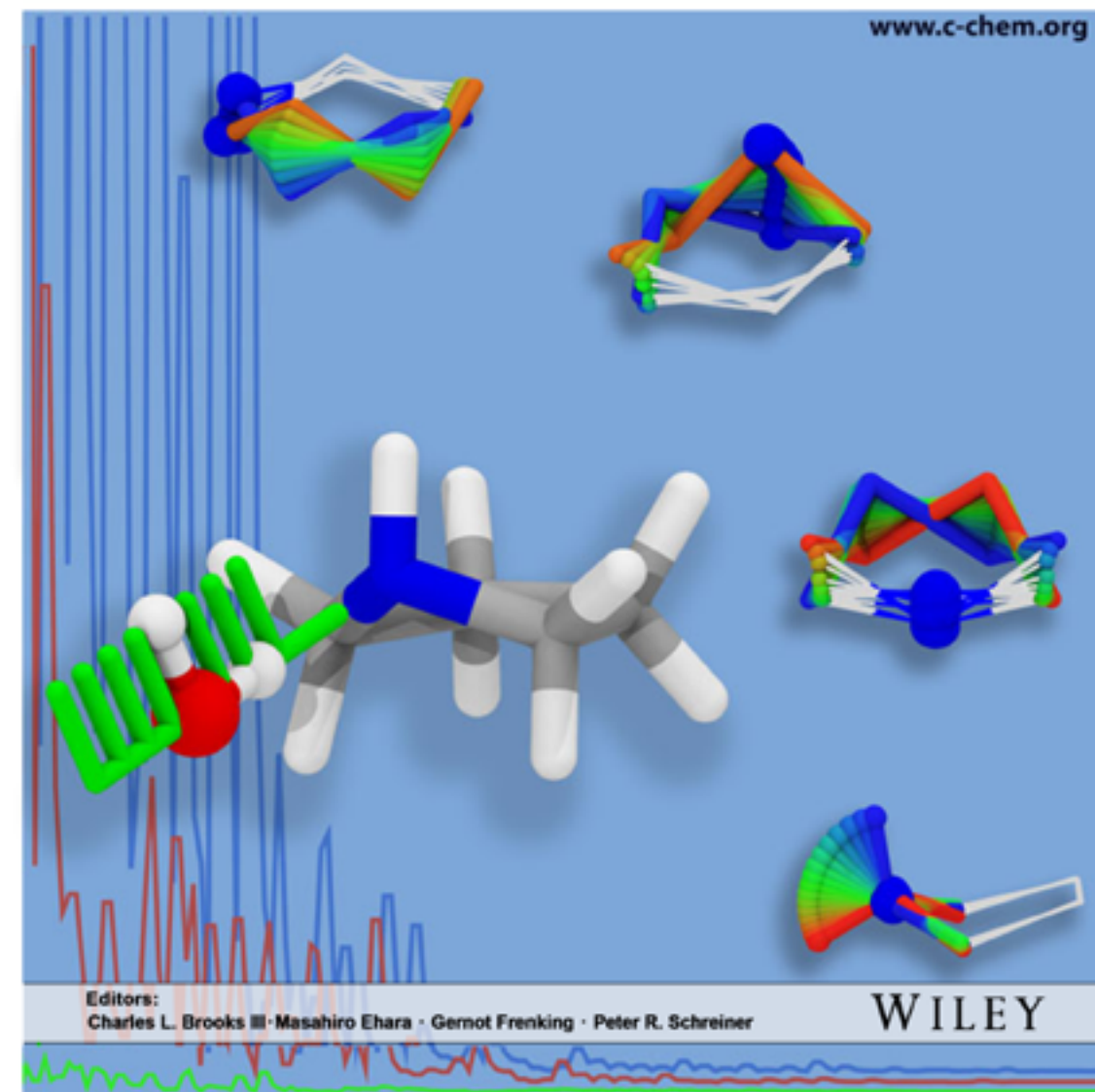
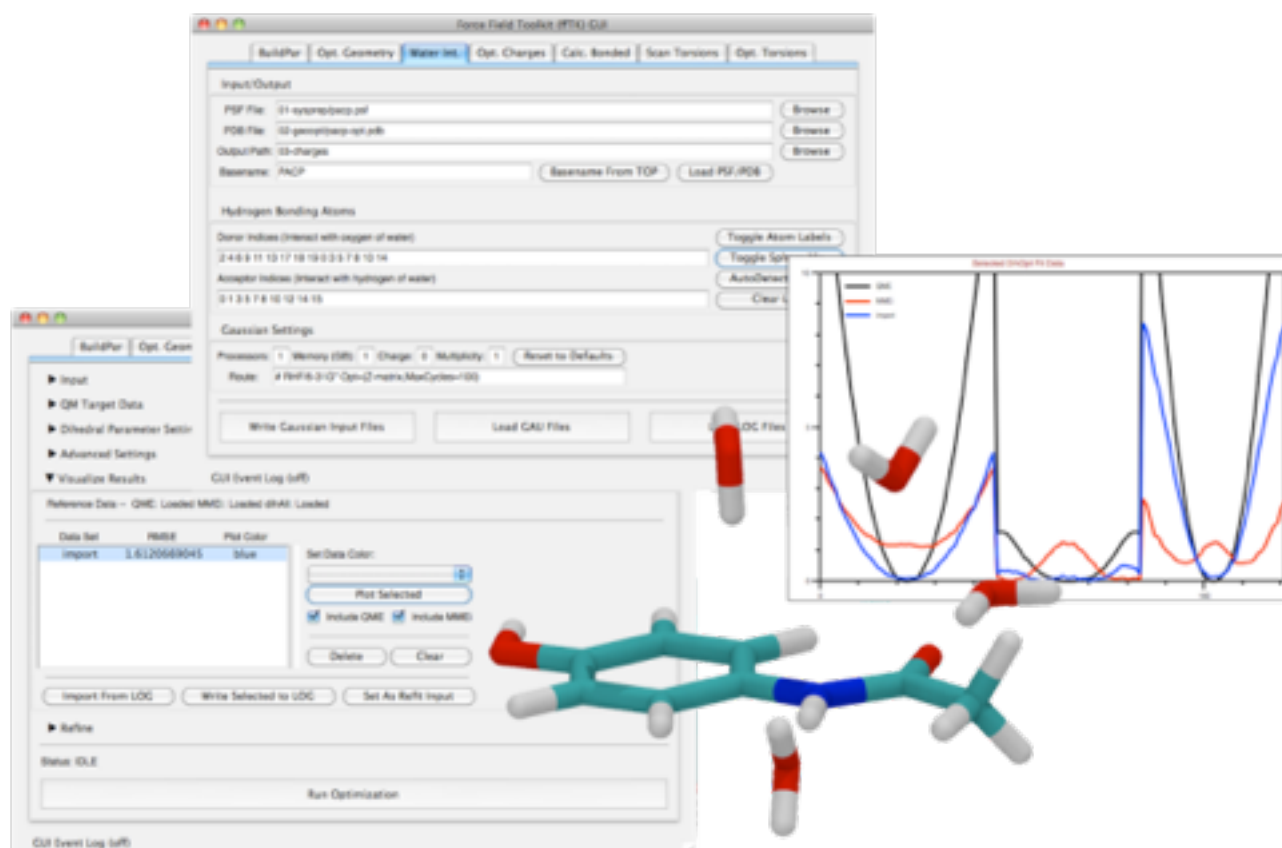
Energies: kinetic and potential



Force field toolkit (FFTK)

$$U(\vec{R}) = \underbrace{\sum_{\text{bonds}} k_i^{\text{bond}} (r_i - r_0)^2}_{U_{\text{bond}}} + \underbrace{\sum_{\text{angles}} k_i^{\text{angle}} (\theta_i - \theta_0)^2}_{U_{\text{angle}}} + \underbrace{\sum_{\text{dihedrals}} k_i^{\text{dihe}} [1 + \cos(n_i \phi_i + \delta_i)]}_{U_{\text{dihedral}}} + \underbrace{\sum_i \sum_{j \neq i} 4\epsilon_{ij} \left[\left(\frac{\sigma_{ij}}{r_{ij}} \right)^{12} - \left(\frac{\sigma_{ij}}{r_{ij}} \right)^6 \right] + \sum_i \sum_{j \neq i} \frac{q_i q_j}{\epsilon r_{ij}}}_{U_{\text{nonbond}}}$$

FFTK aids in the development of parameters in the **MD potential function** for novel molecules, ligands

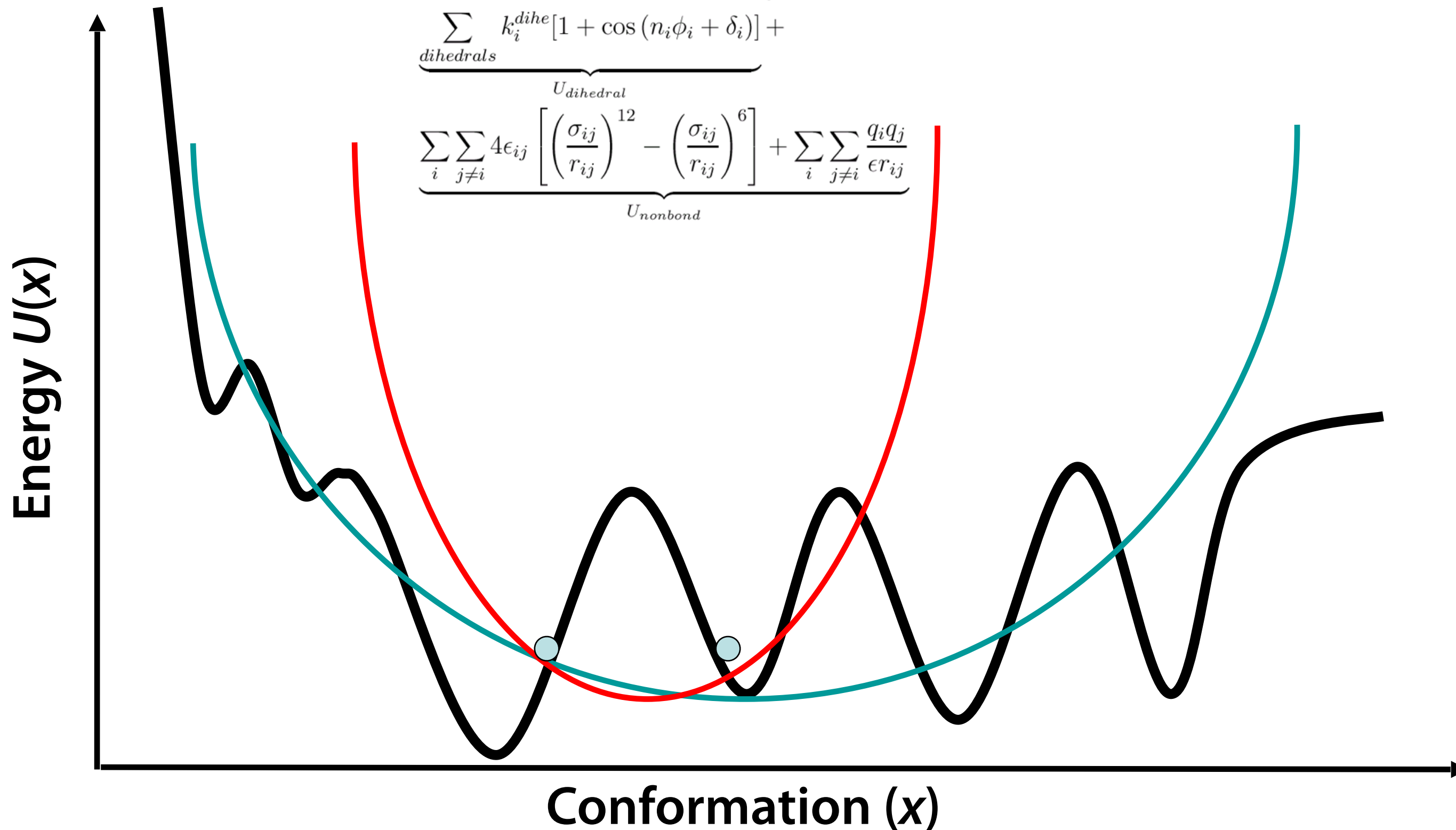


Mayne, Saam, Schulten, Tajkhorshid, and Gumbart. **Rapid parameterization of small molecules using the force field toolkit.** (2013) *J. Comp. Chem.* 34:2757-2770.

<http://www.ks.uiuc.edu/Research/vmd/plugins/fftk/>

Potential Energy (hyper)Surface

$$\begin{aligned}
 U(\vec{R}) = & \underbrace{\sum_{bonds} k_i^{bond} (r_i - r_0)^2}_{U_{bond}} + \underbrace{\sum_{angles} k_i^{angle} (\theta_i - \theta_0)^2}_{U_{angle}} + \\
 & \underbrace{\sum_{dihedrals} k_i^{dihe} [1 + \cos(n_i \phi_i + \delta_i)]}_{U_{dihedral}} + \\
 & \underbrace{\sum_i \sum_{j \neq i} 4\epsilon_{ij} \left[\left(\frac{\sigma_{ij}}{r_{ij}} \right)^{12} - \left(\frac{\sigma_{ij}}{r_{ij}} \right)^6 \right]}_{U_{nonbond}} + \sum_i \sum_{j \neq i} \frac{q_i q_j}{\epsilon r_{ij}}
 \end{aligned}$$



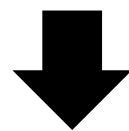
Classical Molecular Dynamics *discretization in time for computing*

$$m_i \frac{d^2 \vec{r}_i}{dt^2} = \vec{F}_i = -\vec{\nabla} U(\vec{R})$$

Use positions and accelerations at time t and the positions from time $t-\delta t$ to calculate new positions at time $t+\delta t$.

$$\begin{aligned} \mathbf{r}(t + \delta t) &\approx \mathbf{r}(t) + \mathbf{v}(t)\delta t + \frac{1}{2}\mathbf{a}(t)\delta t^2 \\ &+ \\ \mathbf{r}(t - \delta t) &\approx \mathbf{r}(t) - \mathbf{v}(t)\delta t + \frac{1}{2}\mathbf{a}(t)\delta t^2 \end{aligned}$$

“Verlet algorithm”

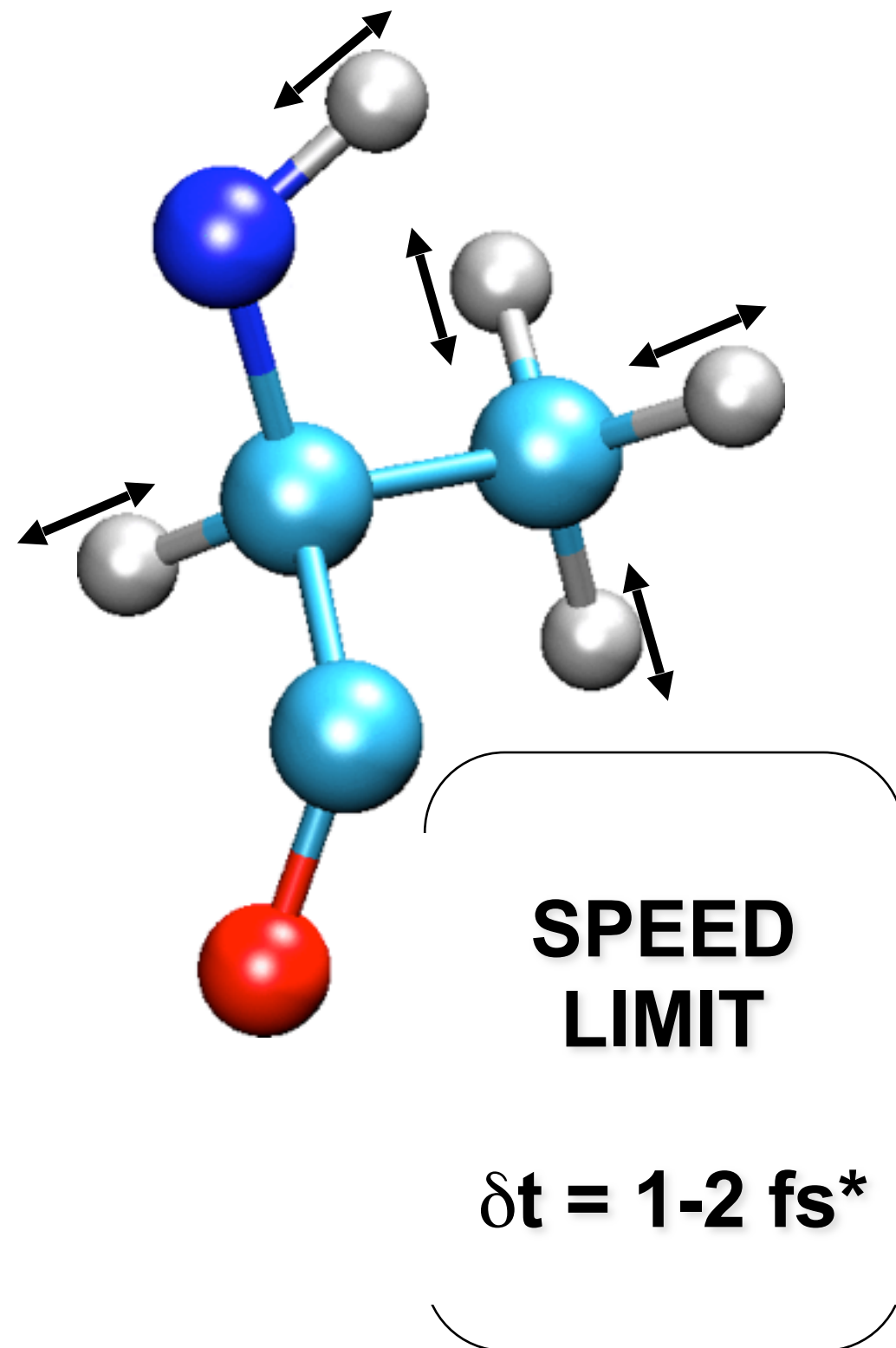


$$-\vec{\nabla} U(\vec{R}) / m_i$$

$$\mathbf{r}(t + \delta t) \approx 2\mathbf{r}(t) - \mathbf{r}(t - \delta t) + \mathbf{a}(t)\delta t^2$$

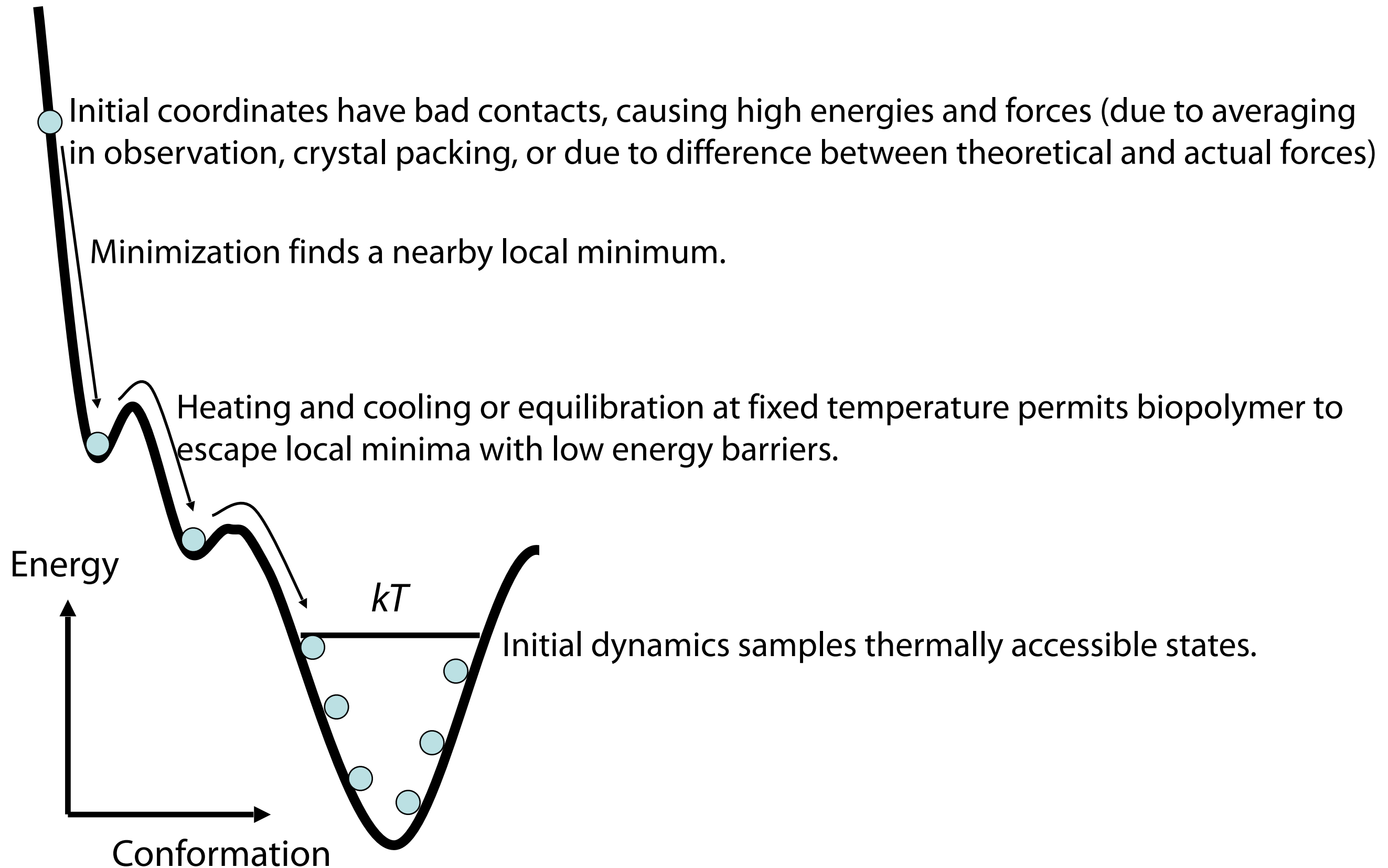
The most serious bottleneck

		<i>steps</i>
Protein folding (typical)	s	10^{15}
Rotation of buried sidechains	ms	10^{12}
Allosteric transitions		
Protein folding (fastest)	μ s	10^9
Hinge bending	ns	10^6
Rotation of surface sidechains	ps	10^3
bond stretching		
Molecular dynamics timestep	fs	10^0

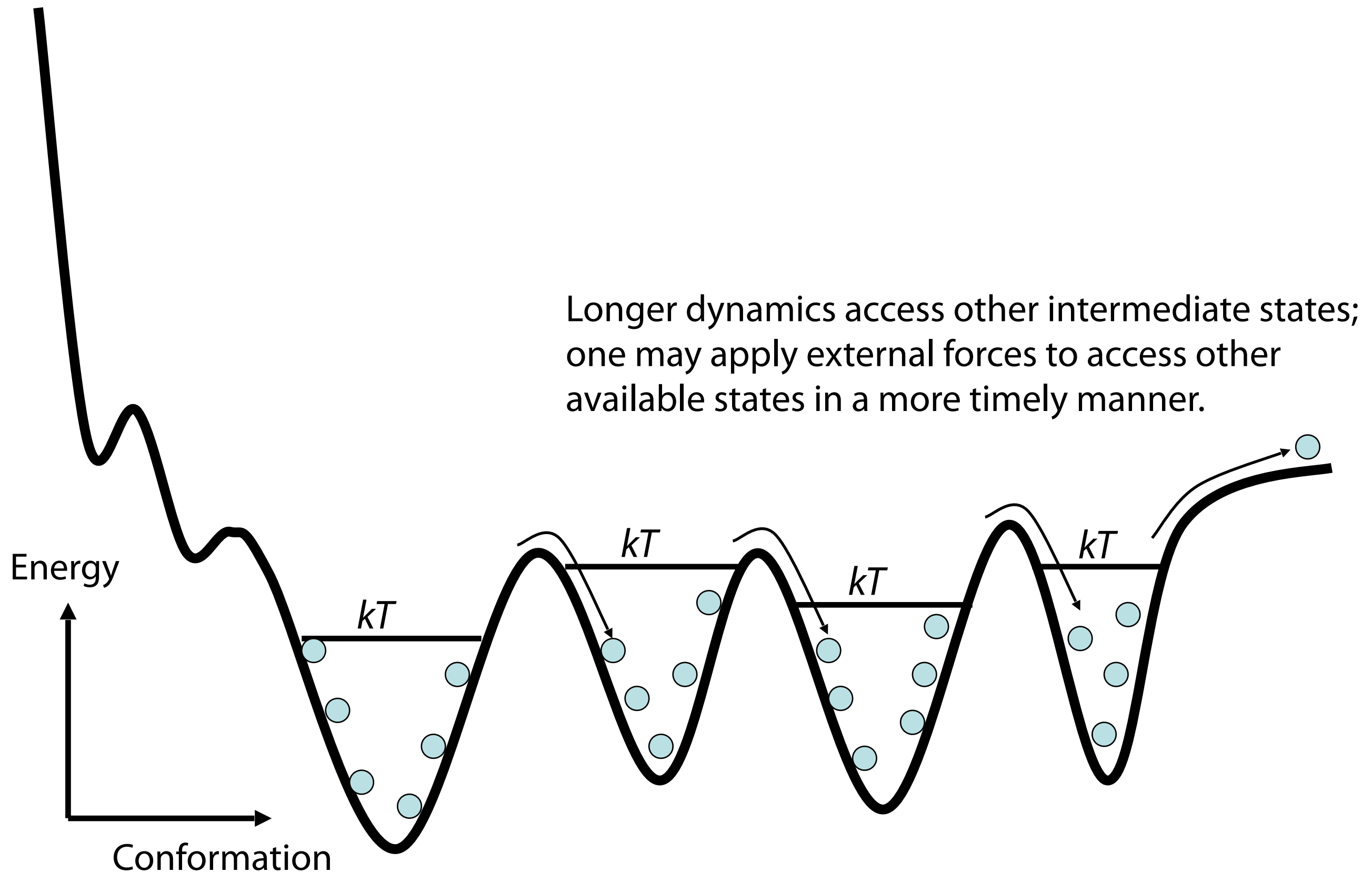


*experimental 4-fs time steps with "hydrogen mass repartitioning" exists, but not in NAMD yet

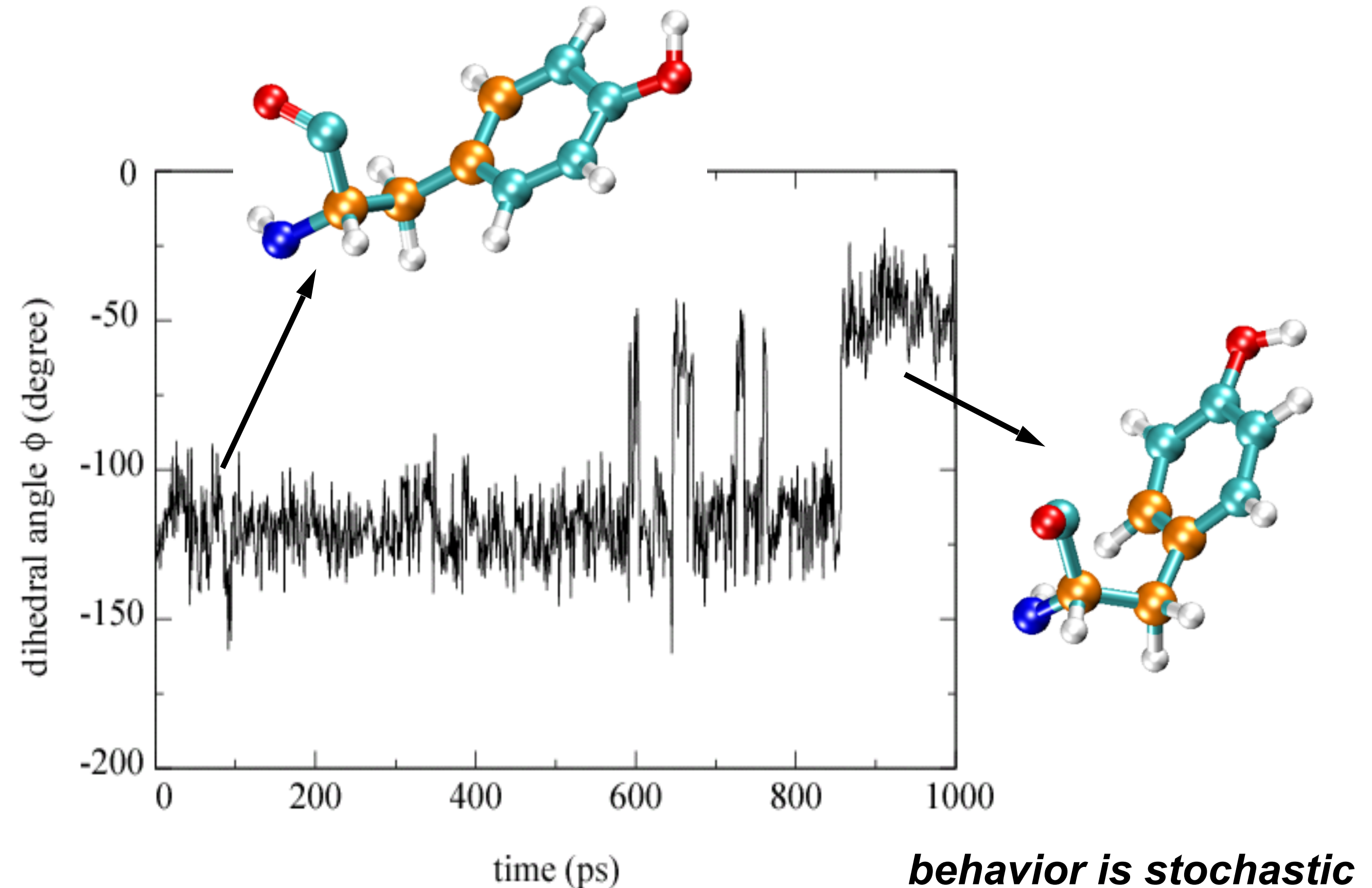
Molecular Dynamics to Sample Energy Landscape



Molecular Dynamics to Sample Energy Landscape



Patience is required to observe molecular events



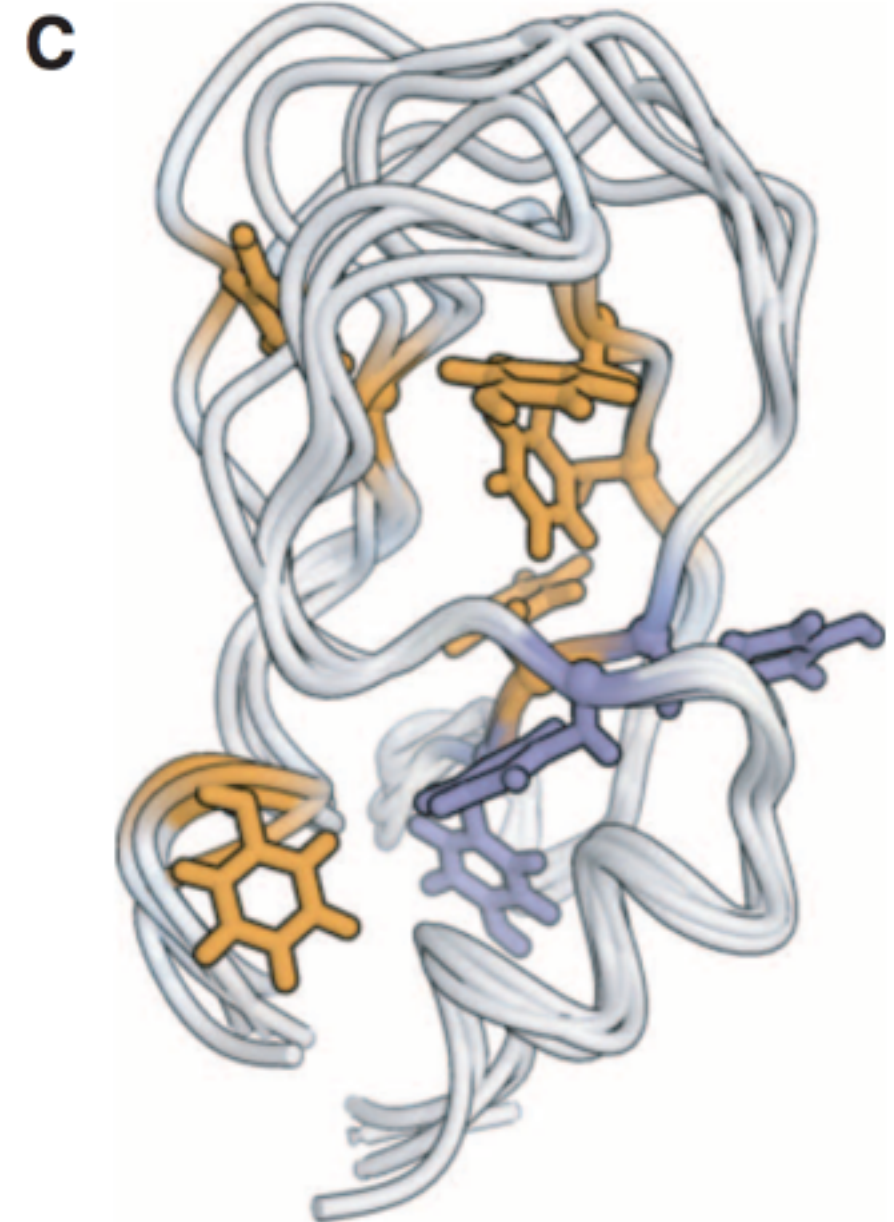
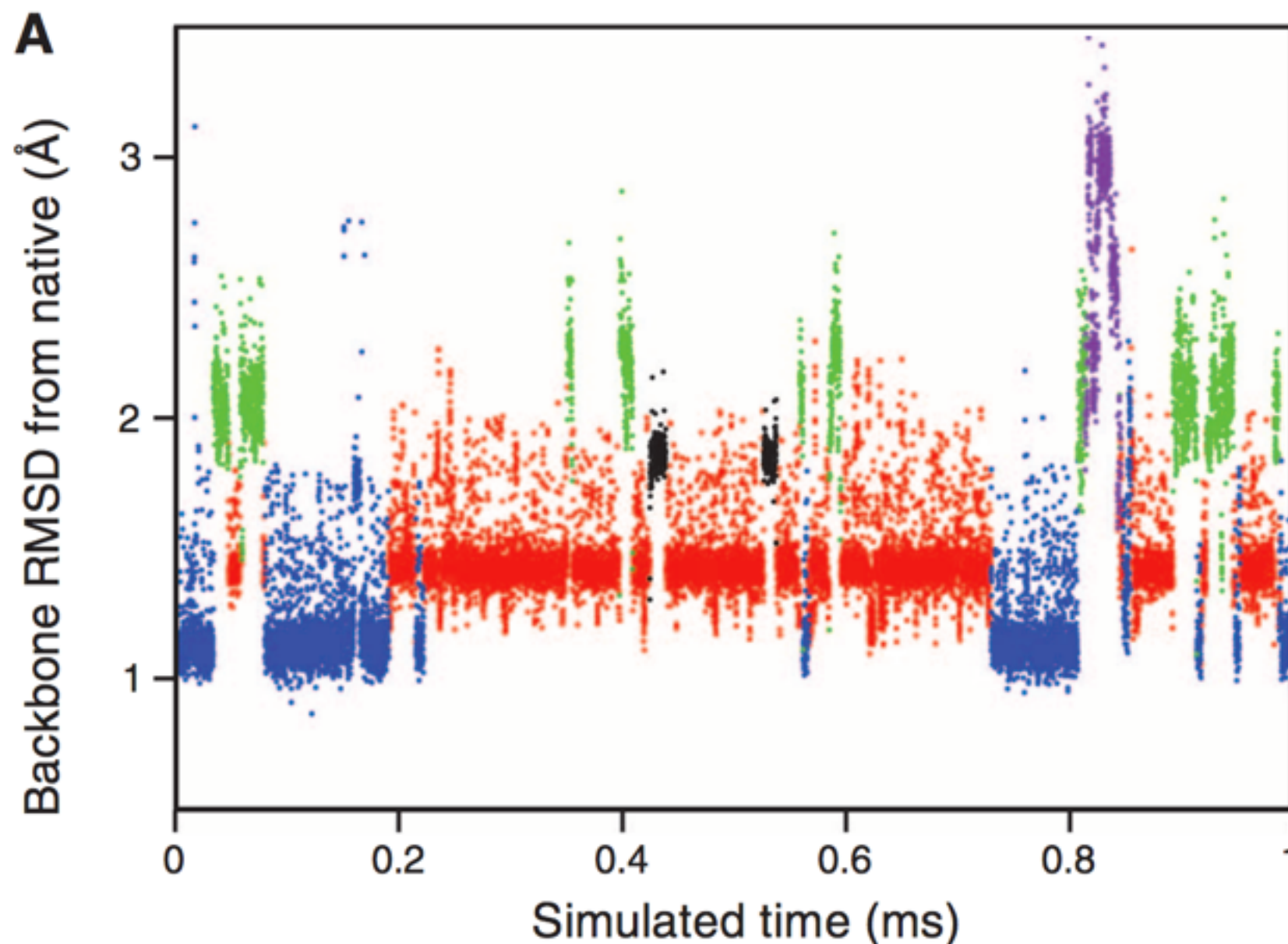
Patience is required to observe molecular events

RESEARCH ARTICLE

Atomic-Level Characterization of the Structural Dynamics of Proteins

David E. Shaw^{1,2,*}, Paul Maragakis^{1,†}, Kresten Lindorff-Larsen^{1,†}, Stefano Piana^{1,†}, Ron O. Dror¹, Michael P. Eastwood¹, Joseph A. Bank¹, John M. Jumper¹, John K. Salmon¹, Yibing Shan¹, Willy Wriggers¹

And there will always be a scale of dynamics you don't observe!



Molecular Dynamics Ensembles

thinking in terms of statistical mechanics

Constant energy, constant number of particles (NE; *only if no periodic boundary conditions*)

Constant energy, constant volume (NVE)

Constant temperature, constant volume (NVT)

Constant temperature, constant pressure (NPT)

Choose the ensemble that best fits your system and start the simulations - for most biomolecular systems, we choose **NPT**

Temperature and Pressure control methods

-to simulate **NVT** (canonical) ensemble, need to duplicate the effect of a large thermal bath around the system

Dynamics governed by the Langevin equation;
gives the **correct ensemble**

$$M\dot{v} = F(x) - \underbrace{\gamma M v}_{\text{damping term}} - \underbrace{\sqrt{2\gamma k_B T M} R(t)}_{\text{random term}}$$
$$\langle \vec{R}(t) \cdot \vec{R}(t') \rangle = 6k_B T \gamma \delta(t - t')$$

-Damping (1/ps) should be enough to maintain temperature without significantly perturbing dynamics (**5 is too much, 1 is probably okay**)

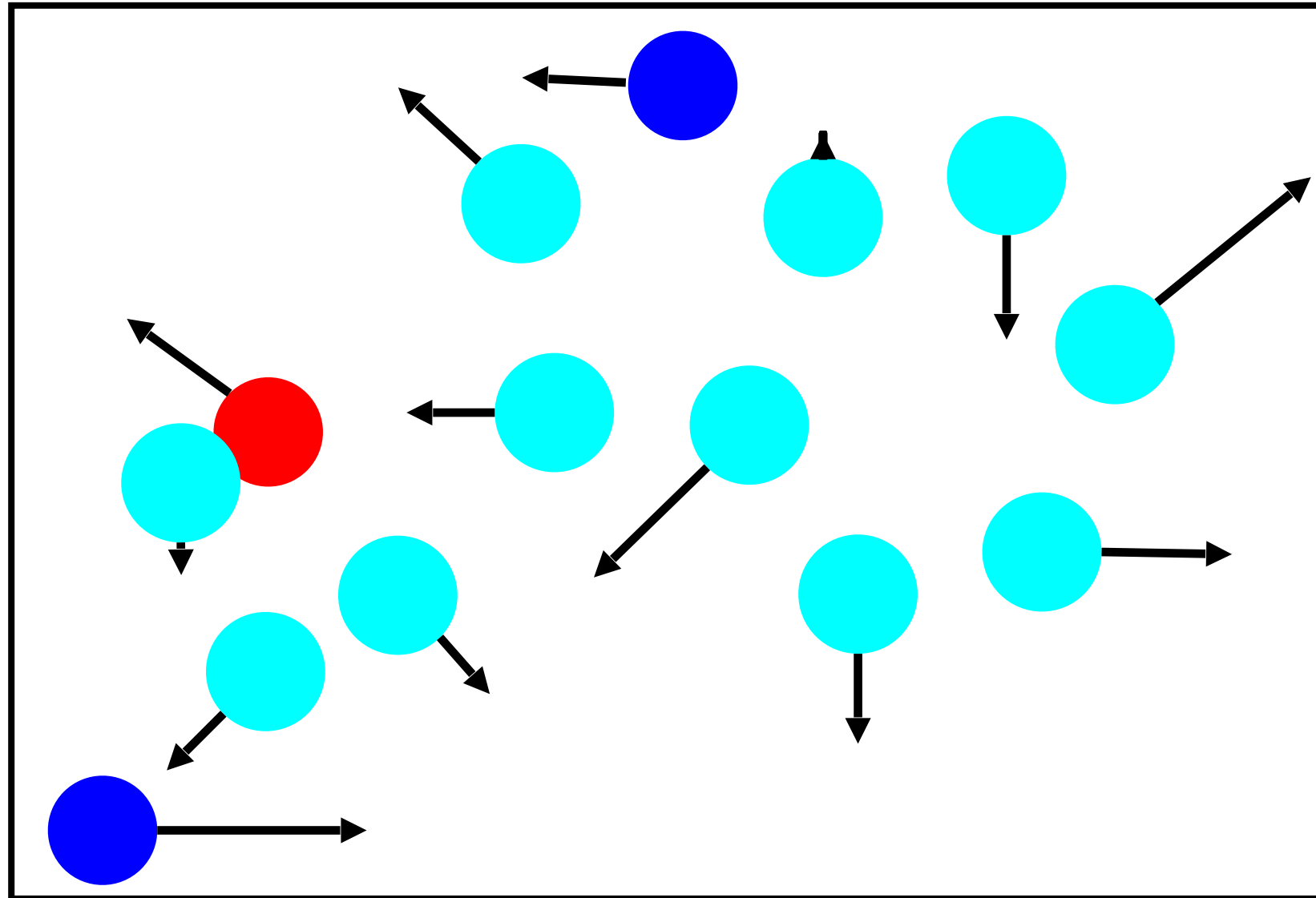
**For non-equilibrium simulations the Lowe-Andersen thermostat, which conserves momentum and does not suppress flows, may be preferred*

NPT (isobaric-isothermal) ensemble adds additional variables to control temperature, pressure which are ultimately integrated out to generate the correct distribution

$$P = \rho k_B T + \frac{1}{dV} \left\langle \sum_{i < j} \mathbf{f}(\mathbf{r}_{ij}) \cdot \mathbf{r}_{ij} \right\rangle$$

In the simulation, pressure is calculated from the virial expansion

Boundary Conditions?

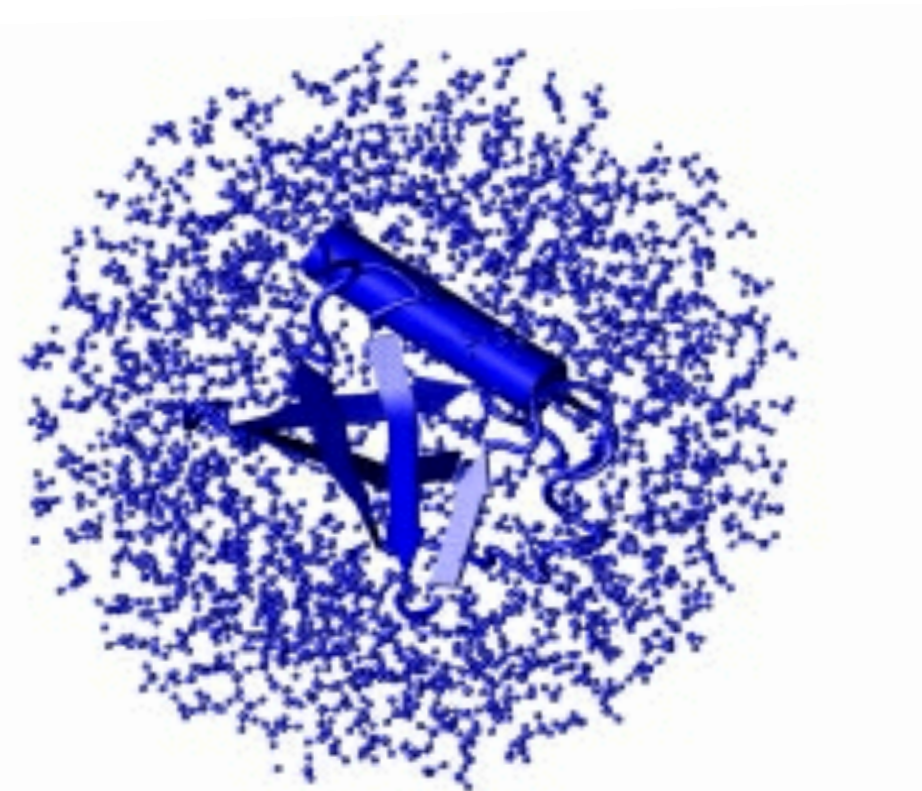
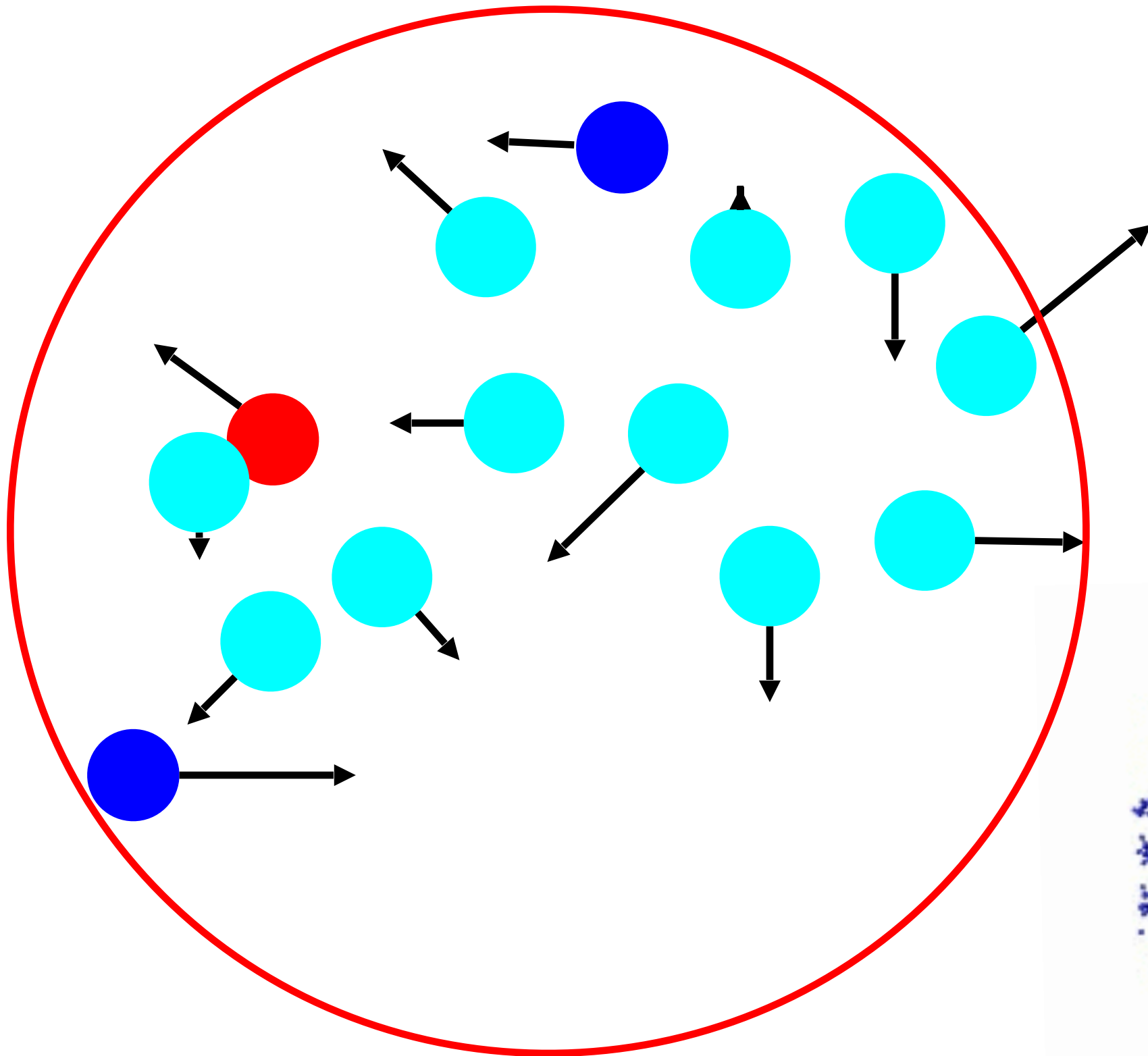


What happens if you put water under vacuum!?

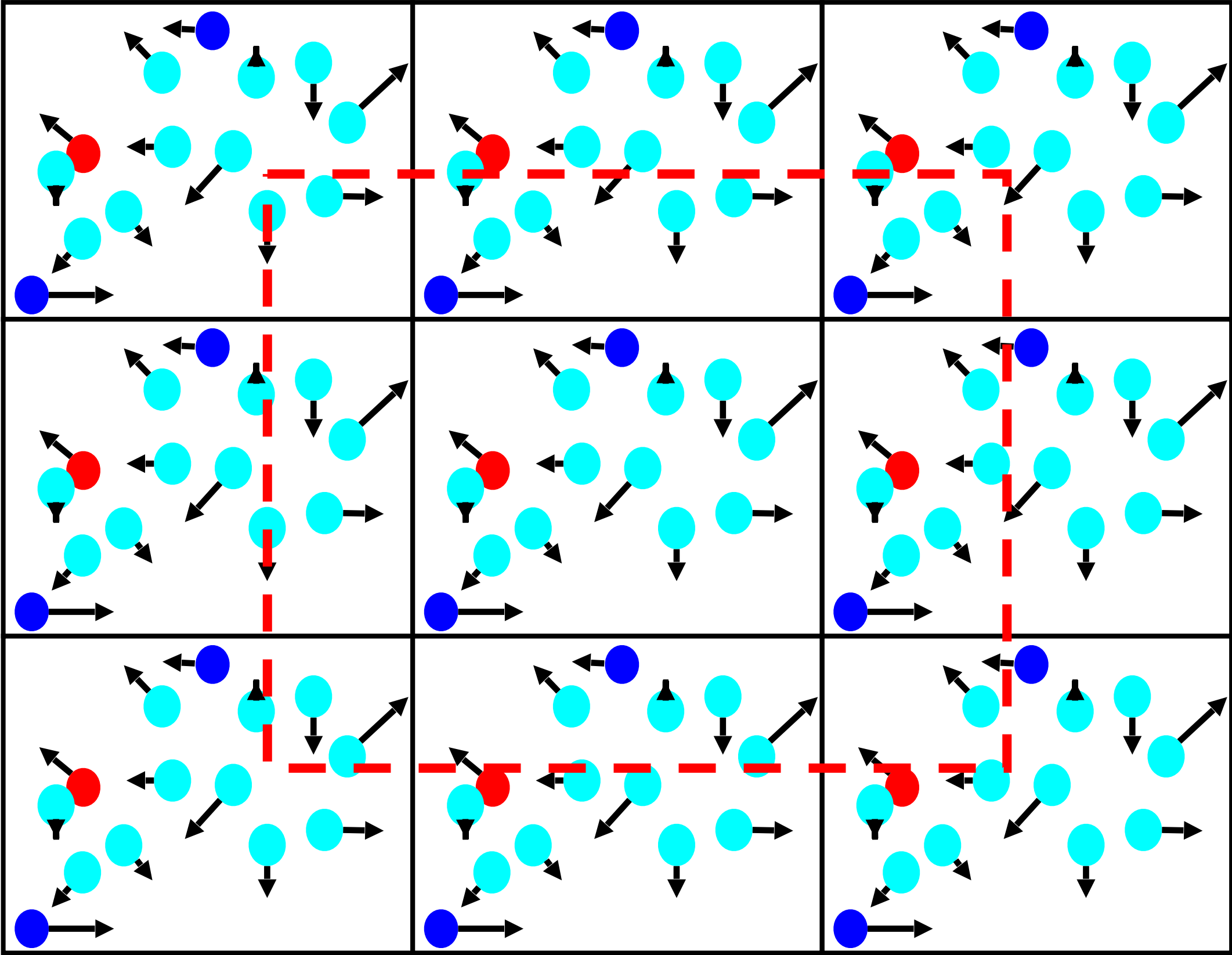
Problems: Density, pressure, boundary effects, ...

One solution: reflective boundaries, not quite good.

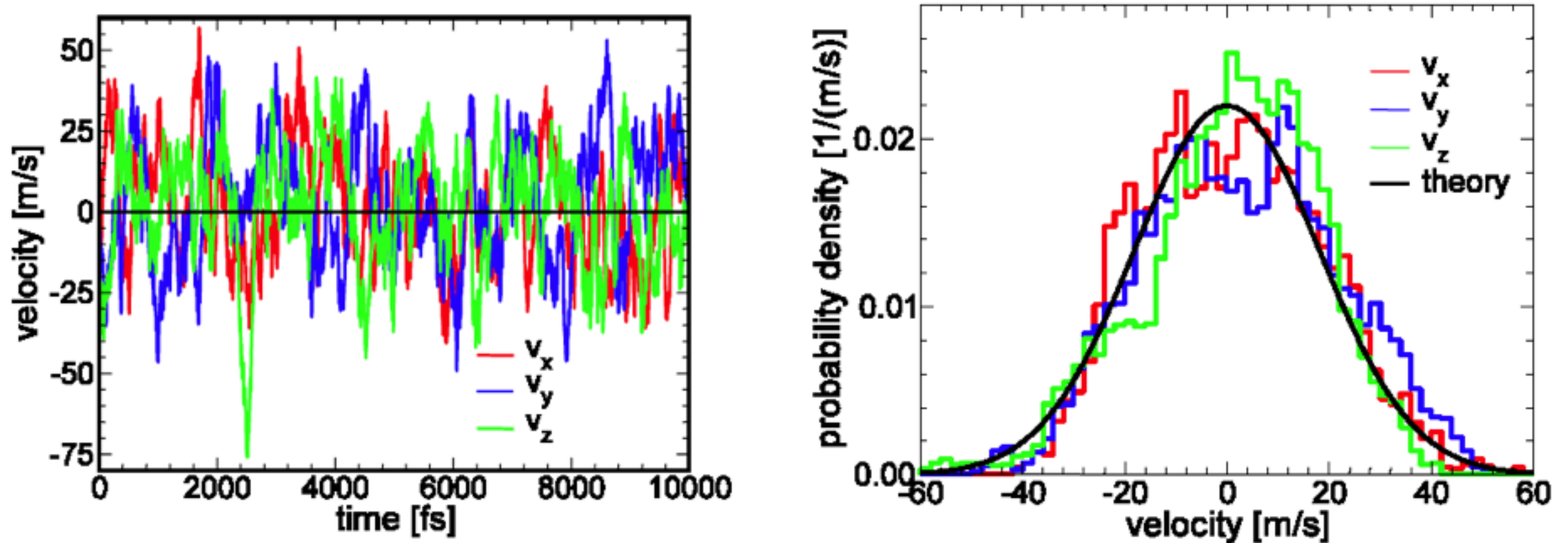
Spherical boundary conditions



Periodic Boundary Conditions



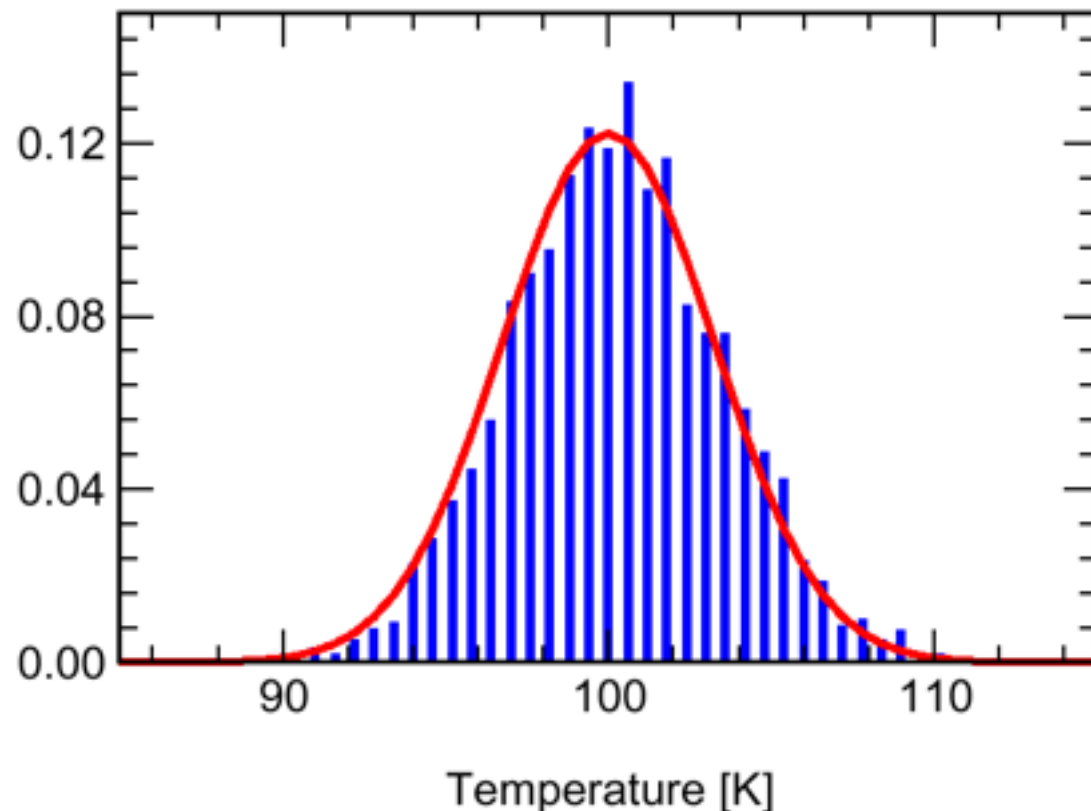
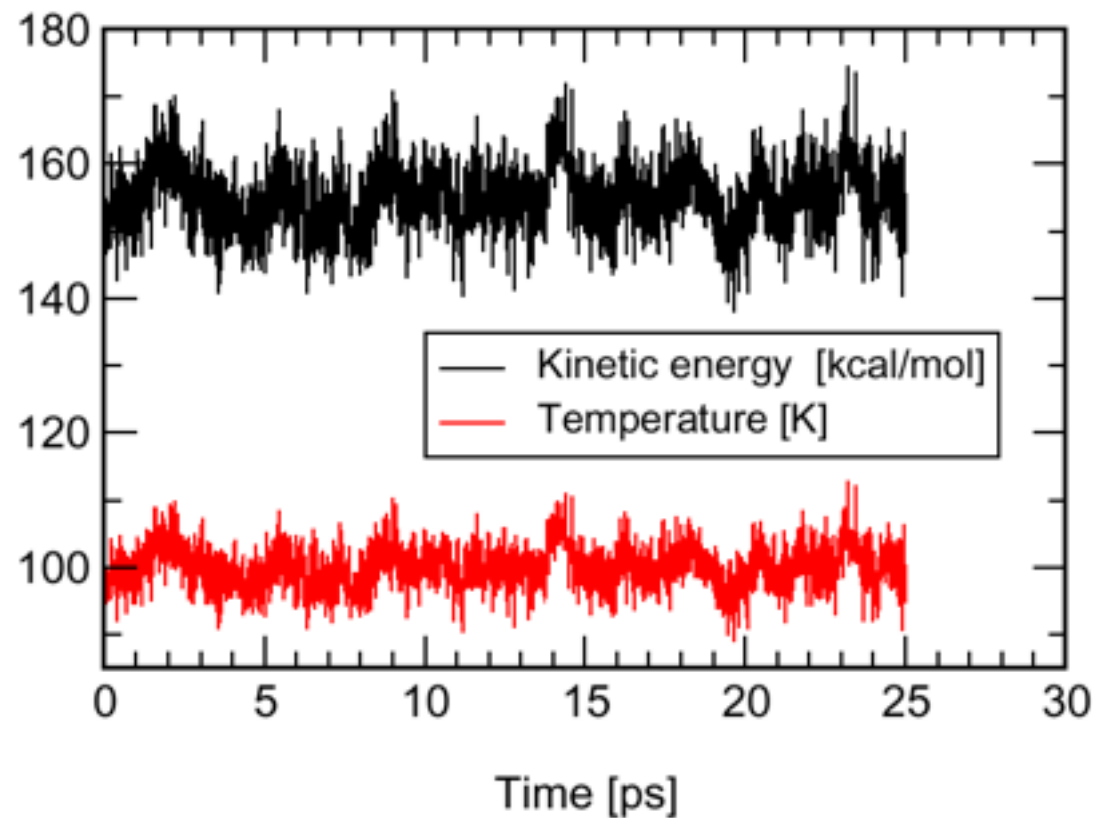
Maxwell Distribution of Atomic Velocities



$$p(v_{\sigma}) = \sqrt{\frac{m}{2\pi k_B T}} \exp \left[-\frac{mv_{\sigma}^2}{2k_B T} \right]$$

$$\sigma = x, y, z$$

Analysis of K , T (free dynamics)



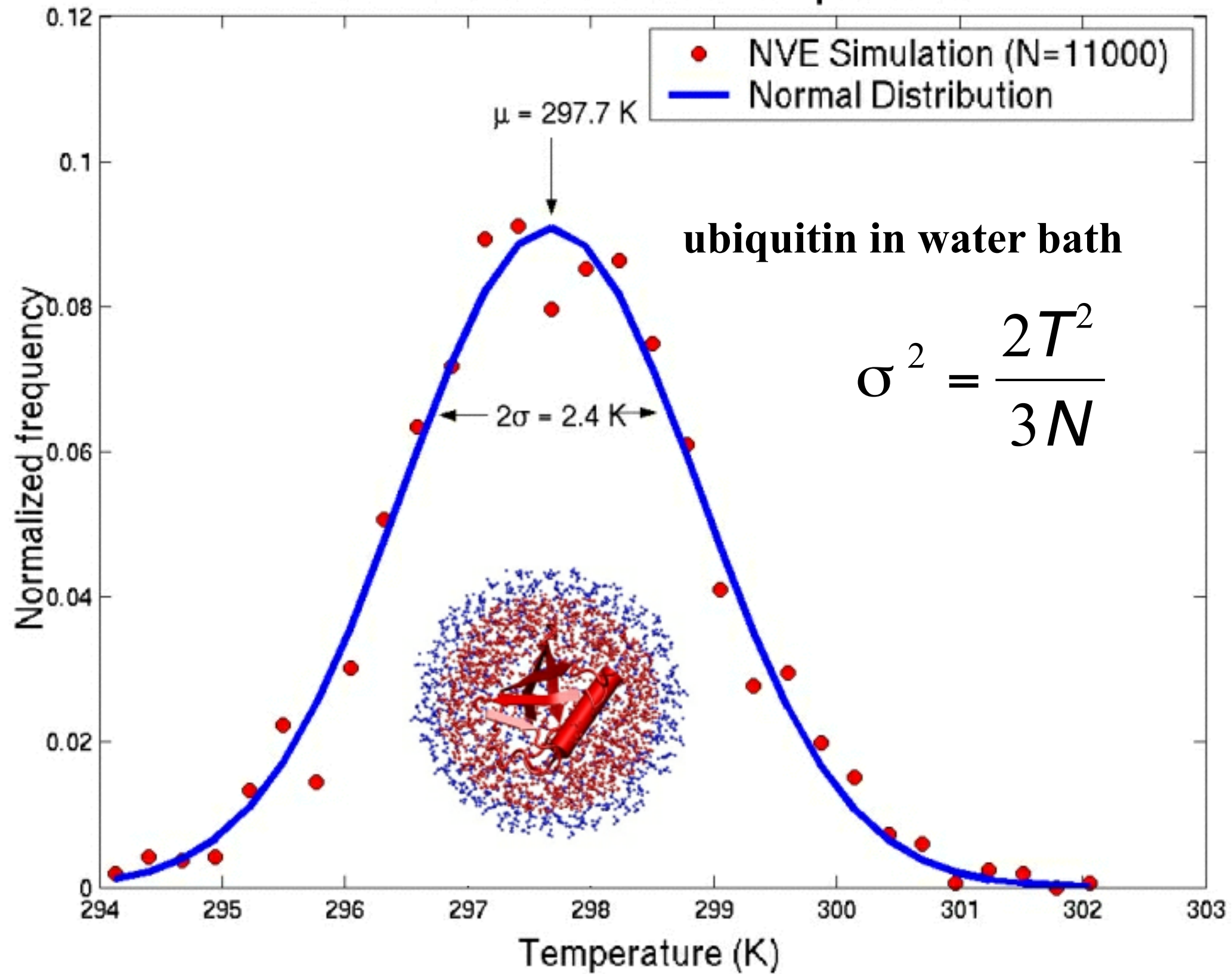
Definition of temperature

$$\left\langle \sum_j \frac{1}{2} m_j v_j^2 \right\rangle = \frac{3}{2} N k_B T$$

$$T = \frac{2}{3N k_B} \left\langle \sum_j \frac{1}{2} m_j v_j^2 \right\rangle$$

The atomic velocities of a protein establish a thermometer.

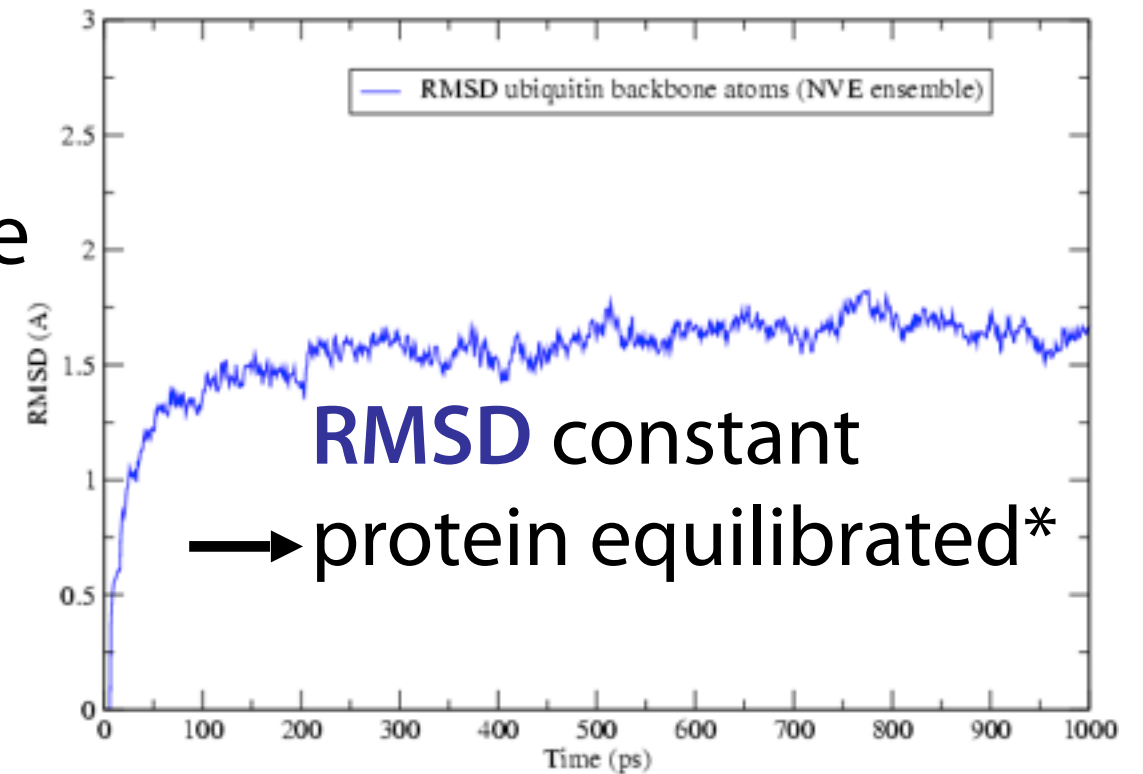
Normal Distribution of Temperatures



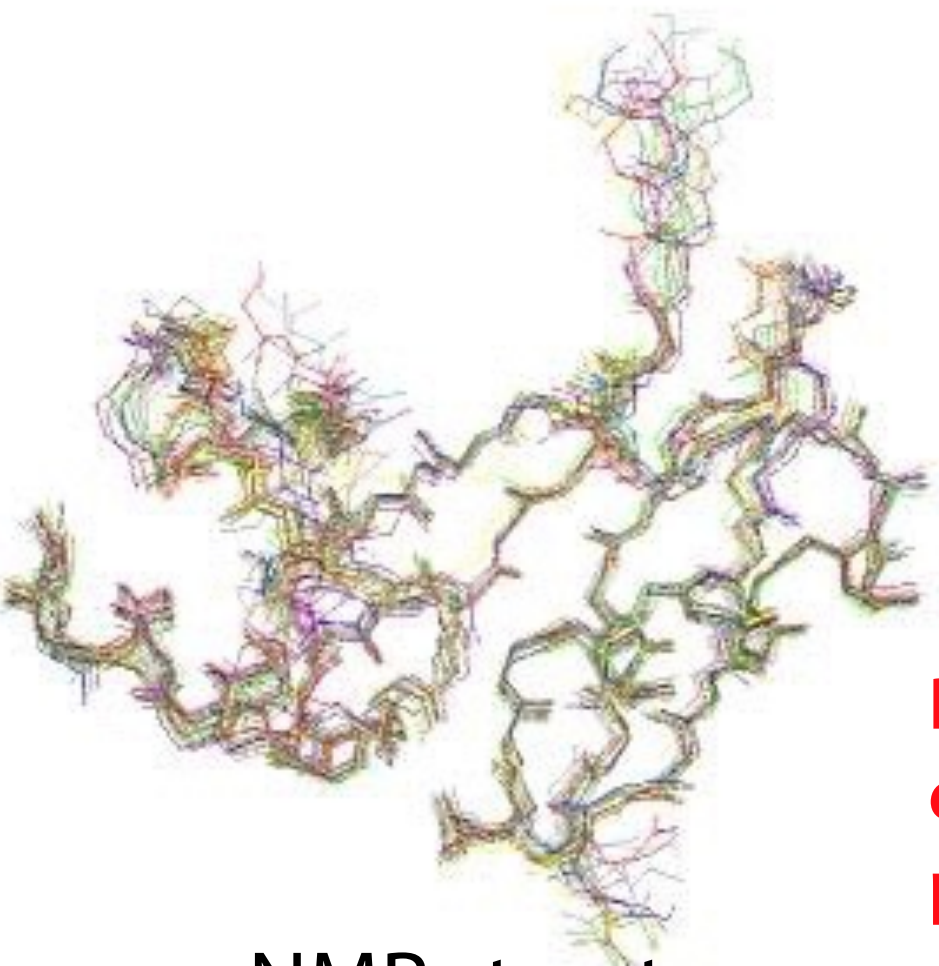
Equilibrium Properties of Proteins

Root Mean Squared Deviation: measure for equilibration and protein flexibility

$$RMSD(t) = \sqrt{\frac{1}{N} \sum_{i=1}^N (R_i(t) - R_i(0))^2}$$



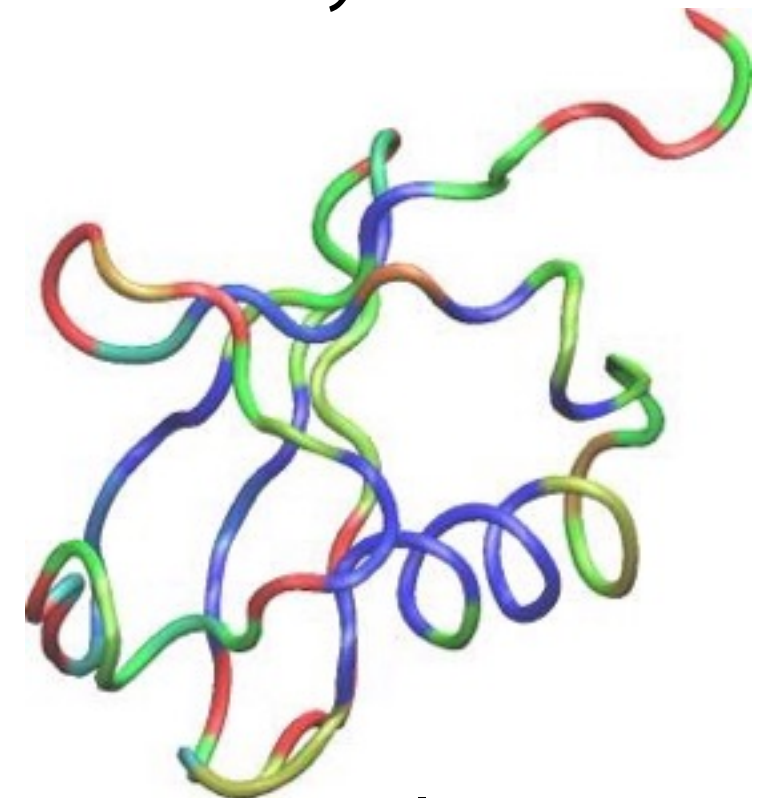
**you never really know for sure*



Ubiquitin

**Protein sequence
exhibits characteristic
permanent flexibility!**

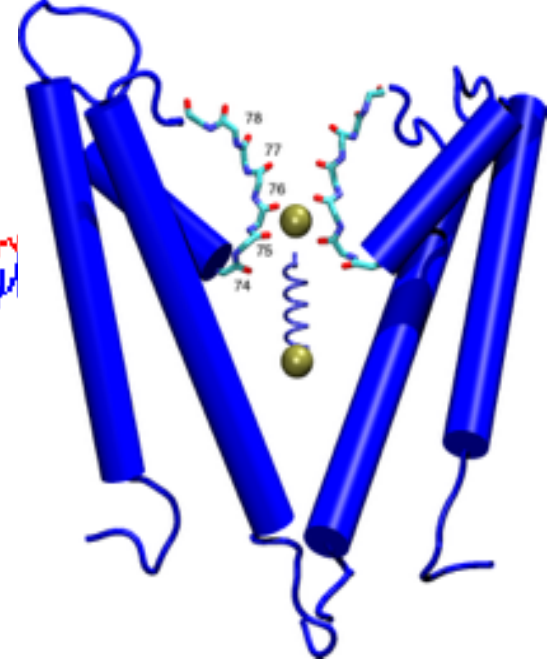
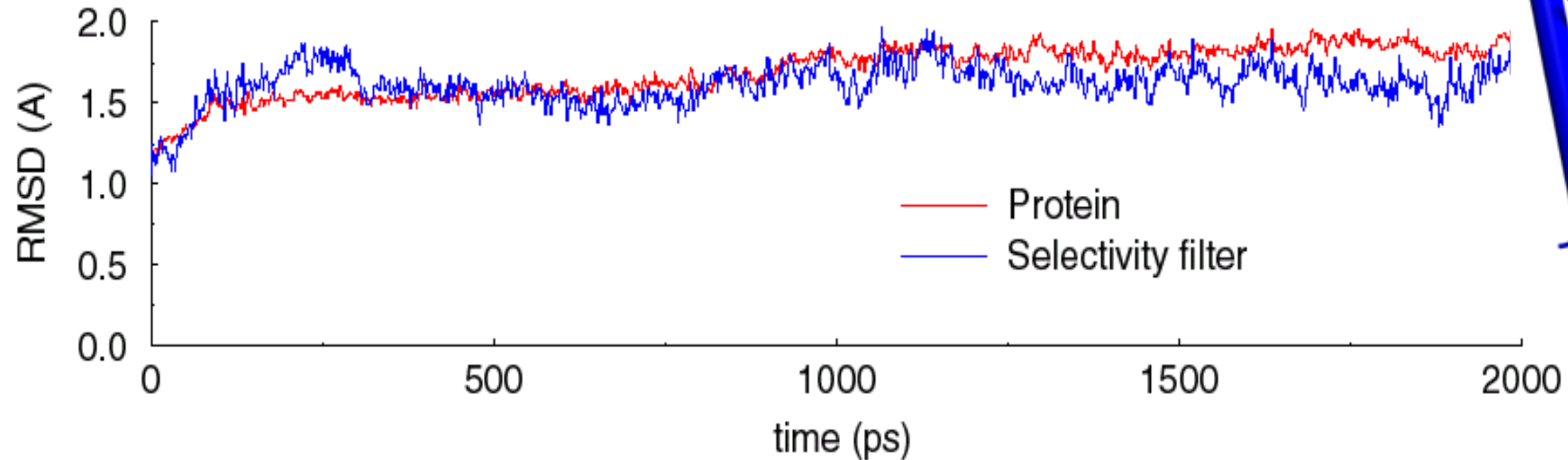
NMR structures
aligned together to see flexibility



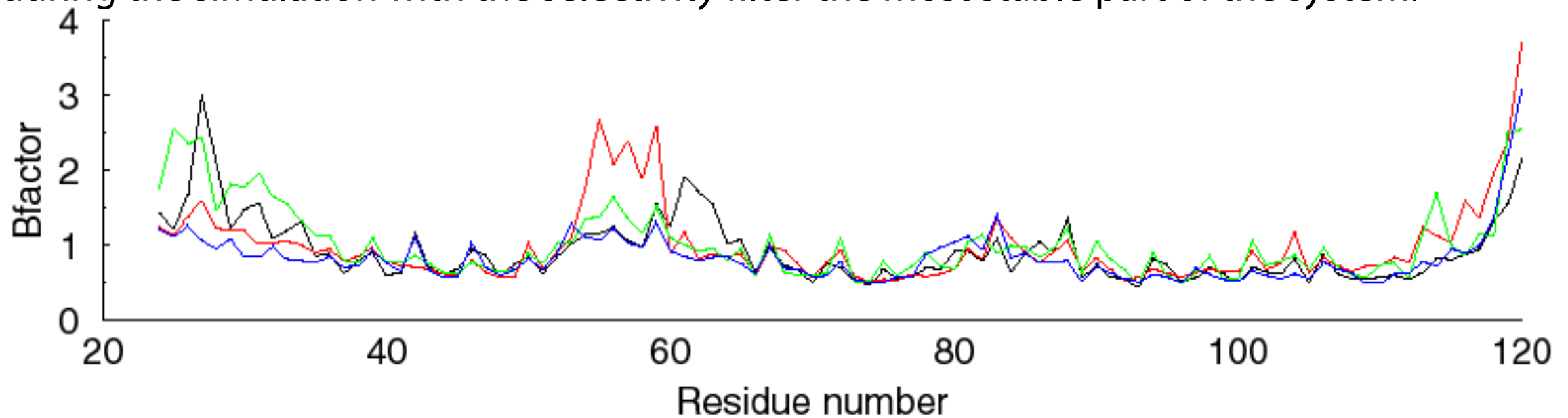
MD simulation

The color represents mobility of the protein through simulation (red = more flexible)

MD Results



RMS deviations for the KcsA protein and its selectivity filter indicate that the protein is stable during the simulation with the selectivity filter the most stable part of the system.



Temperature factors for individual residues in the four monomers of the KcsA channel protein indicate that the most flexible parts of the protein are the N and C terminal ends, residues 52-60 and residues 84-90. Residues 74-80 in the selectivity filter have low temperature factors and are very stable during the simulation.